

Reflector Antenna Discrete Distortion Determination: An Iterative-Field-Matrix Solution



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Abstract

The idealized shapes of satellite reflector antennas are often distorted once they are placed in orbit. The performance of such antennas can be improved by identifying the locations and amount of their surface distortions and then by correcting them using active surface distortion or array feeding. This work presents a method to determine the required discrete surface distortions to correct errors. The algorithm starts by discretizing the entire reflector surface into triangular patches, then by finding a linear relationship between the local distortion and the difference of the distorted and undistorted farfields patterns. A linear system of equations with the distortions as unknowns results when the number of observation points is the same as the number of triangular patches a unique solution is achieved without iteration. The conditioning of the system as a function of the observation direction and the spatial resolution is discussed. The method has been applied to determine thermal, gravitational and random distortions on a reflector antenna.

Introduction

Satellite reflector antennas are often distorted from their idealized shapes once they are placed in orbit. This distortion may be due to thermal gradients, mechanical stresses, damage from launch, or collisions with foreign objects. Variations from the designed shape lead to antenna performance degradation, compromising the antenna radiation pattern. The shape distortions can often be dynamically corrected with on-board actuators or feeding arrays, but it is essential to determine exactly where and how great the distortions are.

The proposed method approximates the nonlinear association between the position of a patch of reflector and the contribution to the field radiated by this patch in the farfield. By assuming a small positional perturbation, the association is linearized, and thus the perturbed position can be inverted given the perturbed pattern. Repeated application of this small perturbation inversion leads to an iterative procedure which quickly approaches the final patch position.

Analysis Technique

The forward model is based on the radiation of electric currents in a free space [1]:

$$\mathbf{E}^s(\mathbf{r}) = -jkZ_0 \int_{\Omega} [g_1 \mathbf{J}(\mathbf{r}') - g_2 (\mathbf{J}(\mathbf{r}') \cdot \hat{\mathbf{R}}) \hat{\mathbf{R}}] G d\Omega \quad (1)$$

$$g_1 = 1 - 1/(kR)^2 - j/(kR), \quad g_2 = 1 - 3/(kR)^2 - j3/(kR)$$

$$\mathbf{R} = \mathbf{r} - \mathbf{r}', \quad R = |\mathbf{R}|, \quad \hat{\mathbf{R}} = \mathbf{R}/R$$

$$G = \frac{\exp(-jkR)}{4\pi R}, \quad k = \omega\sqrt{\mu_0\epsilon_0}, \quad Z_0 = \sqrt{\mu_0/\epsilon_0}$$

The Physical Optics (PO) approximation considers the induced currents to be proportional to the magnetic incident field to the PEC

$$\mathbf{J}(\mathbf{r}') \approx \mathbf{J}_{po}(\mathbf{r}') = 2 \hat{\mathbf{n}}' \times \mathbf{H}^{inc}(\mathbf{r}') \quad (2)$$

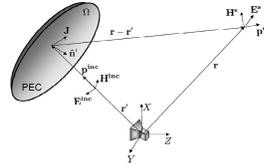
To solve (1), the surface is divided in triangular facets, the electric field contribution of the subdomain i in the observation point l can be expressed as

$$\mathbf{E}_{i,l}^s = \mathbf{E}^s(\mathbf{r}_l, \mathbf{r}'_i) = -jkZ_0 \frac{\exp(-jkR_{il})}{4\pi R_{il}} \int_{\Omega_i} [\mathbf{J}(\mathbf{r}'_i) - \hat{\mathbf{r}}_i (\mathbf{J}(\mathbf{r}'_i) \cdot \hat{\mathbf{r}}_i)] \exp(jk \mathbf{r}'_i \cdot \hat{\mathbf{r}}_i) d\Omega \quad (3)$$

$$\mathbf{r}_l = |\mathbf{r}_l|, \quad \hat{\mathbf{r}}_l = \mathbf{r}_l/r_l$$

The total scattered field is determined by the summation of the contributions of each subdomain

