

Static compression of iron-silicon alloys: Implications for silicon in the Earth's core

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[1] Three iron-silicon alloys ($\text{Fe}_{85}\text{Si}_{15}$, $\text{Fe}_{71}\text{Si}_{29}$, and $\epsilon\text{-FeSi}$) have been studied in a diamond anvil cell at room temperature up to 55 GPa by in situ energy-dispersive X-ray diffraction techniques. A body centered cubic (bcc) to hexagonal close packed (hcp) phase transformation in $\text{Fe}_{85}\text{Si}_{15}$ began at 16 GPa and was completed by 36 GPa. No phase transformations were observed in either $\text{Fe}_{71}\text{Si}_{29}$ or $\epsilon\text{-FeSi}$ at high pressures, even when laser-heated to about 2000 K. The isothermal bulk modulus (K_{0T}) of hcp- $\text{Fe}_{85}\text{Si}_{15}$ is 141 (± 10) GPa with $K'_{0T} = 5.70(\pm 0.60)$ and $V_{02} = 6.882(\pm 0.031)$ cm^3/mol (per molar atom). The K_{0T} of $\text{Fe}_{71}\text{Si}_{29}$ is 199.0 (± 5.3) GPa with $K'_{0T} = 5.66(\pm 0.61)$ and $V_0 = 6.887(\pm 0.014)$ cm^3/mol , and the K_{0T} of $\epsilon\text{-FeSi}$ is 184.7 (± 3.9) GPa with K'_{0T} of 4.75 (± 0.37) and $V_0 = 6.790(\pm 0.007)$ cm^3/mol . Our study indicates that the substitution of Si into iron would lower the density of iron, but significantly changes its compressibility neither in the bcc phase, nor at high pressures in the hcp phase. Upon comparison with the Preliminary Reference Earth Model, the calculated equations of state (EOS) of hcp- $\text{Fe}_{85}\text{Si}_{15}$, using the Mie-Grüneisen EOS, indicate that an outer core containing about 8–10 wt.% Si and inner core containing about 4 wt.% Si in iron would satisfy the seismological constraints. Addition of silicon into iron increases the bulk sound velocity of iron, consistent with silicon being a light element in the Earth's core. **INDEX TERMS:** 1015 Geochemistry: Composition of the core; 3919 Mineral Physics: Equations of state; 3924 Mineral Physics: High-pressure behavior; 3954 Mineral Physics: X ray, neutron, and electron spectroscopy and diffraction; **KEYWORDS:** high pressure, light elements, iron-silicon alloy, Earth's core, X-ray diffraction, equation of state

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1. Introduction

[2] Geophysical and cosmochemical evidence indicates that iron is the most abundant component in the Earth's core. The simplest interpretation of the core is an inner core composed of a crystalline iron-nickel alloy and an outer core of liquid iron-nickel alloyed with a small fraction of lighter element(s) [Birch, 1952]. The density of the outer core is about 10% lower than the density of iron at core pressures and temperatures, while the bulk sound velocity of the core is about 3% higher than that of liquid iron [Birch, 1952; Dziewonski and Anderson, 1981; Ahrens, 1982; Anderson and Ahrens, 1994]. The difference in density between the core and the high-pressure phase of iron indicates the presence of a low-atomic-weight component in the outer core [Birch, 1952]. There

is also evidence that the inner core is less dense than pure iron, and the amount of light element in the inner core may be as much as several weight percent [Jephcoat and Olson, 1987; Anderson and Ahrens, 1994; Stixrude et al., 1997; Fiquet et al., 2001; Mao et al., 2001]. A number of light elements (H, C, N, O, Mg, Si, S) have been considered by various workers [e.g., Poirier, 1994; Hillgren et al., 2000].

[3] Silicon has been proposed to be an important alloying element in the Earth's core, based on its cosmochemical abundance [Ringwood, 1966] and the thermoelastic properties of iron-silicon alloys in shock-wave experiments [Balchan and Cowan, 1966; Matassov, 1977]. Allègre et al. [1995] used the ratios of major and trace elements in the mantle and in meteorite groups to estimate the possible silicon content in the core, and concluded that the core may contain about 7.3 wt.% Si in addition to other light elements. Silicon solubility in liquid iron-rich metal increases with decreasing oxygen fugacity, increasing pressure, and increasing temperature, and a few weight percent of silicon may dissolve into liquid iron at the base of a deep magma ocean (~ 700 km depth, ~ 25 GPa) [Gessmann et al., 2001].

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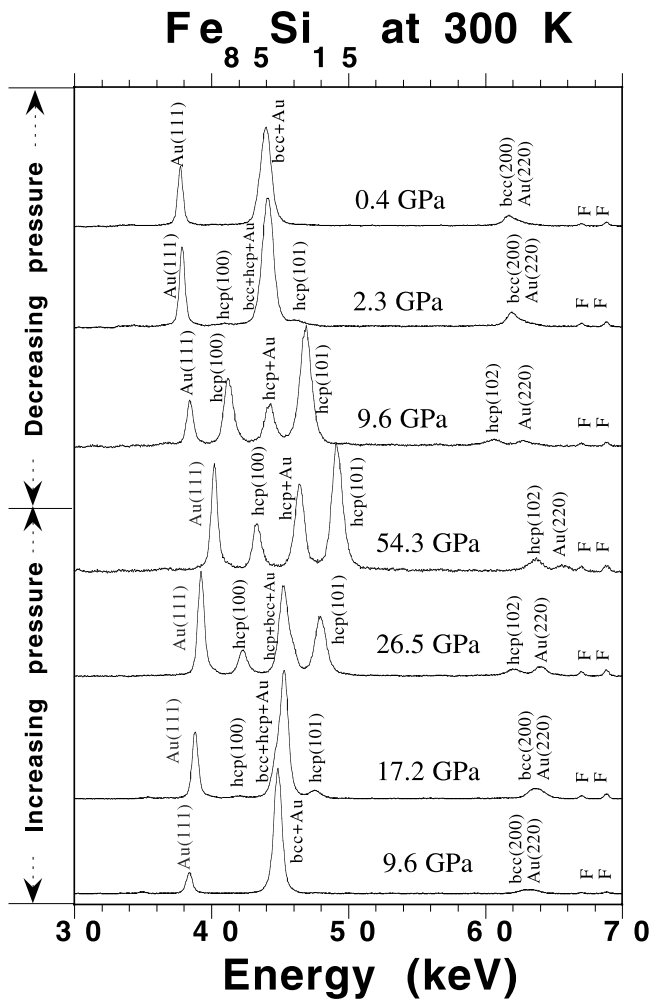


Figure 1. A series of X-ray diffraction patterns of $\text{Fe}_{85}\text{Si}_{15}$ at increasing pressures to 54.3 GPa and then at decreasing pressures at 300 K. Diffraction peaks from the bcc and hcp phases of the alloy are indicated. F, the fluorescence peaks of gold; Au, the internal pressure calibrant.

To generate a core containing ~ 7.3 wt.% Si as estimated by *Allègre et al.* [1995], a low oxygen fugacity during accretion [Ringwood, 1966] or possibly a magma ocean extending to 850 km (~ 30 GPa) [Gessmann et al., 2001] would be required. Studies of chemical reactions between mantle silicates and molten iron have also indicated that Fe-Si alloys may form at the core-mantle boundary (D'' zone) [Knittle and Jeanloz, 1989, 1991]. Shock experiments on Fe-Si alloys suggested that an outer core containing 14–20 wt.% Si in iron is sufficient to satisfy the density deficit of the core [Balchan and Cowan, 1966; Matassov, 1977].

[4] Nevertheless, silicon was excluded as the primary alloying element in the outer core based on the relatively high bulk modulus (K_0) and low-pressure derivative of the bulk modulus (K'_0) of an intermediate compound, iron silicide (ϵ -FeSi) [Knittle and Williams, 1995]. Since the sound velocity of liquid iron is essentially parallel to the Preliminary Reference Earth Model (PREM) velocity profile [Anderson and Ahrens, 1994], a lower value of K'_0 of iron-silicon alloy would not be able to compensate for the

effect of significant Si within the outer core [Williams and Knittle, 1997]. In contrast, a recent static study on the equation of state (EOS) of a body-centered cubic (bcc)- $\text{Fe}_{84}\text{Si}_{16}$ (in atomic percent) alloy up to 9 GPa and 773 K concluded that substitution of silicon in iron would not significantly change the compressibility of the iron-rich Fe-Si alloy [Zhang and Guyot, 1999].

