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Nonslipping Contact Between a Mismatch Film and a Finite-Thickness Graded Substrate

The contact behavior of an elastic film subjected to a mismatch strain on a finite-thickness graded substrate is investigated, in which the contact interface is assumed to be nonslipping and the shear modulus of the substrate varies exponentially in the thickness direction. The Fourier transform method is adopted in order to reduce the governing partial differential equations to integral ones. With the help of numerical calculation, the interfacial shear stress, the internal normal stress in the film and the stress intensity factors are predicted for cases with different material parameters and geometric ones, including the modulus ratio of the film to the substrate, the inhomogeneous feature of the graded substrate, as well as the profile of the contacting film. All the physical predictions can be used to estimate the potential failure modes of the film–substrate interface. Furthermore, it is found that the result of a finite-thickness model is significantly different from the prediction of a generally adopted half-plane one. The study should be helpful for the design of film–substrate systems in real applications. [DOI: 10.1115/1.4031936]

Keywords: contact mechanics, nonslipping contact, thin film, graded substrate, finite-thickness

1 Introduction

The mechanical behavior of film/substrate systems is continuously concerned due to the important application in a large number of fields in recent decades [1]. For example, the system of inorganic films on soft substrates is now a hot topic, which can be used to monitor blood pressure during exercise or fabricate electronic eye cameras [2]. A model of film/substrate system can also be adopted in cell mechanics to investigate the interaction between cells and the surrounding substrates [3]. A precise estimation of the interface behavior in a film/substrate system is therefore an essential issue related to many real applications.

Two typically theoretical models have been proposed to study the interface mechanics of film/substrate systems. The first one represented by Akisanya and Fleck [4] and Yu et al [5] can be called as a fracture model, in which a pre-existing crack is introduced at the edge of a thin film. Of course, a model free from cracks or defects would also be crucial in many practical applications [6]. The second one may be addressed as a contact model, in which contact mechanics is used to find the stress field at the interface and near the edge of the film in order to judge possible forming cracks. The governing integrodifferential equations in the film–substrate contact model belong to singular ones, whose exact close-form solutions do not exist [7]. An infinite power series solution was given by Arutiunian [8] for the contact problem between a finite-length stiffener and a half-plane and an alternative method was to express the governing singular equation in terms of the interfacial shear stress [9]. Later, Erodgan and Civelek [10] explored the contact problem between an elastic reinforcement and a homogeneous elastic plate. Multilayered/multiperiodic films bonded to elastic substrates were discussed by Erodgan and Joseph [11,12]. Using the finite difference method, Hu [13] analyzed the film-edge-induced stress in substrates. Shield and Kim [14] discussed the problem of a thin film bonded to an elastic half space by the beam theory considering both normal and shear stresses. A closed form solution to a film bonded a homogeneous

substrate was obtained by Alaca et al. [6] with the assumption of a nearly rectangular film profile and the Vekua's solution procedure of Prandtl's equation.

All the above models involve homogeneous materials. Recently, graded materials have attracted numerous attentions due to the novel and excellent performance [15]. The traditional contact problem between a stamp (punch) and a graded medium has been explored extensively by many researchers to find the contact stress that may cause crack initiation on the contact surface, such as Giannakopoulos and Suresh [16,17], Choi and Paulino [18], Guler and Erodgan [19,20], Ke and Wang [21,22], Mao et al. [23], etc. Considering the interface adhesive effect, the problem of a rigid cylinder/sphere contacting a graded half-space was successfully discussed by Chen et al. [24,25]. The corresponding nonslipping adhesive model and the axis-symmetrically one were solved elegantly by Guo et al. [26,27].

A more complex but practical model is a film bonded on a graded substrate, especially for a deformable film. Guler et al. [7,28] investigated the problem of a film bonded on a graded and graded coated half-space using the approach of a singular integral equation. Whether the result of a film bonded on a graded substrate with an infinite thickness is proper for systems with a finite thickness in practical applications needs further discussion, since a significant thickness effect has been found in the problem of an elastic reinforcement contacting a homogeneous elastic plate by Erodgan and Civelek [10]. Does the thickness of substrates affect the interface behavior of a film–substrate system?

In this paper, a nonslipping contact model is established in order to answer the above questions, in which the mechanical deformation of the film is modeled as a mismatch strain for convenience and the thickness of the graded substrate is finite. The governing equation is formulated analytically, which can be reduced to an integral one and further solved numerically. The interfacial shear stress, the internal normal stress in the film, and the singularity near the contact edges are mainly focused in order to discuss the interface behavior of the film–substrate system.

2 Contact Model of a Deformable Film and a Finite Thickness Graded Substrate

The two-dimensional contact model of a deformable film bonded a finite thickness graded substrate with a rigid foundation

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Contributed by the Applied Mechanics Division of ASME for publication in the JOURNAL OF APPLIED MECHANICS. Manuscript received October 7, 2015; final manuscript received November 1, 2015; published online November 13, 2015. Editor: Yonggang Huang.

is shown in Fig. 1, where the thickness of the film is h_f and the length is $l_f = 2a$. μ_f is the shear modulus of the film and the Poisson's ratio is ν_f . Here, the thickness of the film h_f is assumed to be a constant, which can also be a function of the x axis [6]. The finite thickness of the graded substrate is h and the shear modulus $\mu_2(y)$ is expressed as

$$\mu_2(y) = \mu_1 \exp(\gamma y) \quad (1)$$

where the subscript "1" denotes the surface layer of the graded substrate, "2" the inner of the substrate. γ is a constant characterizing the substrate material's inhomogeneity, which can be obtained as

$$\gamma = \frac{1}{h} \ln\left(\frac{\mu_3}{\mu_1}\right) \quad (2)$$

Here, the subscript "3" denotes the lower boundary layer of the substrate. The substrate is assumed to have a same constant Poisson's ratio as the film due to a negligible Poisson's ratio effect [16].

