

# IR, Raman, First Hyperpolarizability and Computational Study of 1-chloroethyl Benzene

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(Received: February 03, 2012; Accepted: March 19, 2012)

## ABSTRACT

The FTIR and FT-Raman spectra of 1-chloroethylbenzene were recorded and analyzed. The harmonic vibrational wavenumbers were calculated theoretically using Gaussian03 set of quantum chemistry codes. The calculated wavenumbers (B3LYP) agree well with the observed wavenumbers. The calculated first hyperpolarizability is reported and the title compound is an attractive object for further studies of nonlinear optics.

**Key words:** Chloroethyl, DFT, FT-IR, FT-Raman, hyperpolarizability.

## INTRODUCTION

Chloroethyl derivatives are new agents that have shown promising cytotoxic and antineoplastic activities. The cytotoxicity of many chloroethyl derivatives was demonstrated on several tumor cell lines<sup>1</sup>. Kulkarni *et al.*,<sup>2</sup> reported the quantitative structure activity relationship studies of chloroethyl derivatives. Nonlinear optics deals with the interaction of applied electromagnetic fields in various materials to generate new electromagnetic fields, altered in wavenumber, phase, or other physical properties<sup>3</sup>. Organic molecules able to manipulate photonic signals efficiently are of importance in technologies such as optical communication, optical computing, and dynamic image processing<sup>4,5</sup>. In this context, the dynamic first hyperpolarizability of the title compound is also calculated in the present study. The first hyperpolarizability ( $\beta_0$ ) of this novel molecular system is calculated using B3LYP/6-31G(d) method, based on the finite field approach. In the presence of an applied electric field, the energy of a system is a function of the electric field. First hyperpolarizability is a third rank tensor that can be described by a  $3 \times 3 \times 3$  matrix. The 27

components of the 3D matrix can be reduced to 10 components due to the Kleinman symmetry<sup>6</sup>. The components of  $\beta$  are defined as the coefficients in the Taylor series expansion of the energy in the external electric field. When the electric field is weak and homogeneous, this expansion becomes

$$E = E_0 - \sum_i \mu_i F^i - \frac{1}{2} \sum_{ij} \alpha_{ij} F^i F^j - \frac{1}{4} \sum_{ijk} \beta_{ijk} F^i F^j F^k - \frac{1}{24} \sum_{ijkl} \gamma_{ijkl} F^i F^j F^k F^l + \dots$$

where  $E_0$  is the energy of the unperturbed molecule,  $F^i$  is the field at the origin,  $\mu_i$ ,  $\alpha_{ij}$ ,  $\beta_{ijk}$  and  $\gamma_{ijkl}$  are the components of dipole moment, polarizability, the first hyperpolarizabilities, and second hyperpolarizabilities, respectively. In the present work, vibrational spectroscopic analysis of 1-chloroethyl benzene is reported.

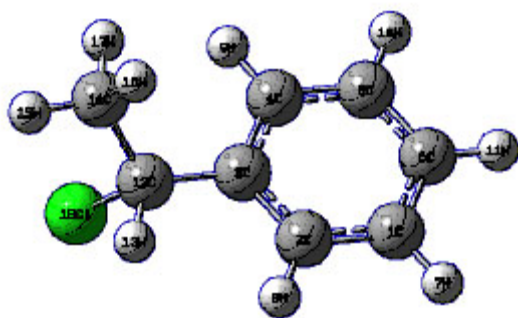
## EXPERIMENTAL

The FT-IR spectrum was recorded on a DR/Jasco FT-IR 6300 spectrometer and the FT-Raman spectrum was obtained on a BRUKER RFS 100/S, Germany. For excitation of the spectrum the emission of a Nd:YAG laser was used, excitation

wavelength 1064 nm, maximal power 150 mW. One thousand scans were accumulated with a total registration time of about 30 min. The spectral resolution after apodization was 2  $\text{cm}^{-1}$ .

### Computational details

Calculations of the title compound were carried out with Gaussian03 software program<sup>7</sup> using the B3LYP/6-31G\* basis sets to predict the molecular structure and vibrational wavenumbers. The DFT hybrid B3LYP functional tends also to overestimate the fundamental modes; therefore scaling factors have to be used for obtaining a considerably better agreement with experimental data. Therefore, a scaling factor of 0.9613 was uniformly applied to the DFT calculated wavenumbers<sup>8</sup>.