[5] Therefore we studied the EOS of three Fe-Si alloys ($\text{Fe}_{85}\text{Si}_{15}$, $\text{Fe}_{71}\text{Si}_{29}$, and ϵ -FeSi) in a diamond anvil cell (DAC) to understand the elastic properties of various Fe-Si alloys under high pressures and to resolve the discrepancy between previous studies. We especially focused on the thermal EOS of iron-rich Fe-Si alloys because silicon alloyed with iron has a large effect upon the phase diagram of iron, and the iron-rich portion of the Fe-Si system is more applicable than ϵ -FeSi to understanding the possible effect of Si on iron under core conditions [Lin et al., 2002]. The thermal EOS and bulk sound

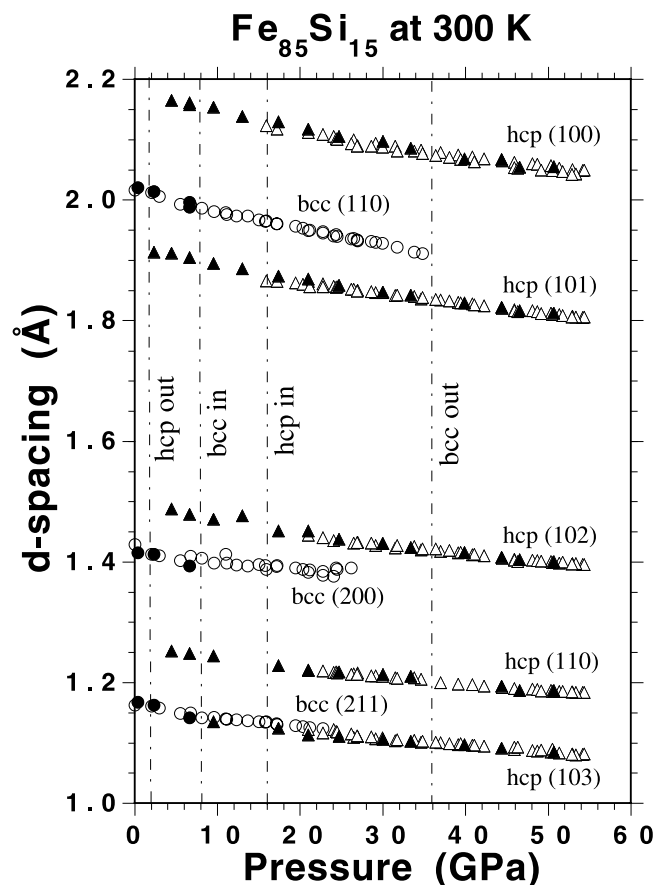


Figure 2. A plot of d -spacing versus pressure for $\text{Fe}_{85}\text{Si}_{15}$ alloy under high pressures at 300 K. When the pressure was elevated above 16 GPa, a transformation to a hcp phase was clearly observed (indicated as hcp in). This phase transition was completed by 36 GPa (indicated as bcc out). The first appearance of the bcc phase during the quench process occurred at about 8 GPa (bcc in), and the hcp phase can still be seen above 2 GPa (hcp out). Open circles, bcc phase in increasing pressure; solid circles, bcc phase in decreasing pressure; open triangles, hcp phase in increasing pressure; solid triangles, hcp phase in decreasing pressure.

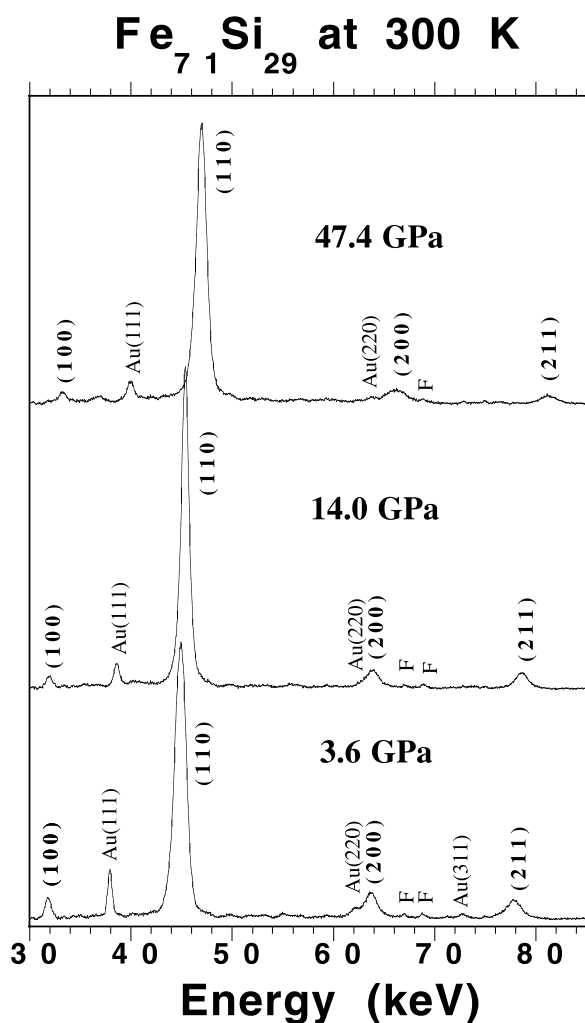


Figure 3. A series of X-ray diffraction patterns of $\text{Fe}_{71}\text{Si}_{29}$ at increasing pressures to 47.4 GPa at 300 K. Diffraction peaks from the $\text{Fe}_{71}\text{Si}_{29}$ phase are labeled in boldface. F, the fluorescence peaks of gold; Au, the internal pressure calibrant.

velocity of hexagonal close packed (hcp)- $\text{Fe}_{85}\text{Si}_{15}$ are calculated under core conditions, using our experimental data and the shock-wave data. The $\text{Fe}_{71}\text{Si}_{29}$ and ϵ -FeSi alloys were also laser heated in a DAC to understand the stability of the phases under high pressures and high temperatures.

2. Experiments

[6] The starting materials, $\text{Fe}_{85}\text{Si}_{15}$ alloy (lot FE166010/6) and $\text{Fe}_{71}\text{Si}_{29}$ alloy (lot FE176010/5), were purchased from Goodfellow Corporation and ϵ -FeSi (iron silicide; lot C18G21) was obtained from Johnson Matthey. Electron microprobe analyses showed that the starting materials contained 7.9 (± 0.3) wt.% Si, 17.0 (± 0.2) wt.% Si, and 33.5 (± 0.3) wt.% Si (each averaged from at least five analyses), respectively. Angle-dispersive X-ray diffraction at ambient conditions showed that the $\text{Fe}_{85}\text{Si}_{15}$ was in bcc structure with $a = 2.8520(\pm 0.0015)$ Å, $\text{Fe}_{71}\text{Si}_{29}$ was in the

ordered B2 structure (space group: $Pm\bar{3}m$) with $a = 2.8108(\pm 0.0020)$ Å, and ϵ -FeSi was in the B20 structure (space group: $P2_13$) with $a = 4.4846(\pm 0.0015)$ Å [Pauling and Soldate, 1948; Khalaff and Schubert, 1974; Kusbaschewski, 1982]. The B20 structure can be viewed as a distortion of the NaCl structure along [111] directions [Pauling and Soldate, 1948].

[7] Stainless steel gaskets were preindented to a thickness of 30 μm and then drilled with a 100- μm hole. The sample was ground into a fine powder with an average grain size of 2 μm with gold powder added as the pressure standard [Heinz and Jeanloz, 1984a], and loaded into a DAC with ethanol-methanol (1:4) as the pressure medium. The ethanol-methanol pressure medium remains hydrostatic to almost 10 GPa [Piermarini *et al.*, 1973]; at high pressures it becomes a quasi-hydrostatic solid [Piermarini *et al.*, 1975]. Energy-dispersive synchrotron X-ray diffraction in a DAC was performed at Beamlines 13IDD and 13BMD, GSECARS sector of the Advanced Photon Source (APS), Argonne National Laboratory (ANL). The X-ray beam was focused to less than a 10- μm spot onto the sample to minimize peak broadening caused by pressure gradients in

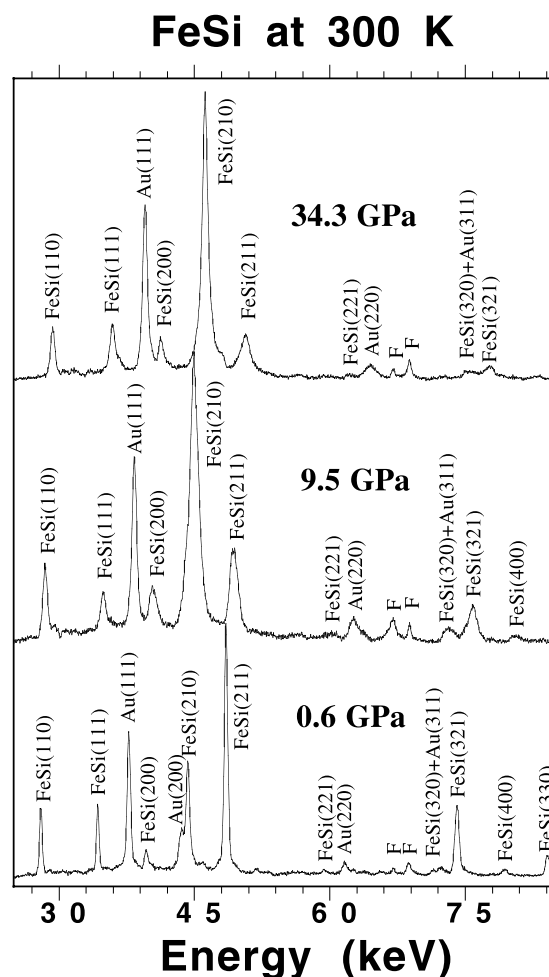


Figure 4. A series of X-ray diffraction patterns of ϵ -FeSi at increasing pressures to 34.3 GPa at 300 K. F, the fluorescence peaks of gold; Au, the internal pressure calibrant.