3 Governing Equation of the Boundary Value Problem

For the present plane contact problem, the constitutive relation of the graded substrate in terms of the displacement components $u(x, y)$ and $v(x, y)$ in x and y directions can be given as

$$\sigma_{xx} = \frac{\mu_2}{\kappa - 1} \left[(1 + \kappa) \frac{\partial u}{\partial x} + (3 - \kappa) \frac{\partial v}{\partial y} \right] \quad (3a)$$

$$\sigma_{yy} = \frac{\mu_2}{\kappa - 1} \left[(1 + \kappa) \frac{\partial v}{\partial y} + (3 - \kappa) \frac{\partial u}{\partial x} \right] \quad (3b)$$

$$\sigma_{xy} = \mu_2 \left[\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right] \quad (3c)$$

where σ_{ij} (i or $j = x$ or y) denotes stress components, $\kappa = 3 - 4\nu$ for the plane strain case and $\kappa = (3 - \nu)/(1 + \nu)$ for the plane stress one.

Substituting Eqs. (3a)–(3c) into the classically equilibrium equation yields

$$\begin{aligned} (\kappa + 1) \frac{\partial^2 u}{\partial x^2} + (\kappa - 1) \frac{\partial^2 u}{\partial y^2} + 2 \frac{\partial^2 v}{\partial x \partial y} + \gamma(\kappa - 1) \frac{\partial u}{\partial y} \\ + \gamma(\kappa - 1) \frac{\partial v}{\partial x} = 0 \end{aligned} \quad (4a)$$

$$\begin{aligned} (\kappa - 1) \frac{\partial^2 v}{\partial x^2} + (\kappa + 1) \frac{\partial^2 v}{\partial y^2} + 2 \frac{\partial^2 u}{\partial x \partial y} + \gamma(3 - \kappa) \frac{\partial u}{\partial x} \\ + \gamma(\kappa + 1) \frac{\partial v}{\partial y} = 0 \end{aligned} \quad (4b)$$

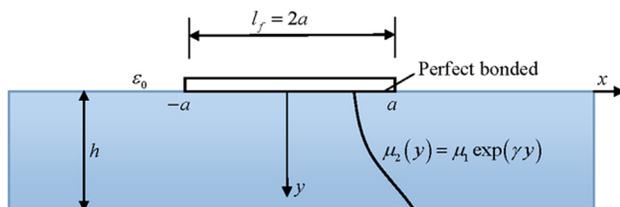


Fig. 1 The perfectly adhesive contact model between an elastic film of length l_f and a finite-thickness graded substrate. The variation of environmental temperature induces a mismatched strain at the contact interface. The shear modulus of the substrate abides by a graded law $\mu_2(y) = \mu_1 \exp(\gamma y)$; h is the thickness of the graded substrate.

The equation of displacements $u(x, 0)$ and $v(x, 0)$ at the upper surface of the graded substrate can be derived after Fourier transformation to x and a rather lengthy mathematic procedure similar to Chen and Chen [29,30]. Here, we skip the details and present only the final formula.

$$\begin{aligned} \frac{\partial u(x, 0)}{\partial x} = \frac{\kappa - 1}{4\mu_1} \sigma_{yy}(x, 0) - \frac{\kappa + 1}{4\pi\mu_1} \int_{-a}^a \frac{\sigma_{xy}(r, 0)}{r - x} dr \\ + \frac{1}{\pi} \int_{-a}^a [K_{11}(x, r) \sigma_{xy}(r, 0) + K_{12}(x, r) \sigma_{yy}(r, 0)] dr \end{aligned} \quad (5a)$$

$$\begin{aligned} \frac{\partial v(x, 0)}{\partial x} = -\frac{\kappa - 1}{4\mu_1} \sigma_{xy}(x, 0) - \frac{\kappa + 1}{4\pi\mu_1} \int_{-a}^a \frac{\sigma_{yy}(r, 0)}{r - x} dr \\ + \frac{1}{\pi} \int_{-a}^a [K_{21}(x, r) \sigma_{xy}(r, 0) + K_{22}(x, r) \sigma_{yy}(r, 0)] dr \end{aligned} \quad (5b)$$

where the kernels $K_{ij}(x, r)$ are bounded functions given in the Appendix.

The film thickness h_f is assumed to be sufficiently small and the film contacts the graded substrate without any initial stress. Only the shear stress σ_{xy} is produced at the contact interface when a mismatch strain is exerted, i.e.,

$$\sigma_{xy}^f(x, 0) = \sigma_{xy}(x, 0) = \begin{cases} -f(x), & -a < x < a \\ 0, & x < -a, x > a \end{cases} \quad (6a)$$

The normal stress in the film in the direction of the y axis is

$$\sigma_{yy}^f(x, 0) = \sigma_{yy}(x, 0) = 0 \quad (6b)$$

Due to a small thickness of the film, the normal stress in the film in the direction of the x axis is assumed to be uniform across the thickness as shown in Fig. 2. The equilibrium equation in the film can be given as

$$\int_{-a}^x \sigma_{xy}^f(t, 0) dt + \sigma_{xx}^f(x, 0) h_f = 0 \quad (7)$$

and

$$\int_{-a}^a \sigma_{xy}^f(t, 0) dt = - \int_{-a}^a f(t) dt = 0 \quad (8)$$

Then, the normal stress in the film can be written as

$$\sigma_{xx}^f(x, 0) = \frac{1}{h_f} \int_{-a}^x f(t) dt \quad (9)$$

and the strain in the film induced by the interface stress is

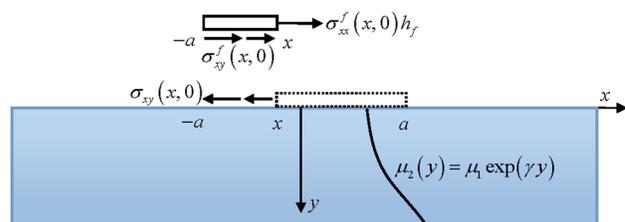


Fig. 2 Schematic of the mechanical behavior of the bonded interface between a deformable film and an elastically graded substrate

$$e_{xx}^f(x, 0) = \frac{\partial u^f}{\partial x} = \frac{1 + \kappa_f}{8\mu_f} \sigma_{xx}^f(x, 0) \quad (10)$$

i.e.

$$e_{xx}^f(x, 0) = \frac{1 + \kappa_f}{8\mu_f} \frac{1}{h_f} \int_{-a}^x f(t) dt \quad (11)$$

where $\kappa_f = 3 - 4\nu_f$ for the plane strain problem and $\kappa_f = (3 - \nu_f)/(1 + \nu_f)$ for the plane stress one.