## RESULTS AND DISCUSSION

The observed IR, Raman and calculated (scaled) wavenumbers and the assignments are given in Table 1. The phenyl CH stretching modes occurs above 3000  $\text{cm}^{-1}$  and is typically exhibited as multiplicity of weak to moderate bands compared with the aliphatic CH stretching<sup>9</sup>. In the present case, the DFT calculations give  $\nu_{\text{CH}}$  modes of the phenyl ring at 3103, 3093, 3084, 3074 and 3066  $\text{cm}^{-1}$ . The bands observed at 3090, 3068  $\text{cm}^{-1}$  in the IR spectrum and at 3064  $\text{cm}^{-1}$  in the Raman spectrum were assigned as CH stretching modes of the phenyl ring. The benzene ring possesses six ring stretching vibrations, of which the four with the highest wavenumbers (occurring near 1600, 1580, 1490 and 1440  $\text{cm}^{-1}$ ) are good group vibrations<sup>10</sup>. With heavy substituents, the bands tend to shift to somewhat lower wavenumbers. In the absence of ring conjugation, the band at 1580  $\text{cm}^{-1}$  is usually

weaker than that at 1600  $\text{cm}^{-1}$ . In the case of C=O substitution, the band near 1490  $\text{cm}^{-1}$  can be very weak. The fifth ring stretching vibration is active near 1315  $\pm$  65  $\text{cm}^{-1}$ , a region that overlaps strongly with that of the CH in-plane deformation<sup>10</sup>. The sixth ring stretching vibration, or the ring breathing mode, appears as a weak band near 1000  $\text{cm}^{-1}$ , in mono, 1,3-di- and 1,3,5-trisubstituted benzenes. In the otherwise substituted benzenes, however, this vibration is substituent sensitive and difficult to distinguish from the ring in-plane deformation<sup>10,11</sup>. For mono substituted benzene, the  $\delta_{\text{Ph}}$  modes are expected in the range<sup>10</sup> of 1285-1610  $\text{cm}^{-1}$ . For the title compound, the phenyl ring stretching modes are assigned at 1494, 1455  $\text{cm}^{-1}$  in the IR spectrum, 1607  $\text{cm}^{-1}$  in the Raman spectrum and at 1598, 1580, 1496, 1457, 1330  $\text{cm}^{-1}$  theoretically. The ring breathing mode is observed at 1006  $\text{cm}^{-1}$  in the Raman spectrum which finds support from the computational value at 999  $\text{cm}^{-1}$ , as expected.

The in-plane CH vibrations are expected in the range of 1015-1300  $\text{cm}^{-1}$  for mono substituted benzenes, respectively<sup>10</sup>. For the title compound, the DFT calculations give 1299, 1192, 1177, 1057, 1025  $\text{cm}^{-1}$  as  $\nu_{\text{CH}}$  modes. These modes are observed at 1303, 1028  $\text{cm}^{-1}$  in the IR spectrum and at 1190, 1168, 1029  $\text{cm}^{-1}$  in the Raman spectrum. The out-of-plane CH deformations<sup>10</sup> are observed between 1000 and 700  $\text{cm}^{-1}$ . Generally, the CH out-of-plane deformations with the highest wavenumbers have a weaker intensity than those absorbing at lower wavenumbers. These  $\tilde{\alpha}_{\text{CH}}$  modes are observed at 971, 911  $\text{cm}^{-1}$  in the IR spectrum and at 972, 782  $\text{cm}^{-1}$  in the Raman spectrum. The DFT calculations give these modes at 994, 966, 918, 843 and 771  $\text{cm}^{-1}$ . In the present case, the out-of-plane CH deformation at 782  $\text{cm}^{-1}$  and the ring deformation at 697  $\text{cm}^{-1}$  in the IR spectrum form a pair of strong bands characteristics of mono substituted benzene<sup>12</sup>.