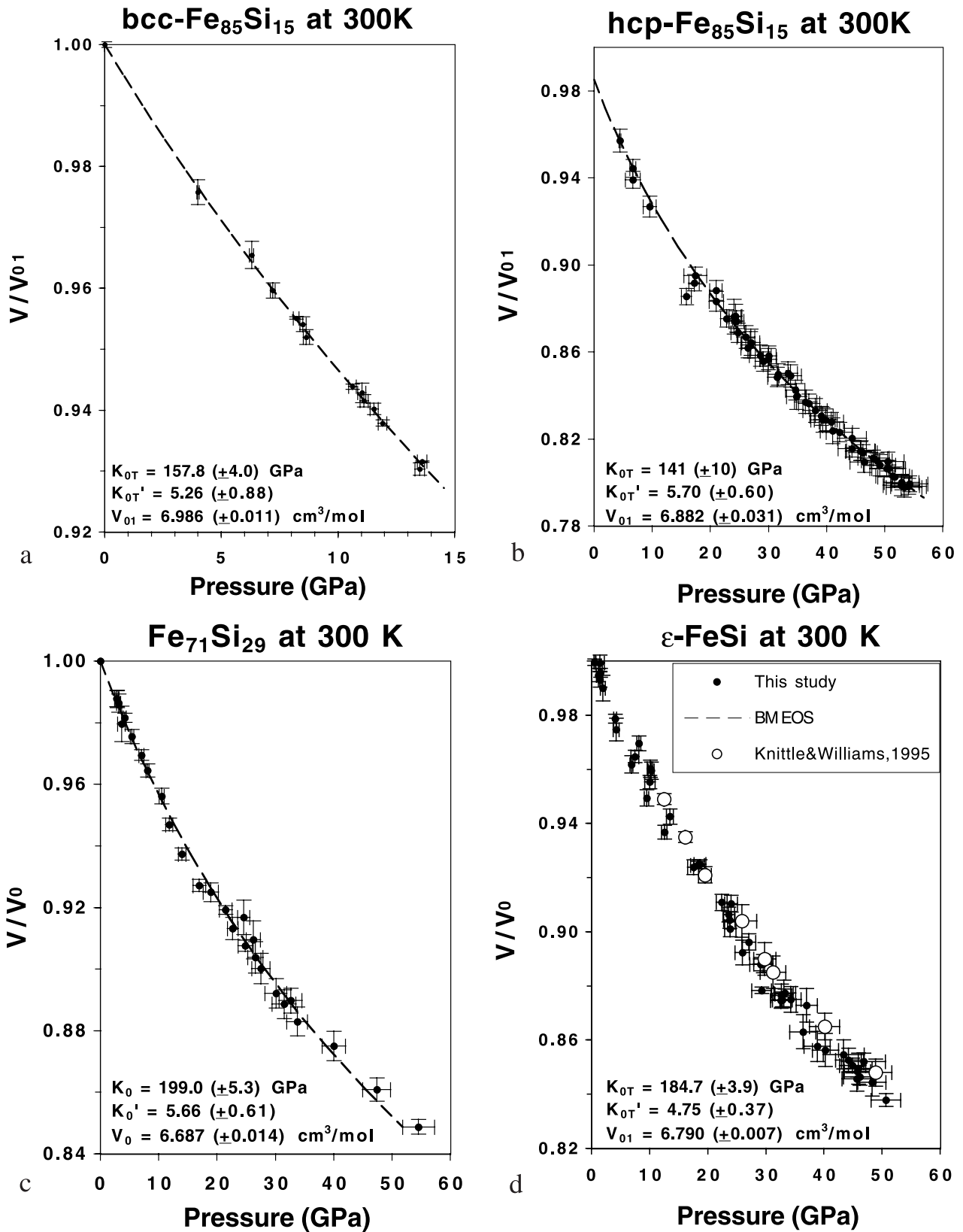


Figure 5. (a) Relative volume of bcc-Fe₈₅Si₁₅ as a function of pressure at 300 K. Dashed line represents a weighted least squares fit. (b) Relative volume of the hcp-Fe₈₅Si₁₅ as a function of pressure at 300 K. (c) Relative volume of the Fe₇₁Si₂₉ as a function of pressure at 300 K. (d) Relative volume of the ε-FeSi as a function of pressure at 300 K (solid circles). Compression data for ε-FeSi collected by *Knittle and Williams* [1995] are also plotted for comparison (open circles). A lower V_0 used by *Knittle and Williams* [1995] may cause the systematic discrepancy between the compression data.

the sample chamber. Two clean-up slits, 30 μm in width, were used to cut off the tails of the focused X-ray beam. The energy calibration of the detector was accomplished using known energies of X-ray emission lines of several elements (Ag, Cd, and Co). The diffracted X rays were collected by a Ge solid state detector at a fixed 2θ angle of about 8° or 12° . The exposure times ranged from 5 to 10 min. The d -spacings of the sample and the internal standard were determined by fitting a Gaussian curve to each diffraction peak. Pressures were calculated from three to five diffraction lines from the set of (111), (200), (220), (311), and (222) for the gold standard [Heinz and Jeanloz, 1984a]. Three to five diffraction peaks from the set of (100), (101), (102), (110), and (103) were used to calculate the unit cell parameters of hcp-Fe₈₅Si₁₅. Three to four diffraction peaks from the set of (100), (110), (200), and (211) were used to calculate the unit cell volume of Fe₇₁Si₂₉. Three to six diffraction lines from the set of (110), (111), (210), (211), (221), and (321) were used to calculate the unit cell parameters of ϵ -FeSi. Since all Fe-Si samples were in simple structures, the cell parameters of the samples were calculated from the d -spacings. Standard deviations of the cell parameters were calculated from the variation of the d -spacings.

[8] Since some of the X-ray diffraction peaks of Au overlapped with those of bcc-Fe₈₅Si₁₅, NaCl was mixed with the sample for use as a pressure calibrant in an externally heated DAC up to 14 GPa to measure the EOS of the bcc-Fe₈₅Si₁₅ in a better way [Birch, 1978; Bassett *et al.*, 1993]. The diffracted X rays were collected at 300 K after the sample was annealed at about 800 K at high pressures to reduce the effects of nonhydrostaticity. Three diffraction peaks, (110), (200), and (211), were used to calculate the unit cell parameters for bcc-Fe₈₅Si₁₅.

[9] The Fe₇₁Si₂₉ alloy (in ordered B2 structure) and the ϵ -FeSi alloy (in B20 structure) were laser heated in a DAC using the methods of Lin *et al.* [2002] to investigate the stability of these phases under high pressures and high temperatures. A sandwiched sample configuration, using NaCl as both the thermal insulator and pressure calibrant, was used in this study [Birch, 1978; Heinz and Jeanloz, 1984b].

3. Results

[10] Fe₈₅Si₁₅, Fe₇₁Si₂₉, and ϵ -FeSi were compressed up to 54.3, 54.5, and 50.7 GPa, respectively (Figures 1–4). A phase transformation in the Fe₈₅Si₁₅ alloy from bcc to hcp was observed to begin near 16 GPa and was completed by 36 GPa (Figures 1 and 2). The first appearance of the bcc phase during quench occurred at about 8 GPa, and the hcp phase can still be seen above 2 GPa (Figure 2). The two-phase (bcc + hcp) zone in Fe₈₅Si₁₅ extends over a wide pressure range of 20 GPa at 300 K, compared with 11 GPa for the same structural transition in iron [Huang *et al.*, 1987]. The volume change ($\Delta V/V_0$) across the bcc-hcp phase transition is about -1.9% or $-0.13 \text{ cm}^3/\text{mol}$.