Using Eq. (5) leads to the strain at the upper surface of the graded substrate

$$e_{xx}(x, 0) = \frac{\partial u(x, 0)}{\partial x} = \frac{\kappa + 1}{4\pi\mu_1} \int_{-a}^a \left[\frac{1}{t-x} - \frac{4\mu_1}{\kappa + 1} K_{11}(x, t) \right] f(t) dt \quad (12)$$

Now consider a film in nonslipping contact with a finite-thickness graded substrate as a reference state. Then, a mismatch strain ε_0 is imposed at the interface, which may be induced by a sudden change

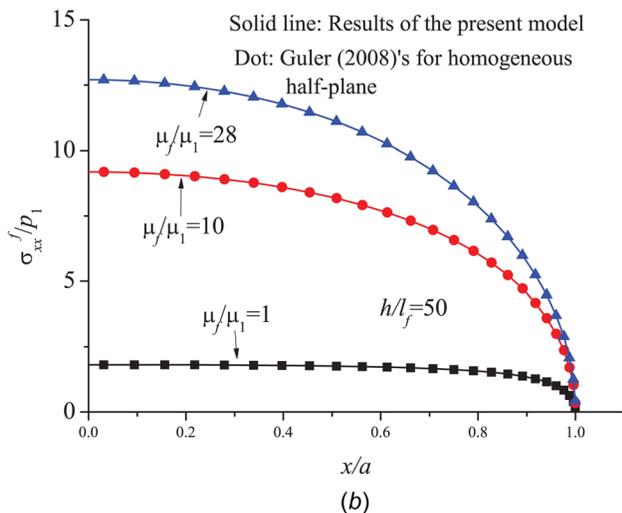
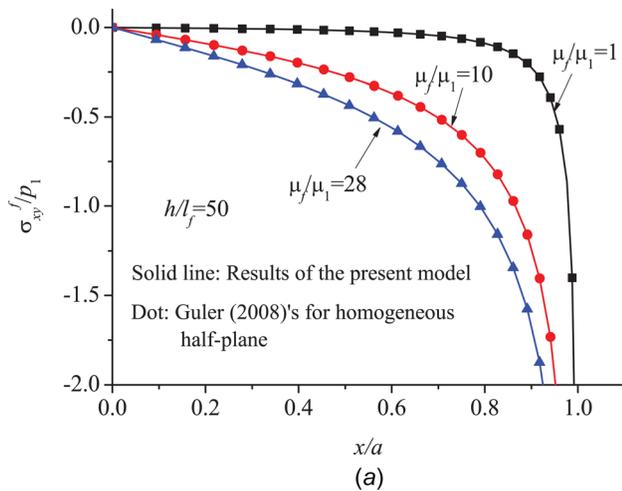


Fig. 3 Comparison of the distribution of the nondimensional interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 for the model of an elastic film in adhesive contact with a homogeneous half-plane and the present one with a relatively thick substrate, where the ratio of the film length to its thickness is $l_f/h_f=32$, μ_f is the shear modulus of the film and μ_1 denotes the shear modulus of the upper surface of the graded substrate

in temperature in the environment. We have $\varepsilon_0 = (\alpha_1 - \alpha_2)\Delta T$, where α_i denotes the thermal expansion coefficient of the film or the substrate, ΔT is the temperature change. The assumption of perfect bonding yields the compatibility condition at the contact interface

$$\varepsilon_{xx}(x, 0) - e_{xx}^f(x, 0) = \varepsilon_0 \quad (13)$$

Substituting Eqs. (11) and (12) into Eq. (13) yields

$$\frac{\kappa + 1}{4\pi\mu_1} \int_{-a}^a \left[\frac{1}{t-x} - \frac{4\mu_1}{\kappa + 1} K_{11}(x, t) \right] f(t) dt - \frac{1 + \kappa_f}{8\mu_f} \frac{1}{h_f} \int_{-a}^x f(t) dt = \varepsilon_0 \quad (14)$$

which is the governing equation of the present model.

4 Solution of the Integral Equation

By introducing the following normalized quantities:

$$t = a\tau, \quad x = as \quad (15)$$

Equations (14) and (8) can be rewritten as

$$\frac{1}{\pi} \int_{-1}^1 \left[\frac{1}{\tau-s} - \frac{4\mu_1 a}{\kappa + 1} K_{11}(s, \tau) \right] f(\tau) d\tau - \frac{4\mu_1}{\kappa + 1} \frac{1 + \kappa_f}{8\mu_f} \frac{a}{h_f} \int_{-1}^s f(\tau) d\tau = \varepsilon_0 \quad (16)$$

$$f(\tau) d\tau = \frac{4\mu_1}{\kappa + 1} \varepsilon_0$$

$$\int_{-1}^1 f(s) ds = 0 \quad (17)$$

Considering the Cauchy-type singular kernel in the integral equation, the solution to Eqs. (16) and (17) may be expressed as [31]

$$f(s) = \frac{1}{\sqrt{1-s^2}} \sum_{n=0}^{\infty} A_n T_n(s) \quad (18)$$

where $T_n(s)$ denotes the Chebyshev polynomial of the first kind of order n , A_n is an unknown constant to be determined.