In aromatic compounds the asymmetric stretching vibrations of  $\text{CH}_3$  are expected in the range 2905-3000  $\text{cm}^{-1}$  and symmetric  $\text{CH}_3$  vibrations in the range<sup>10,13</sup> of 2860-2900  $\text{cm}^{-1}$ . The first of this results from the asymmetric stretching  $\nu_{\text{as}} \text{CH}_3$  mode in which two C-H bonds of the methyl group are extending while the third one is contracting. The second arises from the symmetrical stretching

**Table 1: Vibrational assignments of 1-chloroethyl benzene**

B3LYP/6-31G* $\nu$ (cm <sup>-1</sup> )	IR intensity	Raman activity	IR $\nu$ (cm <sup>-1</sup> )	Raman $\nu$ (cm <sup>-1</sup> )	Assignments
3103	15.22	282.06			$\nu$ CH
3093	35.37	20.65	3090		$\nu$ CH
3084	16.72	89.77			$\nu$ CH
3074	0.66	90.48			$\nu$ CH
3066	6.21	34.32	3068	3064	$\nu$ CH
3042	22.67	55.01	3044	3048	$\nu$ C <sub>12</sub> H <sub>13</sub>
3022	7.58	58.44			$\nu_{as}$ CH <sub>3</sub>
3010	8.70	95.46			$\nu_{as}$ CH <sub>3</sub>
2936	13.22	136.06	2929	2930	$\nu_s$ CH <sub>3</sub>
1598	0.49	84.94		1607	$\nu$ Ph
1580	0.55	6.51			$\nu$ Ph
1496	4.65	3.39	1494		$\nu$ Ph
1482	8.39	21.23		1480	$\delta_{as}$ CH <sub>3</sub>
1468	5.50	19.53			$\delta_{as}$ CH <sub>3</sub>
1457	10.71	1.55	1455		$\nu$ Ph
1406	9.17	6.93			$\delta$ C <sub>12</sub> H <sub>13</sub>
1367	4.39	6.63	1377		$\delta_s$ CH <sub>3</sub>
1330	3.59	4.42			$\nu$ Ph
1299	8.94	9.43	1303		$\nu$ CH
1214	46.08	70.35	1234	1225	$\nu$ CC
1198	15.93	11.43	1200	1203	$\nu$ CC
1192	2.43	7.68		1190	$\delta$ CH
1177	0.24	7.31		1168	$\delta$ CH
1104	7.15	1.95	1090		$\rho$ CH <sub>3</sub>
1057	2.22	3.14			$\delta$ CH
1053	19.51	2.17	1050		$\rho$ CH <sub>3</sub>
1025	5.67	7.40	1028	1029	$\delta$ CH
999	0.71	20.96		1006	$\nu$ Ph
994	0.12	6.17			$\gamma$ CH
966	10.45	14.22	971	972	$\gamma$ CH
963	0.93	1.05			$\gamma$ C <sub>12</sub> H <sub>13</sub>
918	2.13	1.28	911		$\gamma$ CH
843	0.16	4.57			$\gamma$ CH
771	21.12	2.67		782	$\gamma$ CH
769	3.21	8.94			$\delta$ Ph(X)
698	43.01	0.87	697		$\gamma$ Ph
628	3.12	5.87	617	619	$\nu$ CCl
589	44.27	28.15			$\delta$ Ph
560	38.02	15.32	572	576	$\delta$ CCl
506	19.89	12.93	524		$\gamma$ Ph(X)
408	0.07	0.19		415	$\gamma$ Ph
395	14.69	20.44			$\delta$ Ph(X)
293	3.69	1.68		285	$\gamma$ CCl
251	2.39	3.73			$\delta$ CC(X)
245	0.76	1.02			$\gamma$ CC(X)
206	0.62	0.85			tPh
104	2.23	10.33			tCH <sub>3</sub>
39	0.80	9.77			tCCl

$\nu$ -stretching;  $\delta$ -in-plane deformation;  $\gamma$ -out-of-plane deformation;  $\rho$ -rocking; t-torsion; Ph-phenyl ring; X-substituent sensitive; as-asymmetric; s-symmetric.