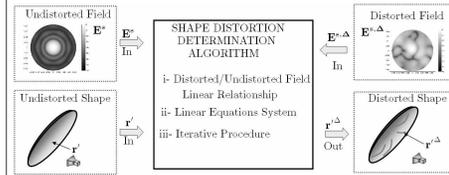
$$\mathbf{E}^s_l = \mathbf{E}^s(\mathbf{r}_l) = \sum_{i=1}^{N_{\Omega}} \mathbf{E}_{i,l}^s \quad (4)$$



Shape Distortion Determination Algorithm

The ideal undistorted shape, given by position vector \mathbf{r}' , produces the scattered field \mathbf{E}^s . Distortions to this ideal shape produce distorted scattered field $\mathbf{E}^{s,\Delta}$.

The problem: Determine the distorted surface, given by a position vector \mathbf{r}'_{Δ} , that produces the distorted scattered field $\mathbf{E}^{s,\Delta}$.



Distorted/Undistorted Fields Linear Relationship

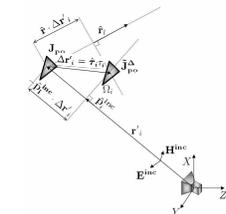
A local displacement produces new induced currents that can be approximated as

$$\tilde{\mathbf{J}}_{po}^{\Delta}(\mathbf{r}'_i + \Delta\mathbf{r}'_i) = \mathbf{J}_{po}(\mathbf{r}'_i) \exp(-jk \hat{\mathbf{p}}_i^{inc} \cdot \hat{\boldsymbol{\tau}}_i \tau_i) \quad (5)$$

The electric farfield produced by this new induced currents can be written as

$$\tilde{\mathbf{E}}_{i,l}^{s,\Delta} = \tilde{\mathbf{E}}^{s,\Delta}(\mathbf{r}_l, \mathbf{r}'_i + \Delta\mathbf{r}'_i) = \mathbf{E}_{i,l}^s \exp(j\phi_{i,l} \tau_i) \quad (6)$$

$$\phi_{i,l} = k (\hat{\mathbf{r}}_l - \hat{\mathbf{p}}_i^{inc}) \cdot \hat{\boldsymbol{\tau}}_i$$



The total scattered field can be written as the summation of the contribution of each facet

$$\tilde{\mathbf{E}}_l^{s,\Delta} = \tilde{\mathbf{E}}^{s,\Delta}(\mathbf{r}_l) = \sum_{i=1}^{N_{\Omega}} \tilde{\mathbf{E}}_{i,l}^{s,\Delta} \quad (7)$$

The relationship between the distorted and undistorted fields is linearized by approximating the exponential term in (6) by the first two terms of a series expansion

$$\tilde{\mathbf{E}}_{i,l}^{s,\Delta} \approx \mathbf{E}_{i,l}^s (1 + j\phi_{i,l} \tau_i) \quad (8)$$

The total scattered field can be written as the summation of the contribution of each facet

$$\tilde{\mathbf{E}}_l^{s,\Delta} \approx \sum_{i=1}^{N_{\Omega}} \mathbf{E}_{i,l}^s \Delta \quad (9)$$

Linear System of Equations

When the number of observation points is coincident with the number of distorted subdomains, a linear system of equations is derived

$$[A] [x] = [b] \quad (10)$$

$$A(l, i) = \tilde{\zeta} \cdot \mathbf{E}_{i,l}^s j\phi_{i,l} \quad (11)$$

$$x(i) = \tau_i \quad (12)$$

$$b(l) = \tilde{\zeta} \cdot (\mathbf{E}_l^{s,\Delta} - \mathbf{E}_l^s) \quad (13)$$

Iterative Procedure

The error introduced in the linearization of the problem imposes that an iterative scheme must be adopted for accurate results

$$\mathbf{r}'_i^{(m)} = \mathbf{r}'_i^{\Delta(m-1)} = \mathbf{r}'_i^{(m-1)} + \tau_i^{(m-1)} \boldsymbol{\tau}_i \quad (14)$$

$$\mathbf{r}'_i^{(0)} = \mathbf{r}'_i \quad (15)$$

$$\tau_i^{(0)} = 0 \quad (16)$$

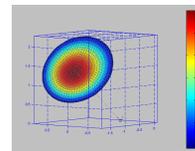
Application Example: Thermal and Random Errors Determination

The method is applied to determine Thermal/Gravitational and Random errors.