Substituting Eq. (18) into Eqs. (16) and (17) yields

$$\frac{1}{\pi} \int_{-1}^1 \sum_{n=0}^{\infty} A_n \frac{T_n(\tau)}{\sqrt{1-\tau^2}} \left[\frac{1}{\tau-s} - \frac{4\mu_1 a}{\kappa + 1} K_{11}(s, \tau) \right] d\tau - \frac{4\mu_1 a}{\kappa + 1} \frac{1 + \kappa_f}{8\mu_f h_f} \int_{-1}^s \sum_{n=0}^{\infty} A_n \frac{T_n(\tau)}{\sqrt{1-\tau^2}} d\tau = \frac{4\mu_1}{\kappa + 1} \varepsilon_0 \quad (19)$$

$$\int_{-1}^1 \sum_{n=0}^{\infty} A_n \frac{T_n(s)}{\sqrt{1-s^2}} ds = 0 \quad (20)$$

Considering the following properties of the Chebyshev polynomials:

$$\int_{-1}^1 \frac{T_n(s)T_m(s)}{\sqrt{1-s^2}} ds = \begin{cases} 0, & m \neq n \\ 1/2, & m = n \geq 1 \\ 1, & m = n = 1 \end{cases} \quad (21)$$

$$\int_{-1}^1 \frac{T_n(\tau)}{(\tau-s)\sqrt{1-\tau^2}} d\tau = \begin{cases} 0, & n = 0 \\ U_{n-1}(s), & n > 0 \end{cases} \quad (22)$$

and

$$\int_{-1}^s \frac{T_n(\tau)}{(\tau-s)\sqrt{1-\tau^2}} d\tau = -\frac{1}{n} U_{n-1}(s) \sqrt{1-s^2} \quad (23)$$

one can readily obtain $A_0 = 0$ and the integral equation may be rewritten as

$$\sum_{n=1}^{\infty} A_n [U_{n-1}(s) + Q_n(s)] = \frac{4\mu_1}{1+\kappa} \varepsilon_0 \quad (24)$$

where

$$Q_n(s) = -\frac{1}{\pi} \int_{-1}^1 \frac{4\mu_1 a}{\kappa+1} K_{11}(s, \tau) \frac{T_n(\tau)}{\sqrt{1-\tau^2}} d\tau + \frac{4\mu_1 a}{\kappa+1} \frac{1+\kappa_f}{8\mu_f h_f} \frac{1}{n} U_{n-1}(s) \sqrt{1-s^2} \quad (25)$$

Truncating the series in Eq. (18) at $n = N$, and selecting proper collocation points as roots of the following Chebyshev polynomials [32]

$$T_N(\tau_i) = 0 \quad (26)$$

the integral equation will be reduced to linear algebraic equations with N unknown constants A_n ($n = 1, 2, \dots, N$).

Based on the solution of Eqs. (24) and (25), the shear stress at the contact surface can be approximately given as

$$\sigma_{xy}^f(x, 0) = -\frac{4\mu_1 \varepsilon_0}{1+\kappa} \sum_{n=1}^N A_n T_n\left(\frac{x}{a}\right) \frac{1}{\sqrt{1-(x/a)^2}} \quad (27)$$

Then, the normal stress in the film can be obtained

$$\sigma_{xx}^f(x, 0) = \frac{4\mu_1 \varepsilon_0 a}{1+\kappa} \sum_{n=1}^N \frac{1}{n} A_n U_{n-1}\left(\frac{x}{a}\right) \sqrt{1-(x/a)^2} \quad (28)$$

The stress intensity factor near the left contact edge can be defined as

$$K_{II} = \lim_{x \rightarrow -a} \sigma_{xy}^f(x, 0) \sqrt{2(x+a)} \quad (29)$$

Substituting Eq. (27) into Eq. (29) yields the stress intensity factor at the left contact edge

$$\frac{K_{II}}{\frac{4\mu_1 \varepsilon_0}{1+\kappa} \sqrt{a}} = -\sum_{n=1}^N A_n T_n(-1) \quad (30)$$

5 Special Case

If we set $\gamma = 0$ and $a \ll h$ in the present model, the graded substrate will be reduced to a homogeneous one and can be looked as a half space. Then, the model is degraded to be a special one, i.e., an elastic film in contact with a homogeneous half space, which has well investigated by Guler [28]. Comparison of the present model and the Guler's one is shown in Fig. 3, where the distribution of the nondimensional interface shear stress σ_{xy}^f/p_1 and the nondimensional normal stress in the film σ_{xx}^f/p_1 is given in Figs. 3(a) and 3(b), respectively, where $p_1 = (4\mu_1 \varepsilon_0 / (1+\kappa))$. The shear modulus of the homogeneous half space is denoted as μ_1 and the ratio of the film length to the thickness of the film l_f/h_f is 32. One can see the prediction of both models agrees very well with each other not only for the system of a soft film on a half space but also for the one of a stiff film on the half space.