[11] The pressure and volume data were analyzed with the Birch-Murnaghan equation of state (BM EOS) using a weighted least squares linear fit of the strain (f) and the normalized stress (F) to obtain values for K_{0T} and K'_{0T}

Table 1. Static Compression Data for bcc-Fe₈₅Si₁₅^a

Run	P , GPa	σ (P), GPa	a , \AA	σ (a), \AA	V/V_0	σ (V/V_0)
Fe8Sieh445	4.00	0.06	2.8288	0.0020	0.9758	0.0021
Fe8Sieh437	6.30	0.09	2.8188	0.0022	0.9655	0.0023
Fe8Sieh310	7.20	0.11	2.8131	0.0012	0.9596	0.0012
Fe8Sieh311	8.20	0.12	2.8087	0.0002	0.9551	0.0002
Fe8Sieh303	8.50	0.13	2.8077	0.0012	0.9541	0.0012
Fe8Sieh429	8.64	0.13	2.8056	0.0012	0.9520	0.0012
Fe8Sieh428	10.62	0.16	2.7976	0.0004	0.9439	0.0004
Fe8Sieh420	11.02	0.17	2.7966	0.0016	0.9428	0.0016
Fe8Sieh401	11.12	0.17	2.7953	0.0010	0.9415	0.0010
Fe8Sieh412	11.55	0.17	2.7940	0.0009	0.9402	0.0009
Fe8Sieh301	11.91	0.18	2.7917	0.0005	0.9379	0.0005
Fe8Sieh406	13.50	0.20	2.7842	0.0010	0.9304	0.0010
Fe8Sieh407	13.60	0.20	2.7854	0.0001	0.9316	0.0001

^a $a_0 = 2.8520(\pm 0.0015) \text{ \AA}$ and $V_0 = 6.985(\pm 0.011) \text{ cm}^3/\text{mol}$ at ambient conditions. The EOS of NaCl in B1 structure was used to determine the pressure [Birch, 1978].

[Birch, 1978; Jeanloz, 1981] (Figure 5 and Tables 1–4). The isothermal EOS parameters are given in Table 5.

[12] The Fe₇₁Si₂₉ and ϵ -FeSi alloys were laser heated up to ~ 1907 K at ~ 49 GPa and to ~ 2219 K at ~ 54 GPa, respectively. No phase transformation was observed in either sample, indicating that these silicon-rich Fe-Si alloys are relatively stable under high pressures and high temperatures.

4. Discussion

4.1. Isothermal Compression at 300 K

[13] Different results were reported in recent studies on the compressibility of ϵ -FeSi [Ross, 1994; Knittle and Williams, 1995; Sarrao *et al.*, 1994; Wood *et al.*, 1995; Guyot *et al.*, 1997; Vocadlo *et al.*, 1999] (Table 6). Most of the studies show that ϵ -FeSi has a bulk modulus between 160 and 180 GPa. Our experimental technique and pressure range are similar to that used by Knittle and Williams [1995], but a significantly higher bulk modulus and lower pressure derivative of the bulk modulus is observed in their study. As shown in Figure 5d, their compression data are systematically higher than that of this study, indicating that the lower, incorrect lattice parameter (see Table 6) of $V_0 = 6.701(\pm 0.032) \text{ cm}^3/\text{mol}$ for ϵ -FeSi under ambient conditions used by Knittle and Williams [1995] significantly affects their results. Analyzing the static compression data of Knittle and Williams [1995] using $V_0 = 6.790(\pm 0.007) \text{ cm}^3/\text{mol}$ from this study, K_{0T} is found to be $155.6(\pm 4.6)$ GPa with $K'_{0T} = 6.31(\pm 0.78)$.

[14] A plot of the molar volume of the Fe-Si alloys versus pressure (Figure 6) reveals that the compression curves of hcp-Fe, hcp-Fe₈₅Si₁₅, and ϵ -FeSi are almost identical to each other when extrapolated over 100 GPa, indicating that substitution of silicon into iron would lower the density of iron, but has little effect on the volume and compressibility of iron alloys at high pressures. This result is consistent with the conclusion of Zhang and Guyot [1999] that substitution of Si in iron would not significantly change the compressibility of the iron-rich Fe-Si alloy.

4.2. Thermal EOS Calculation

[15] To estimate the simultaneous effects of pressure and temperature on the EOS of iron-rich Fe-Si alloys, we have

Table 2. Static Compression Data for hcp-Fe₈₅Si₁₅^a

Run	P , GPa	σ (P), GPa	a , Å	σ (a), Å	c , Å	σ (c), Å	V/V_{01}	σ (V/V_{01})	c/a
Fe79Si636 ^b	4.5	0.2	2.5040	0.0060	4.0890	0.0110	0.9571	0.0053	1.6330
Fe79Si635 ^b	6.7	0.5	2.4980	0.0050	4.0540	0.0070	0.9444	0.0041	1.6229
Fe79Si547 ^b	6.7	1.3	2.4923	0.0050	4.0502	0.0050	0.9392	0.0039	1.6251
Fe79Si634 ^b	9.6	1.1	2.4890	0.0060	4.0080	0.0080	0.9269	0.0048	1.6103
Fe79Si608	15.9 ^c	0.8	2.4475	0.0050	3.9600	0.0050	0.8855	0.0038	1.6180
Fe79Si609	17.2	0.9	2.4535	0.0050	3.9691	0.0050	0.8918	0.0038	1.6177
Fe79Si633 ^b	17.4	1.9	2.4570	0.0050	3.9720	0.0060	0.8951	0.0039	1.6166
Fe79Si545 ^b	21.0	1.1	2.4406	0.0057	3.9947	0.0088	0.8883	0.0046	1.6368
Fe79Si610	21.0	1.2	2.4410	0.0060	3.9710	0.0060	0.8833	0.0045	1.6268
Fe79Si611	22.8	1.0	2.4380	0.0050	3.9450	0.0070	0.8754	0.0039	1.6181
Fe79Si612	24.0	1.2	2.4320	0.0100	3.9690	0.0160	0.8763	0.0080	1.6320
Fe79Si518	24.3	1.1	2.4345	0.0072	3.9612	0.0100	0.8764	0.0056	1.6271
Fe79Si519	24.5	1.3	2.4316	0.0063	3.9599	0.0080	0.8741	0.0049	1.6285
Fe79Si544 ^b	24.7	1.2	2.4350	0.0059	3.9250	0.0068	0.8688	0.0045	1.6119
Fe79Si520	26.1	1.5	2.4262	0.0067	3.9447	0.0088	0.8668	0.0052	1.6259
Fe79Si613	26.5	1.2	2.4280	0.0060	3.9160	0.0070	0.8618	0.0045	1.6129
Fe79Si522	26.9	1.6	2.4260	0.0077	3.9292	0.0106	0.8633	0.0060	1.6196
Fe79Si521	27.0	1.5	2.4290	0.0078	3.9247	0.0108	0.8644	0.0060	1.6158
Fe79Si523	28.6	1.9	2.4208	0.0065	3.9236	0.0081	0.8584	0.0049	1.6208
Fe79Si614	29.1	1.4	2.4220	0.0060	3.9070	0.0070	0.8556	0.0045	1.6131
Fe79Si524	30.0	1.4	2.4201	0.0065	3.9182	0.0082	0.8567	0.0049	1.6190
Fe79Si632 ^b	30.0	1.4	2.4270	0.0060	3.9030	0.0070	0.8582	0.0045	1.6082
Fe79Si615	31.5	1.6	2.4150	0.0060	3.8970	0.0070	0.8485	0.0045	1.6137
Fe79Si525	31.7	1.8	2.4135	0.0065	3.9076	0.0082	0.8497	0.0049	1.6191
Fe79Si543 ^b	33.4	1.7	2.4197	0.0069	3.8899	0.0088	0.8502	0.0052	1.6076
Fe79Si526	33.8	1.9	2.4147	0.0069	3.9006	0.0089	0.8490	0.0052	1.6154
Fe79Si616	34.6	1.7	2.4090	0.0060	3.8890	0.0080	0.8425	0.0045	1.6144
Fe79Si527	34.8	1.7	2.4009	0.0080	3.9013	0.0085	0.8395	0.0059	1.6249
Fe79Si528	36.4	1.9	2.3989	0.0073	3.8965	0.0077	0.8371	0.0054	1.6243
Fe79Si617	37.0	1.8	2.4000	0.0060	3.8890	0.0070	0.8362	0.0044	1.6204
Fe79Si529	38.1	2.1	2.3952	0.0075	3.8898	0.0079	0.8331	0.0055	1.6240
Fe79Si618	39.1	2.0	2.3950	0.0060	3.8790	0.0070	0.8306	0.0044	1.6196
Fe79Si530	39.4	2.1	2.3891	0.0078	3.8893	0.0082	0.8287	0.0057	1.6279
Fe79Si542 ^b	39.9	2.0	2.3912	0.0072	3.8825	0.0075	0.8287	0.0052	1.6237
Fe79Si619	40.8	2.1	2.3930	0.0060	3.8730	0.0070	0.8279	0.0044	1.6185
Fe79Si531	41.0	2.1	2.3850	0.0081	3.8792	0.0086	0.8237	0.0059	1.6265
Fe79Si620	42.2	2.3	2.3880	0.0060	3.8670	0.0070	0.8232	0.0044	1.6193
Fe79Si631 ^b	44.4	2.5	2.3880	0.0060	3.8540	0.0070	0.8204	0.0044	1.6139
Fe79Si621	44.4	2.5	2.3840	0.0050	3.8440	0.0080	0.8156	0.0038	1.6124
Fe79Si622	45.9	2.6	2.3810	0.0060	3.8480	0.0070	0.8144	0.0044	1.6161
Fe79Si533	45.9	2.4	2.3733	0.0064	3.8745	0.0065	0.8147	0.0046	1.6325
Fe79Si534	46.3	2.5	2.3711	0.0052	3.8768	0.0052	0.8137	0.0037	1.6350
Fe79Si541 ^b	46.6	2.5	2.3742	0.0061	3.8466	0.0102	0.8094	0.0047	1.6202
Fe79Si623	48.1	2.7	2.3790	0.0050	3.8400	0.0060	0.8113	0.0036	1.6141
Fe79Si624	48.6	2.7	2.3780	0.0050	3.8370	0.0060	0.8100	0.0036	1.6135
Fe79Si535	49.1	2.8	2.3715	0.0067	3.8486	0.0086	0.8080	0.0049	1.6229
Fe79Si625	50.3	2.8	2.3740	0.0050	3.8330	0.0060	0.8064	0.0036	1.6146
Fe79Si536	50.5	2.8	2.3705	0.0060	3.8596	0.0073	0.8096	0.0044	1.6282
Fe79Si630 ^b	50.7	2.9	2.3740	0.0050	3.8330	0.0060	0.8064	0.0036	1.6146
Fe79Si626	51.4	2.9	2.3710	0.0050	3.8270	0.0060	0.8031	0.0036	1.6141
Fe79Si537	51.8	2.9	2.3701	0.0061	3.8268	0.0102	0.8025	0.0047	1.6146
Fe79Si539 ^b	52.9	3.1	2.3683	0.0065	3.8145	0.0080	0.7987	0.0047	1.6106
Fe79Si627	52.9	3.0	2.3690	0.0050	3.8190	0.0050	0.8001	0.0035	1.6121
Fe79Si538	53.3	3.1	2.3674	0.0063	3.8132	0.0076	0.7978	0.0045	1.6107
Fe79Si628	54.1	3.1	2.3680	0.0050	3.8130	0.0060	0.7982	0.0036	1.6102
Fe79Si629	54.3	3.1	2.3680	0.0050	3.8190	0.0060	0.7994	0.0036	1.6128

