



Fermi

Gamma-ray Space Telescope

Searches for Cosmic-Ray Electron Anisotropies with the *Fermi*-Large Area Telescope

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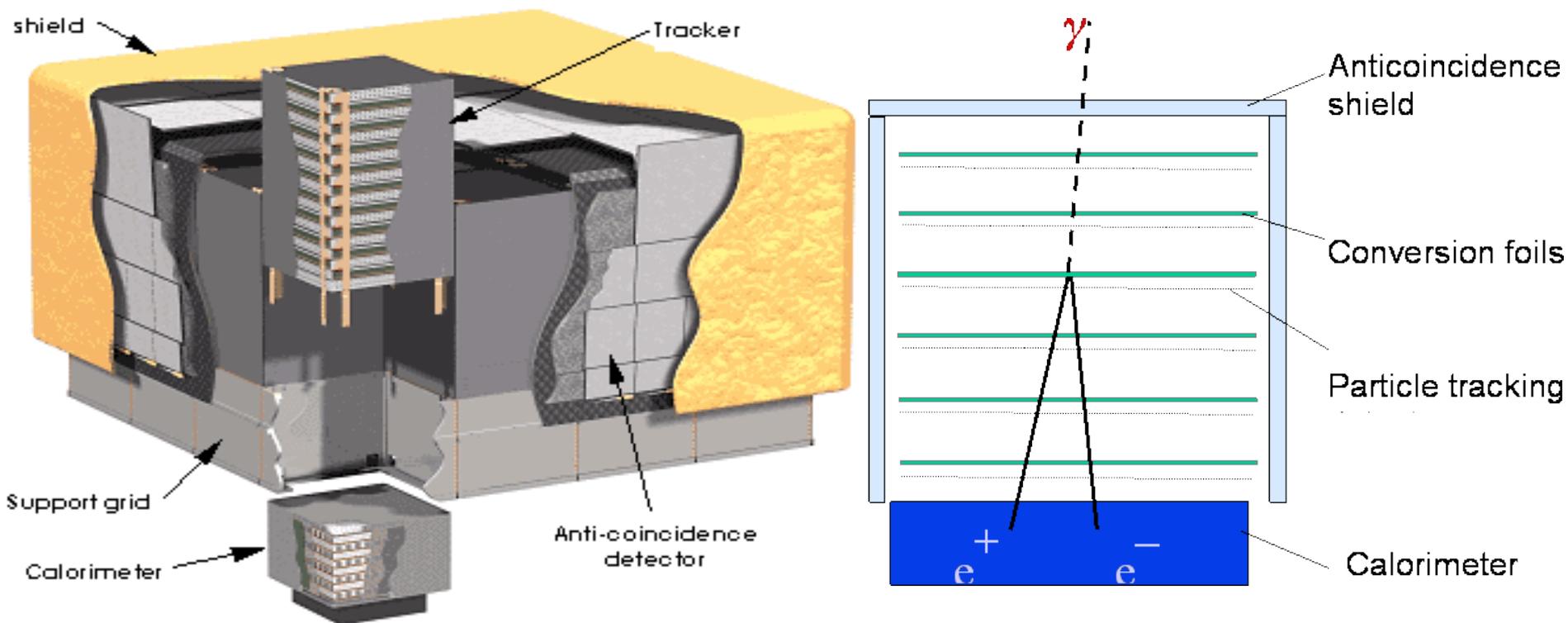
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on behalf of the *Fermi*-LAT
collaboration



- Pair-conversion gamma-ray detector
 - Tracker + Calorimeter + Anticoincidence Shield
- Primary conceived as a MeV/GeV gamma-ray detector but it can also be a detector for Cosmic Ray Electrons and Positrons (CREs)
- Wide field of view ($\sim 2.4\text{sr}$) + Full-sky coverage (every $\sim 3\text{hrs}$)