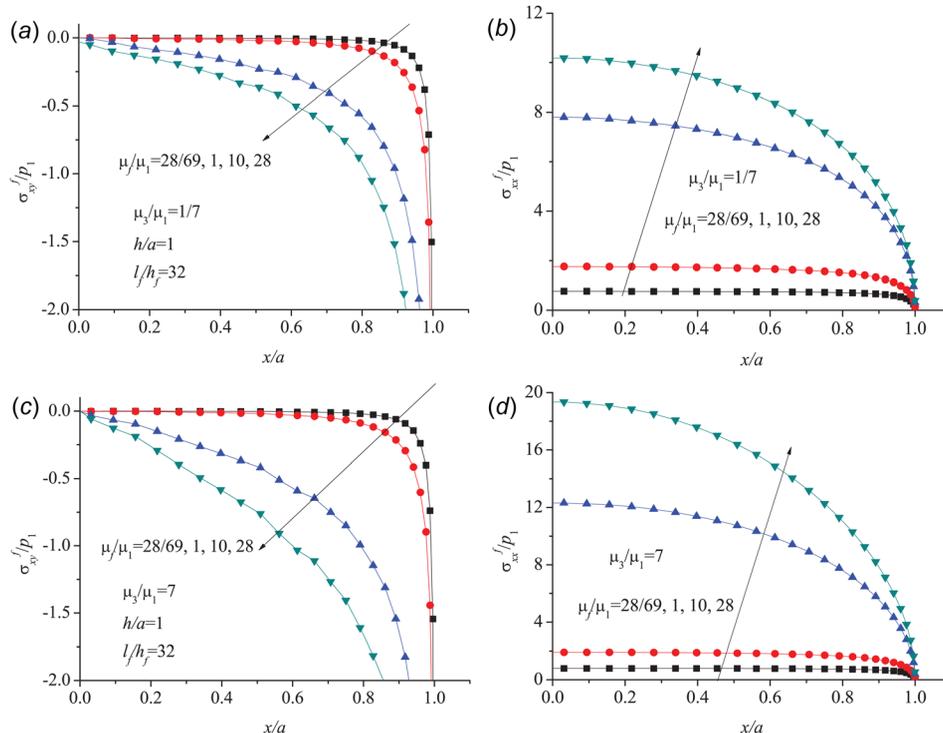


Fig. 4 The distribution of the nondimensional interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 for the model of a deformable film in adhesive contact with a finite-thickness graded substrate with determined parameters $h/a = 1$ and $l_f/h_f = 32$, but with different ratios μ_f/μ_1 . (a) and (b) for $\mu_3/\mu_1 = 1/7$; (c) and (d) for $\mu_3/\mu_1 = 7$.

Table 1 The stress intensity factor $K_{II}/p_1\sqrt{a}$ near the contact edge versus different ratios μ_f/μ_1 in cases with fixed parameters $h/a=1$, $l_f/h_f=32$, and μ_3/μ_1

	$\mu_f/\mu_1 = 28/69$	$\mu_f/\mu_1 = 1$	$\mu_f/\mu_1 = 10$	$\mu_f/\mu_1 = 28$
$K_{II}/p_1\sqrt{a}$ ($\mu_3/\mu_1 = 7$)	0.1782	0.2793	0.7995	1.0831
$K_{II}/p_1\sqrt{a}$ ($\mu_3/\mu_1 = 1/7$)	0.1764	0.2722	0.6476	0.7735

6 Results and Discussion

In the following analysis, the Poisson's ratio of both the film and the substrate is $\nu = 0.3$. Only half of the contact area is analyzed due to the symmetric feature of the model.

The film-substrate system with different parameter combinations is investigated. Figures 4(a) and 4(b) give the distribution of the nondimensional interface shear stress σ_{xy}^f/p_1 and the nondimensional normal stress in the film σ_{xx}^f/p_1 with a fixed $\mu_3/\mu_1 = 1/7$ but for different ratios μ_f/μ_1 , where $p_1 = (4\mu_1\varepsilon_0/1 + \kappa)$. Figure 4(a) shows that the interface shear stress is obviously singular near the contact edges and the singularity increases with an increasing ratio μ_f/μ_1 , which can be also

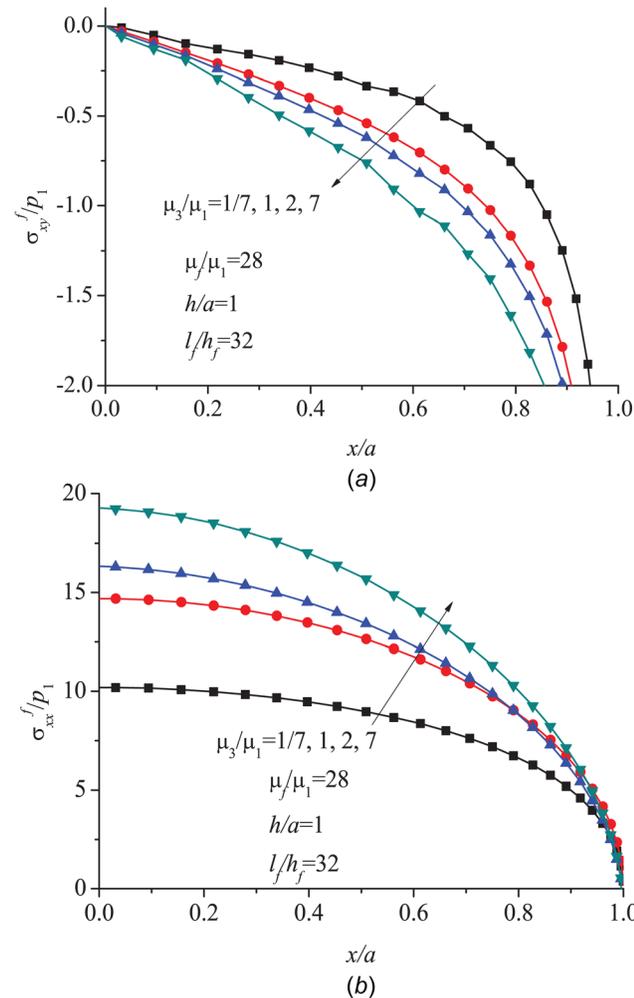


Fig. 5 The distribution of the interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 for the model of a deformable film in adhesive contact with a finite-thickness graded substrate with determined parameters $h/a=1$, $\mu_f/\mu_1=28$, $l_f/h_f=32$, but with different inhomogeneity parameters μ_3/μ_1 of the substrate

Table 2 The stress intensity factors $K_{II}/p_1\sqrt{a}$ near the contact edge versus different ratios μ_3/μ_1 in cases with determined parameters $\mu_f/\mu_1=28$, $h/a=1$, and $l_f/h_f=32$

	$\mu_3/\mu_1 = 1/7$	$\mu_3/\mu_1 = 1$	$\mu_3/\mu_1 = 2$	$\mu_3/\mu_1 = 7$
$K_{II}/p_1\sqrt{a}$	0.7735	0.9409	0.9955	1.0831

found from the variation of the Mode II stress intensity factor $K_{II}/p_1\sqrt{a}$ shown in Table 1. The normal stress in the film distributes nonuniformly in the length direction, the maximum value always emerges at the symmetric axis and the normal stress in the film vanishes at the end as shown in Fig. 4(b), which is consistent with the boundary condition. For different ratios μ_f/μ_1 , the normal stress in the film increases with an increasing value of μ_f/μ_1 . Both the singularity near the contact edge and the normal stress in the film demonstrate that an interface between a relatively soft film and a determined substrate is much reliable. Figures 4(c) and 4(d) give the distribution of the nondimensional interface shear stress σ_{xy}^f/p_1 and the nondimensional normal stress in the film σ_{xx}^f/p_1 for different ratios μ_f/μ_1 but with another fixed value $\mu_3/\mu_1 = 7$, which means an increasing stiffness along with the depth of the graded substrate. Comparing to Figs. 4(a) and 4(b), it is interesting to find that the varying features of both the shear stress singularity near the contact end and the normal stress in the film do not change.