^a $V_{02}/V_{01} = 0.9851 (\pm 0.0032)$. Au was used as the pressure calibrant [Heinz and Jeanloz, 1984a].

^bCollected during the process of decreasing pressure.

^cThis file (Fe79Si608 at 15.9 GPa) was not included in the BM EOS fit because the X-ray peaks of hcp-Fe₈₅Si₁₅ were very weak.

calculated the thermal EOS of hcp-Fe and the hcp-Fe₈₅Si₁₅ alloy using the Mie-Grüneisen EOS. Since bcc Fe-Si alloys have thermodynamic properties similar to those of iron under high pressures and high temperatures [Zhang and Guyot, 1999], it is reasonable to assume that hcp-Fe₈₅Si₁₅ also has thermodynamic parameters similar to those of hcp-Fe. The parameters used in calculating the EOS of hcp-Fe and hcp-Fe₈₅Si₁₅ in this study are listed in Table 7 [Mata-

ssov, 1977; Jeanloz, 1979; Jephcoat et al., 1986; Mao et al., 1990; Uchida et al., 2001]. The Grüneisen parameter (γ_0) and $\partial \ln \gamma / \partial \ln V$ (q) of hcp-Fe measured by shock-wave experiments are used in this study [Jeanloz, 1979]. The reference isotherm of hcp-Fe [Jephcoat et al., 1986; Mao et al., 1990; Uchida et al., 2001] and hcp-Fe₈₅Si₁₅ at 300 K is calculated using the BM EOS [Birch, 1978]. The thermal EOS of Fe and Fe₈₅Si₁₅ is then calculated from the

Table 3. Static Compression Data for Fe₇₁Si₂₉^a

Run	P , GPa	σ (P), GPa	a , Å	σ (a), Å	V/V_0	σ (V/V_0)
Fe17Si401	2.8	0.3	2.7994	0.0025	0.9879	0.0026
Fe17Si402	2.9	0.3	2.7992	0.0026	0.9877	0.0028
Fe17Si403	3.1	0.4	2.7979	0.0028	0.9863	0.0030
Fe17Si502	3.6	0.7	2.7916	0.0055	0.9796	0.0058
Fe17Si404	4.2	0.3	2.7934	0.0014	0.9815	0.0015
Fe17Si405	5.4	0.3	2.7876	0.0022	0.9754	0.0023
Fe17Si406	7.0	0.4	2.7819	0.0017	0.9695	0.0018
Fe17Si407	8.1	0.3	2.7771	0.0020	0.9645	0.0021
Fe17Si408	10.5	0.4	2.7691	0.0025	0.9562	0.0026
Fe17Si409	11.8	0.6	2.7602	0.0020	0.9470	0.0021
Fe17Si410	14.0	0.6	2.7508	0.0019	0.9373	0.0019
Fe17Si412	17.0	1.1	2.7408	0.0020	0.9271	0.0020
Fe17Si413	18.9	1.3	2.7387	0.0030	0.9250	0.0030
Fe17Si414	21.5	1.1	2.7330	0.0014	0.9192	0.0014
Fe17Si415	22.7	0.9	2.7270	0.0036	0.9132	0.0036
Fe17Si507	24.6	1.0	2.7306	0.0055	0.9168	0.0055
Fe17Si416	24.9	1.2	2.7214	0.0018	0.9076	0.0018
Fe17Si423 ^b	26.3	1.2	2.7233	0.0062	0.9095	0.0062
Fe17Si417	26.5	1.3	2.7176	0.0051	0.9038	0.0051
Fe17Si418	27.5	1.5	2.7140	0.0049	0.9002	0.0049
Fe17Si419	30.1	1.9	2.7058	0.0048	0.8921	0.0047
Fe17Si420	31.5	2.1	2.7023	0.0047	0.8886	0.0046
Fe17Si422 ^b	32.6	1.9	2.7035	0.0041	0.8898	0.0040
Fe17Si421	33.8	1.8	2.6966	0.0048	0.8830	0.0047
Fe17Si510	40.0	2.0	2.6884	0.0049	0.8750	0.0048
Fe17Si511	47.4	2.4	2.6739	0.0039	0.8609	0.0038
Fe17Si512	54.5	2.7	2.6613	0.0025	0.8488	0.0024

^a $a_0 = 2.8108(\pm 0.0020)$ Å. Au was used as the pressure calibrant [Heinz and Jeanloz, 1984a].

^bCollected during the process of decreasing pressure.

reference isotherm by use of the Mie-Grüneisen equation [McQueen *et al.*, 1970]:

$$P_i(V) = P_{300}(V) + \frac{\gamma}{V} [E_i(T_i, \Theta) - E_{300}(300, \Theta)], \quad (1)$$

where $P_{300}(V)$ is the pressure at the reference isotherm (300 K), γ is the Grüneisen parameter, V is the volume, Θ is the Debye temperature, and $E_i(T_i, \Theta)$ and $E_{300}(300, \Theta)$ are the thermal energies at temperatures T_i and 300 K, respectively. The thermal energies are calculated using

$$E(T, \Theta) = \left(\frac{9nRT}{x^3} \int_0^x \frac{x^3}{e^x - 1} dx \right) + E_e \quad (2)$$

with

$$x = \frac{\Theta(V)}{T}, \quad (3)$$

$$\gamma = \gamma_0 \left(\frac{V}{V_0} \right)^q = - \frac{d \ln \Theta}{d \ln V}, \quad (4)$$

$$\Theta = \Theta_0 \exp \left\{ \frac{\gamma_0}{q} \left[1 - \left(\frac{V}{V_0} \right)^q \right] \right\}. \quad (5)$$

Here E_e is the electronic contribution of the energy [Boness and Brown, 1986; Anderson and Ahrens, 1994], and the Grüneisen parameter and the Debye temperature are considered to be only functions of volume [Matassov,

1977], i.e., the quasi-harmonic approximation was used. The electronic energy is

$$E_e = \int_0^T C_e dT, \quad (6)$$

where C_e is the electronic contribution to the heat capacity and T is the temperature. C_e data calculated by Boness and Brown [1986] are used both for Fe and Fe₈₅Si₁₅.

