Effects of Force Fields on the Equilibrium State of 5,7-Dihydroxy-4-MethylCoumarin Comparative Study of Molecular and Langevin Dynamic Simulations

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Abstract

We have studied the influence of the two force fields Amber and MM+ in both gas and solution environments. Molecular Dynamic (MD) and Langevin Dynamic (LD) simulations of the 5,7-dihydroxy-4-methylcoumarin were performed with the an efficient program. The geometries, interaction energies, bonds, angles, stretch-bends, electrostatic and the Van der Waals (VDW) interactions were carried out in solution and gas phases. The results show that MD and LD simulations are identical. This comparative study shows that the coumarin acquires the low-energy and has the dipolar characteristic under the MM+ field by using the both methods. This molecule reaches its high stable conformation state in solution environment. So, under MM+ field, the simulation gives best results.

Keywords: AMBER, MM+, molecular dynamics (MD), langevin dynamics (LD), coumarin

1. Introduction

Coumarins are an important class of heterocyclic organic compounds (Borges et al., 2005). They have attracted great interest due to their importance in synthetic organic chemistry, and to their possession of diverse pharmacological properties (Hoult & Payá, 1996). Coumarins consist of a carbon-carbon double bond which is fixed as *trans* conformation as in *trans*-stilbene through a lactone structure. From the molecular structure result a strong fluorescence and photostability in most of coumarin derivatives.

In the late of 1950, Wheelock has shown that substitutions on the coumarin structure shifted the fluorescence band. For example, adding a methyl group to the 4-position of 7-hydroxy- or 7-methoxycoumarin moves the fluorescence spectra. Addition of electron-repelling groups in the 4-, 6-, or 7-position or electron-attracting groups in the 3-position shifts the fluorescence band to longer wavelengths.

As consequence, 4-methylcoumarins have been found to possess choleretic, analgesic, anti-spermatogenic, anti-tubercular and diuretic properties (Goel et al., 2007; Pedersen et al., 2007). Hydroxyl derivatives of 4-methylcoumarin are used as optical brightening, dispersed fluorescent lasers (Shan & Xiao, 2004; Wheelock, 1959). Recently, they are employed as potential bio-antioxidants agents (Borges et al., 2005; Hoult & Payá, 1996; Kumar et al., 2005; Kancheva, 2009).

In our earliest work, we have studied the different properties of some coumarin systems by Molecular dynamics (Mesli & Mahboub, 2010). From these results, we are interested in 5,7-dihydroxy-4-methylcoumarin. For this system, we apply Molecular dynamics (MD) and Langevin dynamics (LD) methods. In this article, we detail the method of simulation using AMBER and MM+ as force fields. We will present here recent results both in water and gas environments and we give geometry optimization, energies and dihedrals angles properties.

In the present work, we describe and characterize the molecular structure of 5,7-dihydroxy-4-methylcoumarin by two methods: MD and LD simulations using two force fields: AMBER and MM+. We discuss the computational chemistry results of the coumarin compound then we compare the two calculations methods and force fields effects.

In section 2, we describe the materials and methods used in this study. In the next section, we have detailed the molecular calculations: i) geometry optimization, and ii) dynamic simulations. In the second part, we discuss the evolution of: i) energies, ii) dihedrals angles, iii) temperatures, and iv) geometry optimization in both vacuum

and water environments. We present the calculation results obtained with the two force fields: AMBER and MM+ by Molecular and Langevin dynamics, then we compare all this results. In the last section we give the conclusion.

2. Materials and Methods

Molecular dynamics (MD) is a useful technique to study the kinetics and thermodynamics properties of proteins and other biomolecules. The success of MD is the study process energy of organic compounds as protein, heterocyclic, and important conformational changes in a small time step size.

Solvent effects can be modeled by Langevin dynamics then a molecular system in a liquid environment can be studied. These simulations can be much faster than molecular dynamics. Langevin dynamics simulates the effect of molecular collisions and the resulting dissipation of energy that occurs in real solvents, without explicitly including solvent molecules. This is accomplished by adding a random force and a frictional force to each atom at each time step (van Gunsteren & Berendsen, 1990; Karplus & Petsko, 1990). Because of the implicit treatment of the solvent, this method is particularly well-suited for studying large molecules in solution (Sabbaghzadeh et al., 2010).

Molecular mechanics (MM) force field is a combination of Coulomb and Lennard-Jones interactions to describe all non-bonded interactions (Ponder & Case, 2003). The potential energy is simple and depends on a large number of empirical parameters, which must be obtained from *ab initio* calculations of the electronic structure in small molecules and experimental data. Though AMBER utilizes implicit hydrogen bonding terms (MacKerrell et al., 1998; Cornell et al., 1995; Weiner et al., 1984).

In vacuum, the system was simulated using Molecular and Langevin dynamics (MD and LD) with 0.0001 ps step. For the LD simulation, the frictional coefficient was chosen to 0.05 ps⁻¹. Temperature was kept constant at 300 K. The starting temperature was taken at 100 K and the step at 20 K.

In water, simulations, the system was placed in a box ($20 \times 20 \times 20 \times 20 \text{ Å}$) containing one molecule of 5,7-dihydroxy-4-methylcoumarin and 254 water molecules and cut-off 4 Å. Optimization of the molecule was realized in periodic boundary conditions (PBC). The compound was solvated by added water molecules. The systems were first energy minimized steps with the conjugate gradient algorithm. Then, the position-restrained MD and LD simulation were run 0.5 ps.