For a fixed ratio $\mu_f/\mu_1 = 28$, the distribution of the interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 is shown in Figs. 5(a) and 5(b), respectively, with different values of the inhomogeneity parameter μ_3/μ_1 . It is found that both the interface shear stress and the normal one in the film increase with an increasing μ_3/μ_1 . Table 2 gives the Mode II stress intensity factor $K_{II}/p_1\sqrt{a}$, which is obviously enhanced by an increasing μ_3/μ_1 . That is to say, for a determined film, a substrate with a gradually decreasing stiffness in the depth direction would be advantageous for resisting interface crack initiation. Such a rule should be very useful for the design of new structures and has already been adopted by many nature materials, for example, bone marrow, whose stiffness becomes larger and larger from the inside to the outside.

Figure 6 gives the effect of thickness of the graded substrate on the distribution of the interface shear stress and the normal stress in the film with a fixed ratio $\mu_f/\mu_1 = 28$. Figures 6(a) and 6(b) correspond to cases of a finite-thickness substrate with an increasing stiffness in the depth direction $\mu_3/\mu_1 = 7$, while Figs. 6(c) and 6(d) correspond to the one with the stiffness decreasing in the depth direction $\mu_3/\mu_1 = 1/7$. The substrate thickness shows a more regular effect on the interface shear stress and the normal stress in the film in the former case, where both stresses increase monotonically with a decreasing thickness of the graded substrate. However, nonmonotonically variation in the interface shear stress and the normal stress in the film can be found in Figs. 6(c) and 6(d). Variation of the Mode II stress intensity factor $K_{II}/p_1\sqrt{a}$ versus the nondimensional substrate thickness h/a is given in Fig. 7 for different μ_3/μ_1 , where the result of a homogeneous half-plane case is also shown for comparison. Obviously, $K_{II}/p_1\sqrt{a}$ increases monotonically with a decreasing thickness h/a in cases of $\mu_3/\mu_1 \geq 1$, while decreases first and then increases for the case of $\mu_3/\mu_1 = 1/7$. All the results demonstrate that the effect of substrate thickness on the interface behavior can not be neglected. A semi-infinite model may not be precise in explaining phenomena found in a finite film-substrate structure. However, if the substrate thickness is large enough, the boundary effect should be weakened in all cases as shown in Fig. 7, both a half-plane model and a finite one may replace each other.

The ratio effect of the film length to its thickness in the finite model on the distribution of the nondimensional interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 is shown in Fig. 8. A fixed value $\mu_3/\mu_1 = 1/7$ is adopted in Figs. 8(a) and

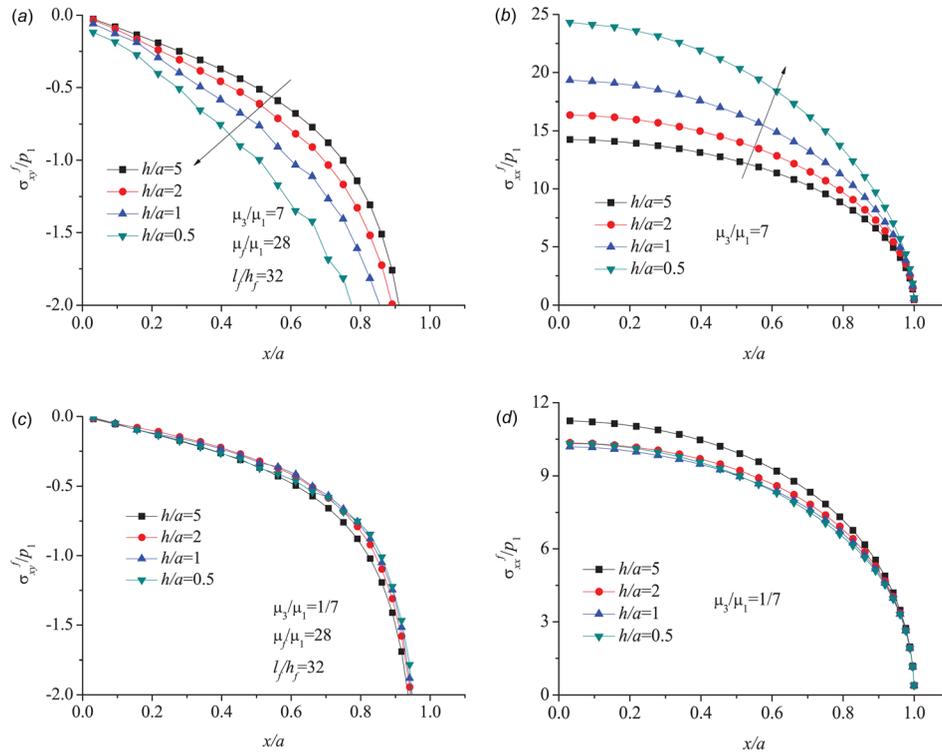


Fig. 6 The distribution of the interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 for the model of a deformable film in adhesive contact with a finite-thickness graded substrate with determined parameters $\mu_f/\mu_1 = 28$, $l_f/h_f = 32$, but with different ratios h/a . (a) and (b) for $\mu_3/\mu_1 = 7$; (c) and (d) for $\mu_3/\mu_1 = 1/7$.