[16] The reference adiabatic EOS is generated by using the BM EOS with $K_{0S} = K_{0T}(1 + \alpha\gamma T)$ and $K'_{0S} \sim K'_{0T}$

Table 4. Static Compression Data for ϵ -FeSi^a

Run	P , GPa	σ (P), GPa	a , Å	σ (a), Å	V/V_0	σ (V/V_0)
FeSi201	0.6	0.1	4.4840	0.0018	0.9996	0.0012
FeSi101	1.2	0.3	4.4768	0.0038	0.9948	0.0025
FeSi134 ^b	1.5	0.7	4.4834	0.0046	0.9992	0.0031
FeSi102	1.5	0.3	4.4743	0.0042	0.9931	0.0028
FeSi219 ^b	1.6	0.3	4.4781	0.0024	0.9957	0.0016
FeSi202	1.9	0.4	4.4695	0.0071	0.9899	0.0047
FeSi103	4.1	0.2	4.4524	0.0025	0.9786	0.0016
FeSi203	4.2	0.5	4.4464	0.0064	0.9747	0.0042
FeSi204	6.9	0.2	4.4267	0.0049	0.9618	0.0032
FeSi104	7.5	0.7	4.4312	0.0035	0.9647	0.0023
FeSi218 ^b	8.1	0.4	4.4387	0.0042	0.9696	0.0028
FeSi205	9.5	0.4	4.4077	0.0046	0.9494	0.0030
FeSi105	10.0	0.9	4.4168	0.0038	0.9553	0.0025
FeSi133 ^b	10.1	0.5	4.4243	0.0048	0.9602	0.0031
FeSi132 ^b	10.3	0.5	4.4232	0.0049	0.9595	0.0032
FeSi206	12.5	0.5	4.3881	0.0040	0.9368	0.0026
FeSi106	13.5	0.6	4.3971	0.0044	0.9426	0.0028
FeSi207	17.5	1.0	4.3678	0.0043	0.9239	0.0027
FeSi108	18.3	1.2	4.3690	0.0021	0.9246	0.0013
FeSi107	18.7	1.0	4.3694	0.0026	0.9249	0.0017
FeSi208	22.4	0.9	4.3472	0.0048	0.9109	0.0030
FeSi109	23.6	1.0	4.3400	0.0043	0.9064	0.0027
FeSi209	23.7	1.0	4.3368	0.0049	0.9044	0.0031
FeSi110	23.9	1.0	4.3316	0.0046	0.9011	0.0029
FeSi215 ^b	24.0	1.1	4.3465	0.0049	0.9104	0.0031
FeSi111	25.9	1.3	4.3175	0.0073	0.8923	0.0045
FeSi210	27.1	1.2	4.3238	0.0048	0.8962	0.0030
FeSi113	29.1	1.8	4.3105	0.0040	0.8880	0.0025
FeSi130 ^b	29.3	1.7	4.2948	0.0019	0.8783	0.0012
FeSi214 ^b	29.6	1.7	4.3108	0.0054	0.8882	0.0033
FeSi112	32.6	1.7	4.2889	0.0049	0.8747	0.0030
FeSi211	32.7	1.8	4.2917	0.0070	0.8764	0.0043
FeSi114	32.9	1.7	4.2890	0.0040	0.8748	0.0024
FeSi213 ^b	33.1	1.6	4.2927	0.0083	0.8770	0.0051
FeSi115	33.3	1.8	4.2932	0.0024	0.8773	0.0015
FeSi212	34.3	1.8	4.2894	0.0077	0.8750	0.0047
FeSi116	36.4	2.2	4.2698	0.0104	0.8631	0.0063
FeSi118	37.0	1.9	4.2859	0.0102	0.8729	0.0062
FeSi129 ^b	38.8	2.3	4.2608	0.0092	0.8576	0.0056
FeSi117	40.2	2.3	4.2586	0.0102	0.8563	0.0062
FeSi119	43.4	3.5	4.2559	0.0090	0.8547	0.0054
FeSi122	44.2	2.2	4.2525	0.0072	0.8526	0.0043
FeSi121	44.9	2.3	4.2492	0.0045	0.8506	0.0027
FeSi124	45.7	2.5	4.2410	0.0076	0.8457	0.0045
FeSi127 ^b	45.8	2.3	4.2454	0.0032	0.8484	0.0019
FeSi128 ^b	45.9	2.4	4.2472	0.0096	0.8494	0.0058
FeSi123	46.2	2.8	4.2420	0.0049	0.8463	0.0029
FeSi120	46.8	3.8	4.2517	0.0051	0.8522	0.0031
FeSi126 ^b	48.3	2.4	4.2387	0.0084	0.8444	0.0050
FeSi125	50.7	2.5	4.2279	0.0040	0.8379	0.0024

^a $a_0 = 4.4846(\pm 0.0015)$ Å. Au was used as the pressure calibrant [Heinz and Jeanloz, 1984a].

^bCollected during the process of decreasing pressure.

Table 5. EOS Parameters of bcc-Fe₈₅Si₁₅, hcp-Fe₈₅Si₁₅, Fe₇₁Si₂₉, and ε-FeSi

Composition	V_0 , cm ³ /mol	Structure	K_{0T}	K'_{0T}	Highest Pressure, GPa
Fe ₈₅ Si ₁₅	6.986 (±0.007)	bcc	157.8 (±4.0)	5.26 (±0.88)	13.6
Fe ₈₅ Si ₁₅	6.882 (±0.031)	hcp	141 (±10)	5.70 (±0.60)	54.3
Fe ₇₁ Si ₂₉	6.687 (±0.014)	B2	199.0 (±5.3)	5.66 (±0.61)	54.5
ε-FeSi	6.790 (±0.007)	B20	184.7 (±3.9)	4.75 (±0.37)	50.7

(Table 7). High-temperature adiabats are then calculated from the reference adiabat by use of the Mie-Grüneisen equation, and the temperature along each adiabat is

$$T = T_0 \exp \left\{ \frac{\gamma_0}{q} \left[1 - \left(\frac{V}{V_0} \right)^q \right] \right\}. \quad (7)$$

[17] The Hugoniot curve is also calculated by use of the Mie-Grüneisen equation as follows [McQueen *et al.*, 1970]:

$$P_H = P_S + \frac{\gamma}{V} (E_H - E_S) \quad (8)$$

with

$$E_H = E_0 + \frac{(P_H + P_0)(V_0 - V)}{2}, \quad (9)$$

$$E_S = E_0 - \int_{V_0}^V P_S dV, \quad (10)$$

where P_H is the Hugoniot pressure, P_S is the pressure along the adiabat passing through $P = 0$ GPa and $T = 300$ K, E_H is the energy along the Hugoniot, E_S is the energy along the adiabat passing through $P = 0$ GPa and $T = 300$ K and E_0 , P_0 , and V_0 are energy, pressure, and volume at the initial state, respectively. Thus the Hugoniot pressure is

$$P_H = \frac{P_S + \frac{\gamma}{V} \int_{V_0}^V P_S dV}{1 - \frac{\gamma}{2} \left(\frac{V_0}{V} - 1 \right)}. \quad (11)$$

The calculated Hugoniot curve from our data is consistent with shock-wave experiments on Fe₈₇Si₁₃

(6.9 wt.% Si) [Marsh, 1980] (Figure 7). The calculated adiabat passing through $P = 0$ GPa and $T = 3000$ K for hcp-Fe₈₅Si₁₅ shows that Fe-rich Fe-Si alloys have a lower density than pure iron under core conditions (Figure 8). The effect of melting on the density was observed to be less than a few percent (within the experimental error of the shock-wave measurements) at pressures relevant to the core [Jeanloz, 1979]. Therefore our calculated EOS of solid hcp-Fe and hcp-Fe₈₅Si₁₅ can be useful in discussing the density deficit of the liquid outer core.

[18] Upon comparison with the PREM density profile (Figure 8), the calculated EOS of Fe₈₅Si₁₅ are consistent with an outer core containing Si as the light alloying element. An outer core containing about 8–10 wt.% Si (assuming negligible volume change upon melting) and an inner core containing about 4 wt.% Si in iron can satisfy the density deficits in the core (Figure 8). Assuming a maximum silicon content in the core of 7.3 wt.% based on geochemical mass balance [Allègre *et al.*, 1995], a small amount of other light element(s) in addition to silicon may be needed to satisfy the density deficit of the outer core. However, the exact amount of the light element in the core depends on the temperature profile of the core and the thermoelastic properties of the alloyed iron [Poirier, 1994; Hillgren *et al.*, 2000].