- The galactic magnetic field isotropizes the direction distribution of GeV-TeV CRs → information on the direction of CR sources is lost.
- Compared to hadronic Cosmic-Rays, CREs lose their energy rapidly.
 - 100 GeV (1TeV) CREs detected at the earth have originated **from relatively nearby locations** at most ~ 1.6 (0.75) kpc away.
 - Likely to have **originated from an anisotropic collection of few nearby sources**.
 - Depending on the propagation through the GMF, some anisotropy in the directions of GeV-TeV CREs might still exist.
- Remember: Studies tried to quantify the effect of nearby older pulsars to the detected CRE spectra and to the CR-Positron fraction (e.g. Geminga, Monogem).
 - They predict anisotropies towards the directions of dominant sites of CRE production.
 - **The discovery of an anisotropy in agreement with the predictions of these studies would be a big step towards revealing the sources of CREs.**



- The directions of CREs are affected by propagation through the magnetic fields of the galaxy, of the heliosphere, and of the earth.
 - The higher the CRE energy → the smaller the effect.
- **Geomagnetic field:**
 - CREs with energies over the geomagnetic cutoff suffer negligible deflections → ***We used a $E > 60 \text{ GeV}$ cut.***
- **Heliospheric Magnetic field:**
 - $E < 100 \text{ GeV}$ → appreciable/strong effects
 - $E > 1 \text{ TeV}$ → negligible effects
 - ◆ We are in between these two extremes.
 - ◆ HMF effects not well known → not obvious what minimum energy threshold to use → decide to go ahead with an $E > 60 \text{ GeV}$ dataset with the caveat that its low-energy part might have been considerably affected by the HMF (isotropized).

The Dataset



- First year of normal Science Operation (starting August 2008)
- ~1.3 million $E > 60 \text{ GeV}$ CRE events
- $\sim 0.1^\circ$ angular-reconstruction accuracy
- 10% energy resolution
- Low contamination with other species
 - Gammas $\rightarrow < 0.1\%$
 - Hadronic CRs $\rightarrow \sim 13\%$ (projected anisotropy under our sensitivity)
- All-sky uniform-exposure dataset \rightarrow allows us to search for anisotropies of ***any angular size (up to dipole) and from any direction in the sky.***

The Anisotropy Search



- We search for anisotropies with no a priori assumptions on the energy, direction, and angular size of the anisotropy.
 - Analyzed multiple subsets of the data: $E > 60\text{GeV}$, $E > 120\text{GeV}$, $E > 240\text{GeV}$, $E > 480\text{GeV}$.
 - Started with an all-sky search but also evaluated the sky towards candidate sources (different trials \rightarrow different sensitivity).
 - Searched for small ($\sim 10^\circ$) to large scale (90°) anisotropies.

The Anisotropy Search



- Search started with the construction of the **“no-anisotropy” skymap**:
 - the skymap expected to be detected on-average in the case of a perfectly isotropic CRE distribution → acts as our null hypothesis
- By comparing the no-anisotropy skymap to the actual skymap we could search for anisotropies in the actual data. Two methods were used:
 - Direct bin-to-bin comparison between the two maps.
 - Spherical Harmonic Analysis.
- As a cross-check, two techniques yielding consistent results were used to construct the “no-anisotropy” skymap:
 - Event-Shuffling & Direct-Integration Techniques.
 - Both techniques were based solely on the data (no MC simulations).

Event-Shuffling Technique



- Starts from the original data set and randomly shuffles the reconstructed directions of events (in the instrument frame).
 - The reconstructed energy and direction distributions (in the instrument frame) remain the same.
 - However, ***any anisotropy in sky coordinates is smeared out.***
- The randomization process is repeated multiple times (25), with each iteration producing a skymap that is statistically consistent with the case of an isotropic CRE direction distribution.
- These skymaps are then averaged, to construct the final no-anisotropy skymap.
- The technique is simple to implement and straight forward. It also has the benefit of automatically taking care of any short-term variations of the detector's effective area.

Direct-Integration Technique



- In general, the rate of events from some direction in instrument coordinates (θ, φ) is equal to the all-sky rate $R_{\text{allsky}}(t)$ times the probability that an event is reconstructed at that direction $P(\theta, \varphi, t)$.
- Based on the above, given a value for these two variables and the pointing information of the instrument, we can construct an associated skymap.
- What we want to do is to find the value of these variables that corresponds to the case of a perfectly isotropic CRE direction distribution, and using this value construct the no-anisotropy skymap. Which is this value?
 - ✗ As an anisotropy passes through the LAT's FOV, it creates fluctuations in the *instantaneous value* of these variables.
 - ✓ However, their *averaged-over-multiple-orbits value* remains constant, since any anisotropy events are averaged out.