8(b), while $\mu_3/\mu_1 = 7$ is used in Figs. 8(c) and 8(d). Both cases show that the interface shear stress decreases with an increasing ratio of l_f/h_f . The corresponding Mode II stress intensity factor $K_{II}/p_1\sqrt{a}$ is given in Table 3, which also exhibits a decreasing singularity with an increasing ratio l_f/h_f . It means that, in cases with a fixed film length, the thinner the film, the stronger the interface strength is; while for a fixed film thickness, the shorter the film, the more possible failure at the contact end it is. However, in cases with a fixed film thickness, Figs. 8(b) and 8(d) show that the normal stress in the film increases with an increasing film length no matter μ_3/μ_1 is larger or smaller than 1. Therefore, a longer film is more likely to break at site of the symmetric axis than a shorter one in cases of films with a fixed thickness.

All the above analysis is based on an assumption of a uniform film thickness. If the profile of the film in the thickness direction is a function of the x axis, the present model may be transformed to be the model by Alaca et al. [6]. In Alaca et al. [6], the profile function of the film in the thickness direction was adopted as

$$h_f(x) = h_{f0} \sqrt{1 - \frac{x^2}{a^2}} \left(1 + \lambda \frac{x^2}{a^2} \right) \quad (31)$$

where h_{f0} denotes the film thickness at $x = 0$. λ is a shape parameter. Tuning the parameter λ would lead to different film profiles as shown in Fig. 9, which may change from a convex shape to a concave one with an increasing λ . With $\lambda = 0.9$, $\gamma = 0$ and $h \gg a$, the present model can be reduced to the one by Alaca et al. [6], where the stress intensity factor was defined as

$$K_2 = \lim_{x \rightarrow -a} \sigma_{xy}^f(x, 0) \sqrt{2\pi(x+a)} \quad (32)$$

Substituting $\mu_f/\mu_1 = 28$, $l_f/h_{f0} = 32$ into Eq. (32) and using the thickness profile in Eq. (31) yield $K_2 = 0.4683$, which is consistent well with the result of Alaca et al. The requirement that the thickness of the film is much less than its length should be satisfied since the normal stress in the film is assumed to distribute uniformly in the thickness direction.

7 Summary

A nonslipping contact model between a mismatched film and a finite-thickness graded substrate is investigated in the present paper. Fourier transform method is used to transform partial differential equations to the integral ones. Distributions of the interface shear stress, the normal stress in the film, and the stress intensity factors near the contact edge are analyzed. It is found that the interface behavior would be influenced significantly by material

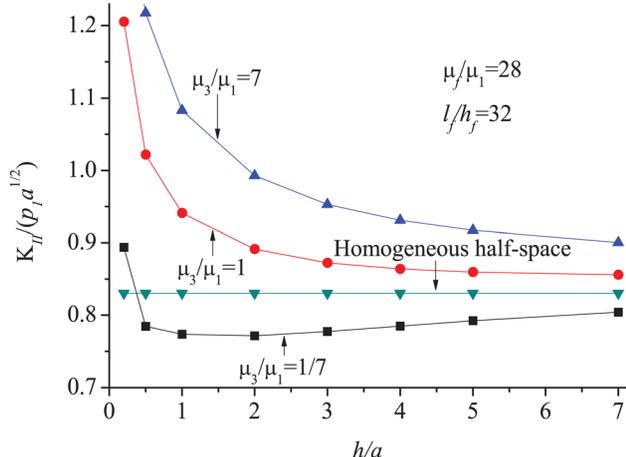


Fig. 7 Variation of the Mode II stress intensity factor $K_{II}/p_1\sqrt{a}$ versus the ratio of the substrate thickness to half of the film length h/a for different ratios μ_3/μ_1 , where $\mu_f/\mu_1 = 28$

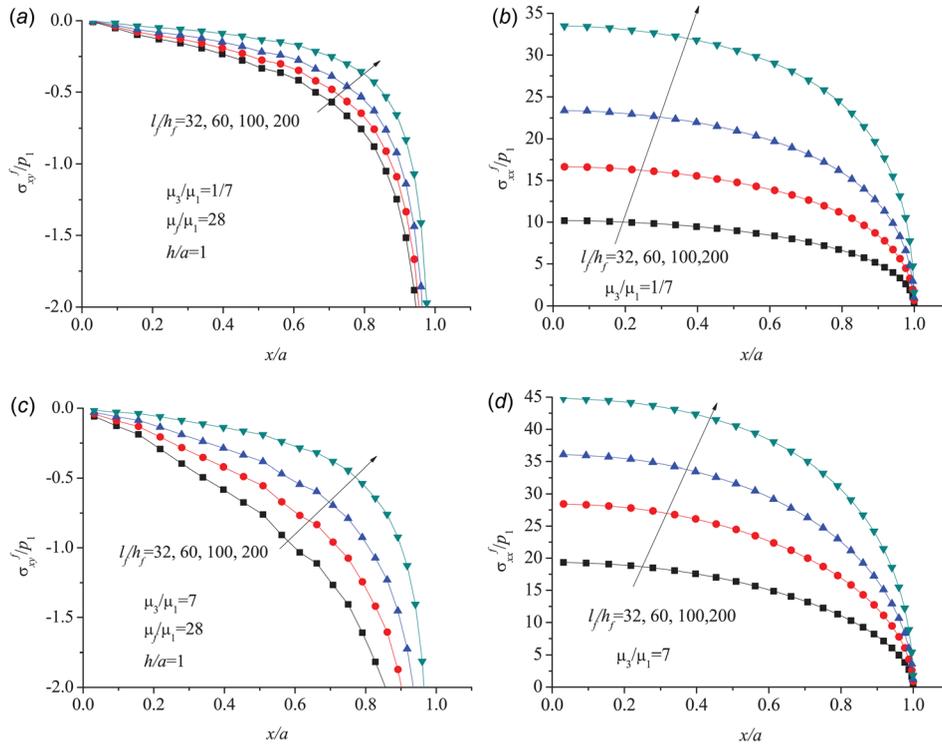


Fig. 8 The distribution of the interface shear stress σ_{xy}^f/p_1 and the normal stress in the film σ_{xx}^f/p_1 for the model of a deformable film in adhesive contact with a finite-thickness graded substrate with determined parameters $\mu_f/\mu_1 = 28$, $h/a=1$, but with different values l_f/h_f . (a) and (b) for $\mu_3/\mu_1 = 1/7$; (c) and (d) for $\mu_3/\mu_1 = 7$.