4.3. Bulk Sound Velocity Calculation

[19] The bulk sound velocity of the core appears to be 3% higher than that of pure liquid iron [Ahrens, 1982; Anderson and Ahrens, 1994]. The change of the bulk sound velocity across the liquid outer core to the solid inner core is less than 1% [Dziewonski and Anderson, 1981], suggesting that melting probably has small effect on the bulk sound velocity of iron under core conditions. High-pressure X-ray diffraction experiments on liquid Fe-Si alloy show that, in contrast to liquid Fe-S alloy [Sanloup *et al.*, 2000], the presence of silicon in iron has a negligible effect on the compressibility of iron [Sanloup *et al.*, 2001]. Therefore if silicon is the dominant light element in the core, Fe-Si alloys must have higher bulk sound velocity

Table 6. Comparison of the Elastic Parameters of ε-FeSi

Reference	Pressure Range, GPa	K_0 , GPa	K'_0	V_0 , cm ³ /mol	Technique
This study	51	185 (±4)	4.8 (±0.4)	6.790 (±0.007)	X-ray diffraction in DAC
Ross [1994]	7	176 (±3)	4 (fixed)	n.a.	single-crystal X-ray diffraction
Knittle and Williams [1995]	50	209 (±6)	3.5 (±0.4)	6.701 (±0.032)	X-ray diffraction in DAC
Sarrao <i>et al.</i> [1994]	room pressure	173 ^a	n.a.	n.a.	ultrasonic spectroscopy
Wood <i>et al.</i> [1995]	9	160 (±1)	4 (fixed)	6.794 (±0.01)	neutron diffraction
Guyot <i>et al.</i> [1997]	8	172 (±3)	4 (fixed)	6.805 (±0.005)	X-ray diffraction in multianvil apparatus
Vocadlo <i>et al.</i> [1999]	calculation	227	3.9	6.692	first-principle calculations

^aThe adiabatic bulk modulus.

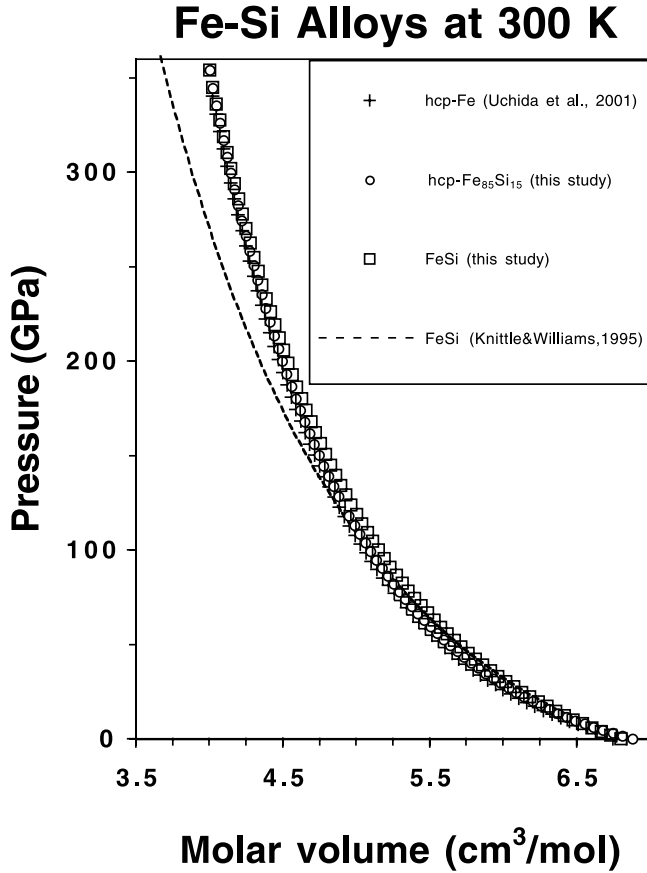


Figure 6. Comparison of molar volume (per unit molar atom) versus pressure at 300 K for hcp-Fe, hcp-Fe₈₅Si₁₅, and ϵ -FeSi. Crosses, hcp-Fe [*Jephcoat et al.*, 1986; *Mao et al.*, 1990; *Uchida et al.*, 2001]; open circles, hcp-Fe₈₅Si₁₅; open squares, ϵ -FeSi in B20 structure (extrapolated from this study); dashed line, ϵ -FeSi extrapolated from the data given by *Knittle and Williams* [1995].

than that of iron under core conditions. The bulk sound velocity is defined as

$$V_{\phi} = \sqrt{\frac{K_S}{\rho}}, \quad (12)$$

where ρ is the density and K_S is the adiabatic bulk modulus. The adiabatic bulk modulus of iron-rich Fe-Si alloys at

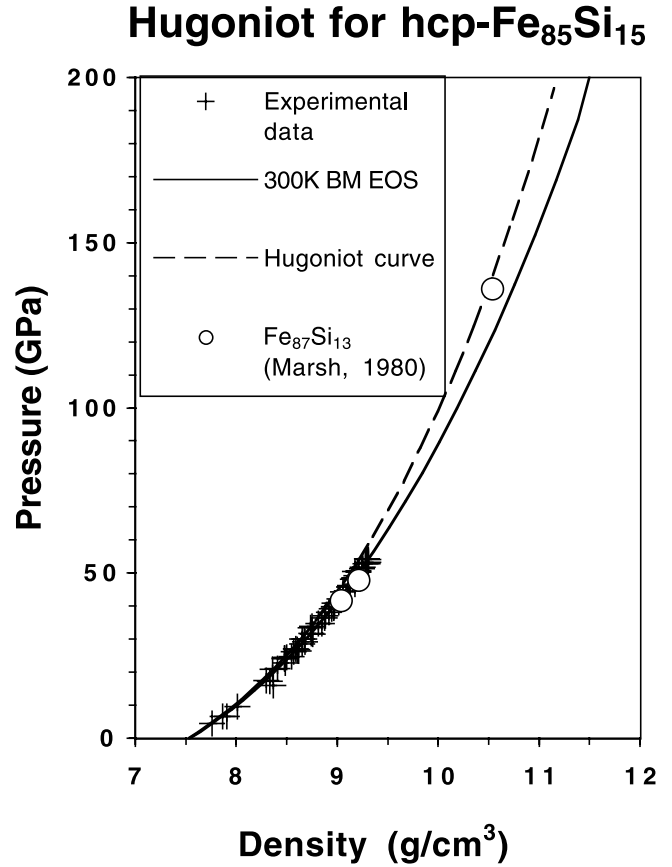


Figure 7. Hugoniot for hcp-Fe₈₅Si₁₅ alloy. Solid line, 300 K BM EOS of the hcp-Fe₈₅Si₁₅; dashed line, calculated Hugoniot; open circles, shock-wave data for Fe₈₇Si₁₃ [*Marsh*, 1980]; crosses, static compression data for hcp-Fe₈₅Si₁₅ (this study).

ambient conditions can be calculated from ultrasonic measurements of the elasticity of these alloys [*Guinan and Beshers*, 1968; *Alberts and Wedepohl*, 1971; *Routbort et al.*, 1971; *Machova and Kadeckova*, 1977]. As shown in Figure 9, the adiabatic bulk modulus of the Fe-Si alloys in the bcc structure is independent of the silicon content and the shear modulus increases slightly with increasing the silicon concentration in the iron-rich Fe-Si alloys. As expected from the decrease in density of Fe-Si alloys, the bulk sound velocity, compressional wave velocity, and shear wave velocity of the iron-rich bcc Fe-Si alloys increase with

Table 7. Thermodynamic Properties of hcp-Fe and hcp-Fe₈₅Si₁₅ at Ambient Conditions

Variable	hcp-Fe	Remark	hcp-Fe8wt.%Si	Remark
Molar volume (V_0)	6.84 (± 0.01) cm ³	<i>Uchida et al.</i> [2001]	6.88 (± 0.03) cm ³	this study
Isothermal bulk modulus (K_{0T})	135 (± 19) GPa	<i>Uchida et al.</i> [2001]	141 (± 10) GPa	this study
Pressure derivative of K_{0T} (K'_{0T})	6.0 (± 0.4)	<i>Uchida et al.</i> [2001]	5.70 (± 0.6)	this study
Adiabatic bulk modulus (K_{0S}) ^a	140 (± 20) GPa	calculated	146 (± 10) GPa	calculated
Heat capacity (C_v) ^b	23.1 (± 6.2) J/mol/K	calculated	24.3 (± 5.8) J/mol/K	calculated
Thermal expansion (α) ^c	5.5 (± 0.4) $\times 10^{-5}$ K ⁻¹	<i>Uchida et al.</i> [2001]	5.5 (± 0.4) $\times 10^{-5}$ K ⁻¹	assumed
Debye temperature (Θ)	464 K	<i>Matassov</i> [1977]	421 K	<i>Matassov</i> [1977]
Grüneisen parameter (γ_0) ^c	2.2 (± 0.5)	<i>Jeanloz</i> [1979]	2.2 (± 0.5)	assumed
$\partial \ln \gamma / \partial \ln V$ (q) ^c	1.62 (± 0.37)	<i>Jeanloz</i> [1979]	1.62 (± 0.37)	assumed

^a K_{0S} is calculated from $K_{0S} = K_{0T} (1 + \alpha \gamma T)$.