Antenna configuration

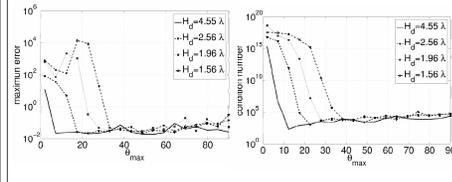
CHARACTERISTIC PARAMETERS OF THE SIMULATED EXAMPLE

Main Reflector	
Reflector diameter	$D = 1.68$ m
Focal length	$F = 1.832$ m
Offset height	$H = 1.43$ m
Feed	
Frequency	$f = 8.45$ GHz
Electric field polarization	$\psi = 45^\circ$
Offset angle	$\theta_0 = 43.13$
Subtended angle	$\theta_s = 63.11$
Taper edge height	$h_e = 12.08$
Taper edge radius	$r_e = 12.08$



Conditioning of the field matrix

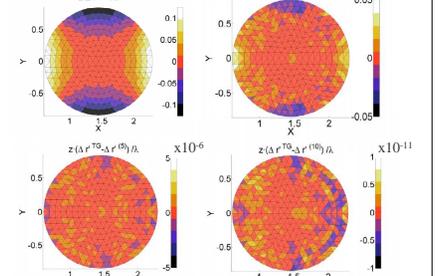
The condition number of the matrix depends on the spatial resolution, described in terms of the averaged size of the subdomains, as well as maximum observation angle existing between an observation point and the antenna boresight direction.



Shape variation

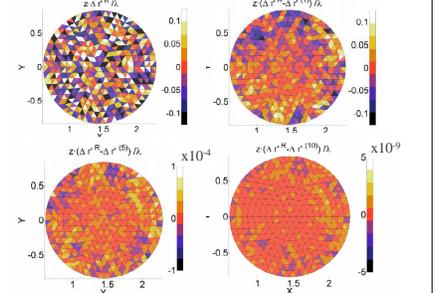
The Thermal/Gravitational errors can be expressed analytically by

$$\Delta\mathbf{r}_i^{TG} = \hat{\mathbf{z}} \delta_{max} \left(\frac{2\rho_i}{D} \right)^3 \cos(2\phi_i) z(\Delta r^{TG}, \Delta r^{(1)}) \hat{\mathbf{r}}_i \quad (17)$$

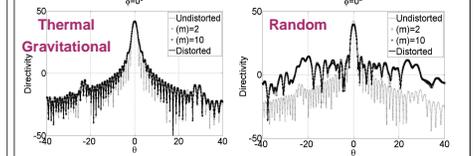


The Random errors are characterized by an uniform distribution

$$\Delta\mathbf{r}_i^R = \hat{\mathbf{z}} \delta_{max} X_i \quad X_i \sim U(-1, 1) \quad (18)$$



Pattern variation



Conclusions and technology transfer

A new method to determine the discrete distortions on reflector antennas have been presented. The method is based on the discretization of the reflector surface into triangular patches. By linearizing the difference between the distorted and undistorted far field patterns a linear system of equations is constructed. The application of the method for correcting thermal/gravitational and random surfaces yielding accurate results has been presented. The method can be applied to improve the radiation performance of on-board satellite reflector antennas by using active surfaces or phased arrays [2].

References

[1] J. A. Martínez-Lorenzo, A. G. Pino, I. Vega, M. Arias, and O. Rubio, "Exact induced-current analysis of reflector antennas," *IEEE Antennas and Propagation Magazine*, vol. 47, no. 2, pp. 92-100, 2005.
 [2] A. G. Pino, J. A. Martínez-Lorenzo, M. A. Arias, and C. Compositio, "Design and analysis of an adjustable subreflector for the hybrid mechanical-electronic pointing system at the satellite q/v band," in *Proc. IEEE Antennas and Propagation Society International Symposium (AP-S'06)*, Albuquerque, N.M., USA, July 2006, pp. 2451-2454.