$\nu_s\text{CH}_3$  in which all three of the C-H bonds extend and contract in phase. The asymmetric stretching modes of the methyl group are calculated (DFT) to be 3022, 3010  $\text{cm}^{-1}$  and the symmetric mode at 2936  $\text{cm}^{-1}$ . The bands observed at 2929  $\text{cm}^{-1}$  in the IR spectrum and 2930  $\text{cm}^{-1}$  in the Raman spectrum are assigned as stretching modes of the methyl group. Two bending can occur within a methyl group. The first of these, the symmetrical bending vibration, involves the in-phase bending of the C-H bonds. The second, the asymmetrical bending vibration, involves out-of-phase bending of the C-H bonds. The asymmetrical deformations<sup>10</sup> are expected in the range 1400-1485  $\text{cm}^{-1}$ . The calculated values (DFT) of  $\delta_{\text{as}}\text{CH}_3$  modes are at 1482, 1468  $\text{cm}^{-1}$ . In many molecules, the symmetric deformations  $\delta_s\text{CH}_3$  appears with an intensity varying from medium to strong and expected in the range<sup>10</sup> 1380  $\pm$  25  $\text{cm}^{-1}$ . The band observed at 1377  $\text{cm}^{-1}$  in the IR spectrum is assigned as the  $\delta_s\text{CH}_3$  mode. The DFT calculations give  $\delta_s\text{CH}_3$  mode at 1367  $\text{cm}^{-1}$ . Aromatic molecules display<sup>10</sup> a methyl rock in the neighborhood 1045  $\text{cm}^{-1}$ . The second rock<sup>10</sup> in the region 970 $\pm$ 70  $\text{cm}^{-1}$  is more difficult to find among the C-H out-of-plane deformations. In the present case, these  $\nu\text{CH}_3$  modes are calculated at 1104 and 1053  $\text{cm}^{-1}$ . The band observed at 1090, 1050  $\text{cm}^{-1}$  in the IR spectrum assigned as rocking modes of the methyl group. The methyl torsions often assigned in the region<sup>10</sup> 185  $\pm$  65  $\text{cm}^{-1}$ .

The aliphatic CCl bands absorb<sup>13</sup> at 830-560  $\text{cm}^{-1}$  and putting more than one chlorine on a carbon atom raises the CCl wavenumber. The  $\text{CCl}_2$  stretching mode is reported<sup>11,13</sup> at around 738  $\text{cm}^{-1}$  for dichloromethane and scissoring mode  $\delta\text{CCl}_2$  at around 284  $\text{cm}^{-1}$ . Pazdera *et al.*<sup>14,15</sup> reported the CCl stretching mode at 890  $\text{cm}^{-1}$ . For 2-cyanophenylisocyanide dichloride, the  $\nu\text{CCl}$  stretching mode is reported at 870  $\text{cm}^{-1}$  (IR), 877  $\text{cm}^{-1}$  (Raman) and at 882  $\text{cm}^{-1}$  theoretically.[16]. Arslan *et al.*<sup>17</sup> reported  $\nu\text{CCl}$  at 683 (experimental) and at 711, 736, 687, 697  $\text{cm}^{-1}$  theoretically. The deformation bands of CCl are reported at 431, 435,

441 and 443  $\text{cm}^{-1}$ . In the present case, the band at 617 (IR), 619 (Raman) and 628  $\text{cm}^{-1}$  (DFT) is assigned as the stretching mode of CCl. The deformation bands of the CCl are assigned below 400  $\text{cm}^{-1}$ . The C-C stretching modes are assigned at 1214, 1198 (DFT), 1234, 1200 (IR) and 1225, 1203  $\text{cm}^{-1}$  (Raman) as expected<sup>10</sup>. The substituent sensitive modes of the phenyl ring is also identified and assigned (Table 1).

Using the x, y and z components, the magnitude of the dynamic first hyperpolarizability can be calculated by

The complete equation for calculating the magnitude of the dynamic first hyperpolarizability from the Gaussian03 output is given as follows<sup>18</sup>.

$$\beta = [(\beta_{xxx} + \beta_{yyy} + \beta_{zzz})^2 + (\beta_{yyy} + \beta_{yzz} + \beta_{yxx})^2 + (\beta_{zzz} + \beta_{zxx} + \beta_{zyy})^2]^{1/2}$$

To calculate the dynamic first hyperpolarizability, the origin of the Cartesian coordinate system was chosen as the centre of mass of the compound. The calculated first hyperpolarizability of the title compound in an electric field of value 0.001 au is  $1.56 \times 10^{-30}$  esu.. We conclude that the title compound is an attractive object for future studies of non linear optical properties.

## CONCLUSION

The FTIR and FT-Raman spectra of 1-chloroethylbenzene were recorded and analyzed. The harmonic vibrational wavenumbers were calculated theoretically using Gaussian03 set of quantum chemistry codes. The calculated wavenumbers (B3LYP) agree well with the observed wavenumbers. The calculated first hyperpolarizability is reported and the title compound is an attractive object for further studies of nonlinear optics. The predicted infrared intensities and Raman activities are reported

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