^b C_v is calculated from $C_v = \alpha K_{0T} V / \gamma$.

^cThermal expansion coefficient, Grüneisen parameter and $\partial \ln \gamma / \partial \ln V$ of hcp-Fe₈₅Si₁₅ are assumed to be the same as that of hcp-Fe.

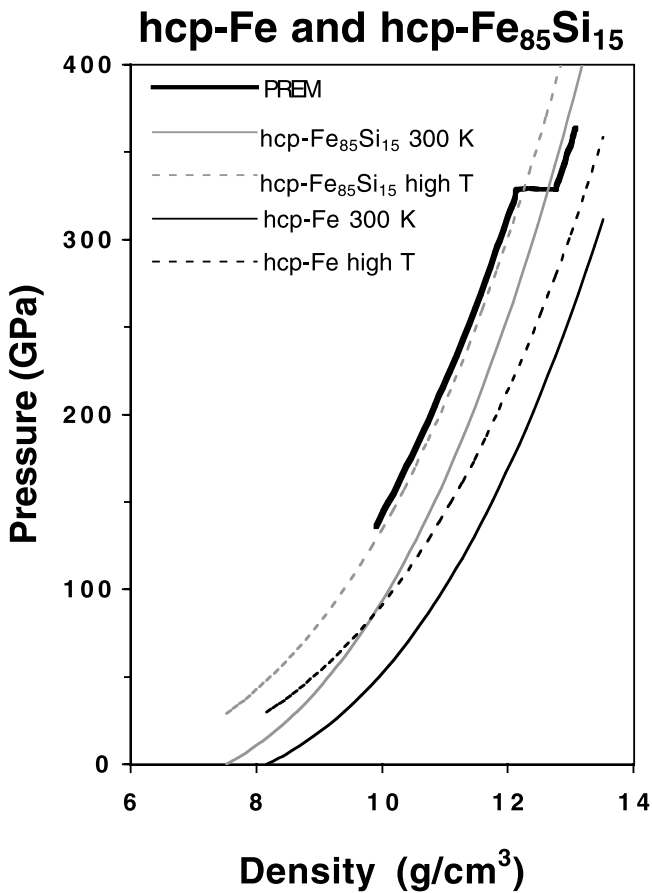


Figure 8. Calculated adiabats of hcp-Fe and hcp-Fe₈₅Si₁₅ compared to the PREM model. The temperature for the adiabat of the hcp-Fe₈₅Si₁₅ passing through $P = 0$ GPa and $T = 3000$ K at the core-mantle boundary is about 4948 K and is about 6300 K at inner core-outer core boundary. Thick black line, PREM; thin black line, an adiabat passing through $P = 0$ GPa and $T = 300$ K for hcp-Fe; dashed black line, an adiabat passing through $P = 0$ GPa and $T = 3000$ K for hcp-Fe; thin gray line, an adiabat passing through $P = 0$ GPa and $T = 300$ K for hcp-Fe₈₅Si₁₅; dashed gray line, an adiabat passing through $P = 0$ GPa and $T = 3000$ K for hcp-Fe₈₅Si₁₅.

increasing Si content (Figure 10). Addition of 8 wt.% Si increases the bulk sound velocity of bcc-Fe by approximately 3% (Figure 10). It was shown previously that substitution of silicon in hcp-Fe would lower the density, but would not significantly change the compressibility at high pressures; therefore addition of silicon into hcp-Fe would also increase the bulk sound velocity, consistent with silicon being a candidate light element in the core [Anderson and Ahrens, 1994].

[20] The calculated bulk sound velocity of hcp-Fe₈₅Si₁₅ is higher than that of hcp-Fe under core conditions. However, the extrapolation of the bulk modulus to core conditions is compromised by the error on K'_0 as well as the neglect of higher-order coefficients (terms beyond fourth order in strain) in the BM EOS. Precise experimental measurements on the elasticity and thermodynamic properties of Fe-Si alloys at higher pressure are needed in order to constrain the

bulk sound velocity of Fe-Si alloys under core conditions in a better way.

5. Conclusions

[21] Candidate light elements in the core are restricted by the following three constraints: lower density than pure iron, higher bulk sound velocity than iron, and freezing point depression of iron under core conditions [Poirier, 1994; Yang and Secco, 1999; Hillgren et al., 2000]. It has been demonstrated in this study that substitution of silicon in iron would lower the density, but would not significantly change the compressibility of iron at high pressures. Silicon, as the sole light alloying component in the core, can produce a sufficiently large reduction in the density. Addition of silicon into iron increases the bulk sound velocity of iron, consistent with silicon being a candidate light element in the core. Finally, it is expected from the phase diagram of Fe-Si system at ambient pressure [Kusbaschewski, 1982] that addition of silicon into iron lowers the melting point of iron under these conditions. The melting temperatures observed for Fe17wt.%Si up to 5.5 GPa are lower than those of pure iron [Yang and Secco, 1999].

[22] Despite the previous exclusion of Si as the primary alloying element in the outer core based on the relatively high K_0 and low K'_0 of ϵ -FeSi [Knittle and Williams, 1995; Williams and Knittle, 1997], our results based on the elasticity of three Fe-Si alloys (Fe₈₅Si₁₅, Fe₇₁Si₂₉, and ϵ -FeSi) indicate that silicon could be a significant light element

Elastic moduli of Fe-Si alloys at 300 K

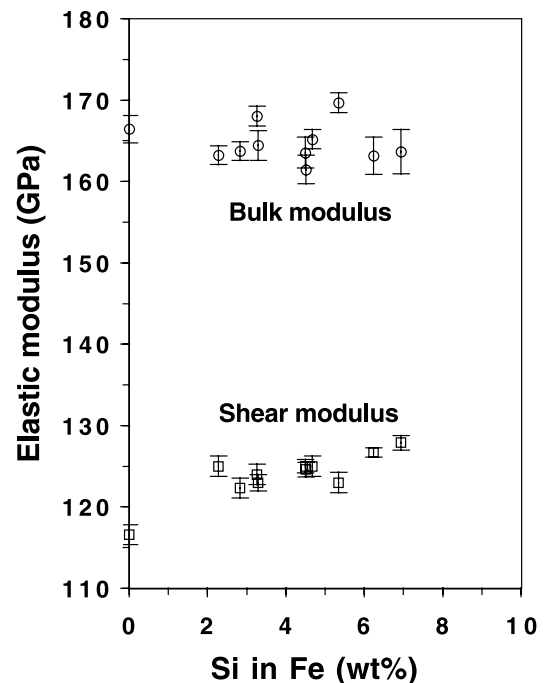


Figure 9. Adiabatic bulk moduli (open circles) and shear moduli (open squares) calculated from ultrasonic measurements of the elasticity of the iron-rich Fe-Si alloys at ambient conditions [Guinan and Beshers, 1968; Alberts and Wedepohl, 1971; Routbort et al., 1971; Machova and Kadeckova, 1977]. All alloys were in bcc structure.

V_ϕ , V_p , V_s of Fe-Si alloys at 300 K

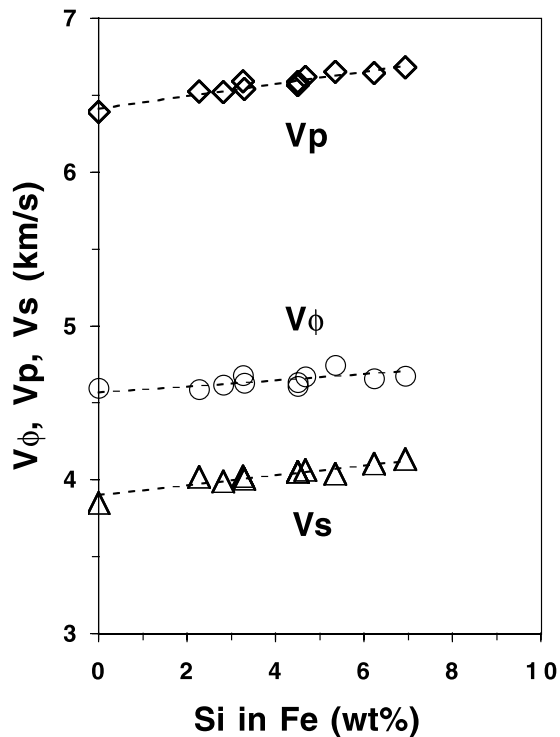


Figure 10. Bulk sound velocities (open circles), compressional wave velocities (open diamonds), and shear wave velocities (open triangles) of iron-rich Fe-Si alloys calculated from ultrasonic measurements at ambient conditions [Guinan and Beshers, 1968; Alberts and Wedepohl, 1971; Roubort et al., 1971; Machova and Kadeckova, 1977].

in the core. However, the possible presence of other light elements cannot be excluded on the basis of these results.

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