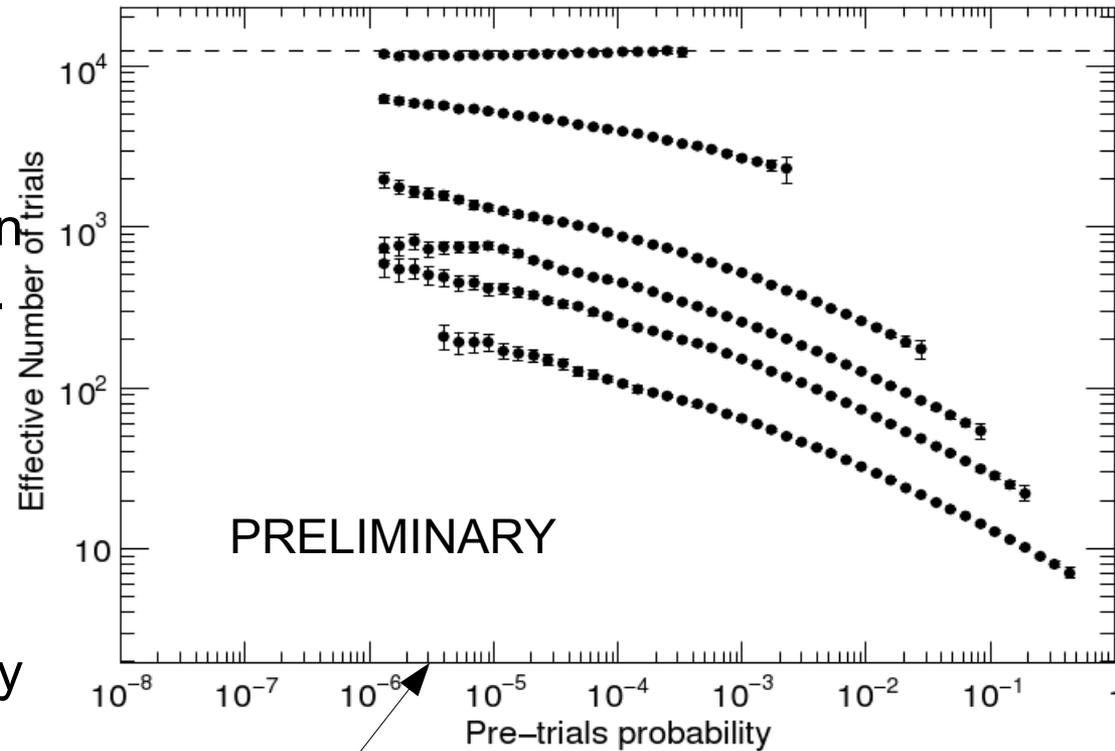


- By construction, started with $\sim 1^\circ$ *independent-bins* skymaps.
 - Can't search for tens-of-degs anisotropies with such skymaps (spillover effects reduce sensitivity).
- Produced several *integrated* skymaps, each with a different integration radius: 10° , 30° , 45° , 60° , 90° .
- Then, a bin-to-bin comparison between the actual and the predicted for the case of no anisotropy skymaps was performed:
 - For each pair of bins in the actual and the no-anisotropy skymap with contents $n_{\text{sig},i}$ and $n_{\text{iso},i}$ respectively, we calculated the probability of detecting a number of events at least as small as $n_{\text{sig},i}$ while expecting $n_{\text{iso},i}$.
 - For the Event-Shuffling technique we used Li & Ma significances. For the Direct-Integration technique we used simple Poisson probabilities.



- We calculated the effective number of trials involved in evaluating the 12,288 possible directions (number of bins) in an integrated significance skymap.
- These data were produced by simulating randomized significance skymaps and counting the fraction of such skymaps (P_{post}) that a probability less or equal than (P_{pre}) was found.

$$T_{\text{eff}} = \frac{\log(1 - P_{\text{post}})}{\log(1 - P_{\text{pre}})}$$



- Each curve shows the effective number of trials for a different integration radius.
- From top to bottom: independent-bins case, 10° , 30° , 45° , 60° , 90° integration radius.