Table 3 The stress intensity factors $K_{II}/p_1\sqrt{a}$ near the contact edge versus different ratios l_f/h_f in cases with determined parameters $\mu_f/\mu_1 = 28$, $h/a=1$, and μ_3/μ_1

	$l_f/h_f = 32$	$l_f/h_f = 60$	$l_f/h_f = 100$	$l_f/h_f = 200$
$K_{II}/p_1\sqrt{a}$ ($\mu_3/\mu_1 = 7$)	1.0831	0.9166	0.7672	0.5738
$K_{II}/p_1\sqrt{a}$ ($\mu_3/\mu_1 = 1/7$)	0.7735	0.7048	0.6304	0.5113

parameters of the film and the graded substrate. Appropriate choice of the film stiffness, the stiffness variation law of the graded substrate, as well as the length scales of the film and substrate, could improve resistance of an interface crack initiation or

the film breakage. The results should be helpful for the design of film–substrate systems in practical applications.

Acknowledgment

P.C. thanks the support of China Postdoctoral Science Foundation Funded Project (No. 2015M580486), NSFC (No. 11402292), Natural Science Foundation of Jiangsu Province (No. BK20140179), Jiangsu Planned Projects for Postdoctoral Research Funds (No. 1501009B), the Fundamental Research Funds for the Central Universities (No. 2014QNA73). S.C. gratefully acknowledges the support by the NSFC (Nos. 11372317 and 11125211), the Nano-project (2012CB937500), and the CAS/SAFEA International Partnership Program for Creative Research Teams.

Appendix

The kernels $K_{ij}(x, r)$ in Eqs. (5a) and (5b) are expressed as

$$K_{11}(x, r) = \int_0^{+\infty} \left[\alpha N_{11}(\alpha) + \frac{\kappa + 1}{4\mu_1} \right] \sin[\alpha(r - x)] d\alpha \quad (A1a)$$

$$K_{12}(x, r) = -i \int_0^{+\infty} \left[\alpha N_{12}(\alpha) - i \frac{\kappa - 1}{4\mu_1} \right] \cos[\alpha(r - x)] d\alpha \quad (A1b)$$

$$K_{21}(x, r) = -i \int_0^{+\infty} \left[\alpha N_{21}(\alpha) + i \frac{\kappa - 1}{4\mu_1} \right] \cos[\alpha(r - x)] d\alpha \quad (A1c)$$

$$K_{22}(x, r) = \int_0^{+\infty} \left[\alpha N_{22}(\alpha) + \frac{\kappa + 1}{4\mu_1} \right] \sin[\alpha(r - x)] d\alpha \quad (A1d)$$

where $N_{jk}(\alpha)$ ($j, k = 1, 2$) are the corresponding four elements in matrix $N(\alpha)$

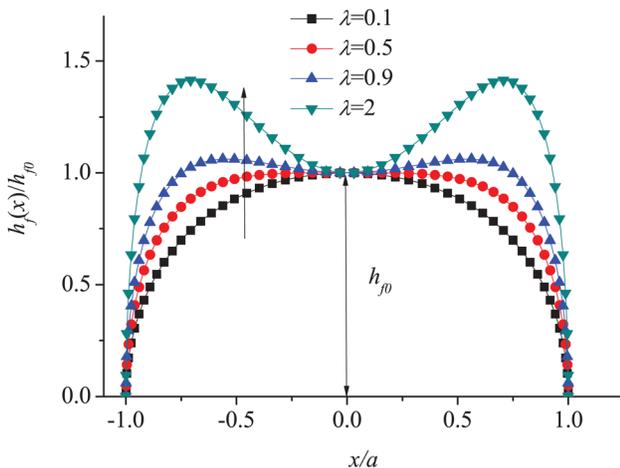


Fig. 9 The profile of the bonded film in the thickness direction with different shape parameters λ

$$\mathbf{N}(\alpha) = \begin{bmatrix} m_1 & m_2 \\ 1 & 1 \end{bmatrix} \cdot \left[\begin{array}{c} 1/(m_1 n_1 - i\alpha) \\ 1/\{(\kappa - 1)[-i\alpha m_1(3 - \kappa) + n_1(1 + \kappa)]\} \\ 1/(m_2 n_2 - i\alpha) \\ 1/\{(\kappa - 1)[-i\alpha m_2(3 - \kappa) + n_2(1 + \kappa)]\} \end{array} \right]^{-1} \quad (\text{A2})$$

Coefficients in Eq. (A2) are as follows:

$$n_j(\alpha) = -\frac{\gamma}{2} + \sqrt{\frac{\gamma^2}{4} + \alpha^2 - i(-1)^j \gamma \alpha \left(\frac{3 - \kappa}{1 + \kappa}\right)^{1/2}}, \quad \text{Re}(n_j) > 0, \quad j = 1, 2 \quad (\text{A3})$$

$$n_j(\alpha) = -\frac{\gamma}{2} - \sqrt{\frac{\gamma^2}{4} + \alpha^2 + i(-1)^j \gamma \alpha \left(\frac{3 - \kappa}{1 + \kappa}\right)^{1/2}}, \quad \text{Re}(n_j) < 0, \quad j = 3, 4 \quad (\text{A4})$$

and $m_j(\alpha)$ for each $n_j(\alpha)$ ($j = 1, \dots, 4$) is

$$m_j(s) = \frac{(\kappa - 1)(n_j^2 + \gamma n_j) - \alpha^2(\kappa + 1)}{\alpha[2n_j + \gamma(\kappa - 1)]} \quad (\text{A5})$$

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