- Effective number of trials + number of events in the dataset → sensitivity.
- The sensitivity becomes worse for smaller integration radii because the effective number of trials increases.
- The sensitivity becomes worse for higher energies because of the worse statistics.
- ***Sensitivity for a dipole anisotropy ranges from ~10% to a fraction of a pct.***

PRELIMINARY

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- Fractional excess needed for the bin-to-bin comparison method to detect an anisotropy with a 3σ post-trials significance.

A Significance Skymap



- A **pre-trials** significance map produced by a bin to bin comparison of the no-CRE-anisotropy to the actual skymap
- Integration radius 45° & Energy $>60\text{GeV}$
- Because of the large number of trials all the observed fluctuations are ***post-trials insignificant***.

PRELIMINARY

Results of the bin-to-bin comparison



- **Curves:** Correspondence between a pre and a post-trials significance. Left to right 90°, 60°, 40°, 30°, 10° integration radius.
- **Markers:** Best significance in each of the different tests.

PRELIMINARY

All the results were post-trials insignificant.



- An evaluation of the maps towards interesting locations (Geminga, Vela, Monogem pulsars, and the Galactic and anti-galactic centers) also yielded null results. Best post-trials significance $\sim 1.6\sigma$ from anti-galactic center.



- Spherical harmonic analysis of a fluctuation map produced by dividing the signal with the no-anisotropy skymap.
 1. The fluctuation map was first expanded in the basis of spherical harmonics, producing a set of a_{lm} coefficients.
 2. The average variance of these coefficients was used to construct an angular power spectrum.
 3. The power spectrum was compared to the expected spectrum of a white-noise data set (applicable for an isotropic sky).
 4. An increased power at a multipole l would indicate the presence of an anisotropy in the data would half-angular scale $\sim 90^\circ$.



- **Markers:** Power spectra from both techniques for $E > 60 \text{ GeV}$.
- **Curves:** Ranges that show the 3σ and 5σ integrated-probability fluctuations of a white-noise spectrum.
- All the spectra were consistent with being mere statistical fluctuations of an isotropic signal.

PRELIMINARY

Conclusion



- We searched for anisotropies in the reconstructed directions of ~ 1.3 million CREs detected by the *Fermi*-LAT during its first year of operation.
- Both methods used to perform the search resulted to null results.
- Upper limits on the degree of anisotropy were set (not presented).
- The analysis and results to be published soon.

THANK YOU!

Contamination



TABLE I. Geometry factor, residual contamination, number of counts before background subtraction, and the flux J_E multiplied by E^3 . The statistical error is followed by the systematic error. The latter does not include the effect due to the uncertainty in the absolute energy scale (see text).

Energy (GeV)	GF ($\text{m}^2 \text{sr}$)	Residual contamination	Counts	$E^3 J_E$ ($\text{GeV}^2 \text{s}^{-1} \text{m}^{-2} \text{sr}^{-1}$)
23.6–26.0	1.65	0.04	478 929	$151.6 \pm 1.2^{+7.3}_{-8.3}$
26.0–28.7	2.03	0.05	502 083	$152.6 \pm 0.9^{+6.2}_{-7.3}$
28.7–31.7	2.35	0.05	487 890	$151.4 \pm 0.8^{+5.1}_{-6.5}$
31.7–35.0	2.59	0.09	459 954	$151.3 \pm 1.8^{+5.2}_{-6.5}$
35.0–38.8	2.67	0.07	385 480	$149.6 \pm 0.7^{+4.4}_{-5.8}$
38.8–43.1	2.72	0.08	330 061	$150.2 \pm 0.7^{+4.5}_{-6.0}$
43.1–48.0	2.76	0.10	276 105	$148.6 \pm 0.7^{+4.9}_{-6.2}$
48.0–53.7	2.79	0.11	233 877	$146.5 \pm 0.7^{+4.9}_{-6.1}$
53.7–60.4	2.77	0.12	194 062	$145.5 \pm 0.7^{+5.0}_{-7.1}$
60.4–68.2	2.76	0.13	155 585	$143.2 \pm 0.7^{+5.6}_{-6.8}$
68.2–77.4	2.73	0.14	126 485	$141.9 \pm 0.8^{+5.6}_{-7.0}$
77.4–88.1	2.71	0.14	100 663	$140.8 \pm 0.8^{+6.2}_{-7.0}$
88.1–101	2.68	0.15	77 713	$139.0 \pm 0.9^{+6.4}_{-6.8}$
101–116	2.64	0.16	61 976	$139.0 \pm 0.9^{+6.4}_{-7.2}$
116–133	2.58	0.17	46 865	$139.4 \pm 1.0^{+6.9}_{-7.2}$
133–154	2.52	0.17	35 105	$139.5 \pm 1.2^{+7.2}_{-7.4}$
154–180	2.44	0.17	27 293	$140.8 \pm 1.3^{+6.9}_{-7.4}$
180–210	2.36	0.18	19 722	$142.3 \pm 1.5^{+7.1}_{-7.4}$
210–246	2.27	0.18	13 919	$140.9 \pm 1.7^{+7.4}_{-6.8}$
246–291	2.14	0.18	10 019	$140.9 \pm 1.9^{+7.5}_{-6.7}$
291–346	2.04	0.18	7 207	$140.4 \pm 2.2^{+6.7}_{-7.0}$
346–415	1.88	0.18	4 843	$139.4 \pm 2.6^{+7.0}_{-7.2}$
415–503	1.73	0.19	3 036	$134.0 \pm 3.1^{+9.3}_{-7.5}$
503–615	1.54	0.20	1 839	$127.4 \pm 4.1^{+8.7}_{-8.6}$
615–772	1.26	0.21	1 039	$115.8 \pm 4.8^{+15.2}_{-10.9}$
772–1000	0.88	0.21	544	$114.4 \pm 6.5^{+19.1}_{-17.8}$

Energy Resolution

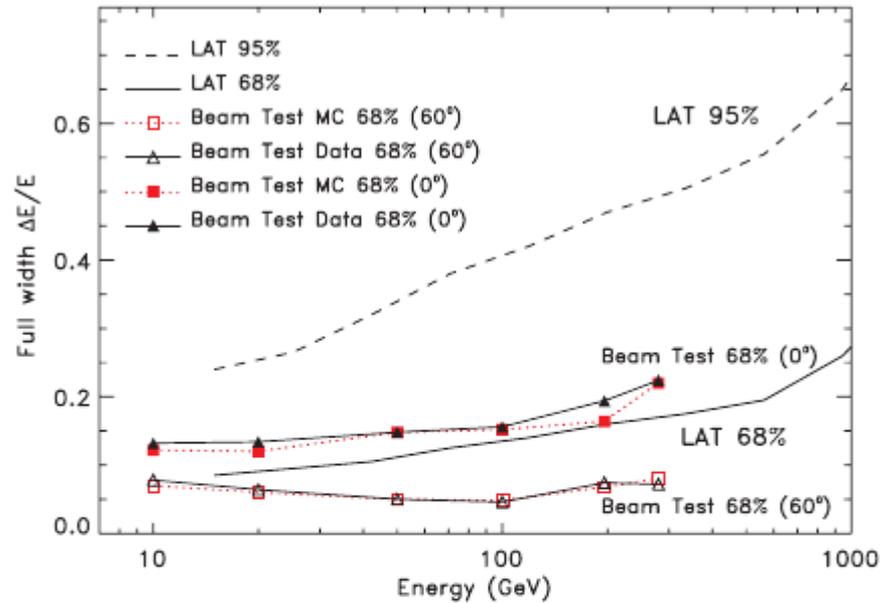


FIG. 1 (color online). Energy resolution for the LAT after electron selection; the full widths of the smallest energy window containing the 68% and the 95% of the energy dispersion distribution are shown. The comparison with beam test data up to 282 GeV and for on-axis and at 60° incidence shown in the figure indicates good agreement with the resolution estimated from the simulation.