Figure 1. Geometry structure of 5,7-dihydroxy-4-methylcoumarin: (a)In vacuum, (b)In water (254 TIP3P water molecules)

3. Results and Discussion

3.1 Geometry Optimization

The geometry optimizations of 5,7-dihydroxy-4-methylcoumarin are realized in vacuum and in water (in PBC with a 253 TIP3P water molecules); using two force fields: AMBER and MM+. The results are given in Table 1. As consequence, the solvated system has lower energies than the isolated structure with both the force fields.

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						Geometry					
State	Force	Bond	Angle	Dihedral	VDW	Stretch-bend	H-bond	Electrostatic	Energy	Gradient	
	field	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/Å.mol)	
Isolated	AMBER	0.5981	1.341	1.9394	5.5091	-	-0.0001	0	9.3881	0.0919	
	MM+	1.0137	4.2602	-8.4056	8.16342	-0.0507	-	0.3398	5.3208	0.0996	
	AMBER	33.8598	37.1948	2.5932	134.435	-	-0.0002	-4143.67	-3935.5886	0.0969	
Solvated	MM+	1 0684	9 4177	-8 2320	-203 361	-0 1741		-481 124	-682 4055	0.0832	

Table 1. Geometry optimization properties of 5,7-dihydroxy-4-methylcoumarin at 300 K with MM+ and AMBER before dynamic simulations

Trans (gauche) conformation of the coumarin structure is confirmed by the two dihedral angles $(C_{10}-O_1-C_2-O_{11})$ (coco), and $C_4-C_3-C_2-O_{11}$) and the distance bonds. In the solvated coumarin, the torsion angle $C_4-C_3-C_2-O_{11}$ changes. So, we obtained an anhedral angle value of -176.406° (Table 2a). This situation is due of the strong collisions between coumarin and water under AMBER.

Table 2a. Structural properties of 5,7-dihydroxy-4-methylcoumarin with AMBER

	Coumarin										
		Isolate	ed					Solvat	ed		
Angle	θ(°)	Dihedral	Φ (°)	Bond	D (Å)	Angle	θ(°)	Dihedral	Φ (°)	Bond	D (Å)
10-1-2	121.855			14-19	2.635	10-1-2	122.067			14-19	2.643
1-2-3	119.921			13-18	2.588	1-2-3	119.652			13-18	2.604
11-2-3	120.323			11-22	2.559	11-2-3	120.815			11-22	2.576
1-2-11	119.746			13-15	2.520	1-2-11	119.510			13-15	2.617
22-3-2	119.529			13-16	3.814	22-3-2	120.010			13-16	3.843
6-13-20	113.535	10-1-2-11	179.706	13-17	2.507	6-13-20	113.235	10-1-2-11	172.303	13-17	2.457
8-14-21	113.372	4-3-2-11	179.735	14-18	2.637	8-14-21	113.320	4-3-2-11	-176.406	14-18	2.630
				10 11	3.48					10 11	3.473
				4 11	3.556					4 11	3.548

Table 2b. Structural properties of 5,7-dihydroxy-4-methylcoumarin with MM+

	Coumarin										
		Isolate	ed					Solvate	ed		
Angle	θ (°)	Dihedral	Φ (°)	Bond	D (Å)	Angle	θ(°)	Dihedral	Φ (°)	Bond	D (Å)
10-1-2	119.341			14-19	2.640	10-1-2	119.291			14-19	2.644
1-2-3	120.407			13-18	2.572	1-2-3	120.676			13-18	2.573
11-2-3	120.330			11-22	2.491	11-2-3	120.013			11-22	2.486
1-2-11	118.966			13-15	2.472	1-2-11	119.307			13-15	2.449
22-3-2	118.363			13-16	3.762	22-3-2	120.013			13-16	3.753
6-13-20	120.236	10-1-2-11	179.927	13-17	2.447	6-13-20	120.253	10-1-2-11	179.890	13-17	2.447
8-14-21	120.147	4-3-2-11	179.986	14-18	2.637	8-14-21	119.943	4-3-2-11	179.681	14-18	2.631
				10 11	3.425					10 11	3.427
				4 11	3.451					4 11	3.448

In vacuum, the minimum potential energies $E_{MD,min}$ and $E_{LD,min}$ calculated for the geometry optimization using AMBER and MM+ force fields are identical (respectively 9.30 Kcal/mol, and 5.31 Kcal/mol). We note that the dihedral factor in potential energy is big in MD simulation than in LD when using AMBER field (table 3a). For the last situation, the system considers the friction coefficient which affects the structure geometry of molecule.

In water, we have obtained the same minimum potential energies $E_{MD,min}$ and $E_{LD,min}$ for both AMBER and MM+ force fields. Variations are visible for bonds, angle, and dihedral energies. The much important difference is happened in the Van der Waals energies.

In MD or LD calculations, the collisions between solute-solvent and solvent-solvent are very important and increase the energy in MM+ field than in AMBER one where the rubbing are present due to the friction coefficient (VDW = 5.4, VDW = 8.16 Kcal/mol) (Tables 3a,b).

Table 3a. Geometry optimization properties of 5,7-dihydroxy-4-methylcoumarin at 300 K with MM+ and AMBER after LD and MD simulations in vacuum

Method	Force	Bond	Angle	Dihedral	VDW	Stretch-bend	H-bond	Electrostatic	Energy	Gradient
	field	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/ Å.mol)
MD	AMBER	0.5646	1.2342	2.1058	5.4013	-	-0.0001	0	9.3058	0.0887
	MM+	1.0135	4.2698	-8.4277	8.1692	-0.0492	-	0.3404	5.3160	0.0935
LD	AMBER	0.5827	1.3156	1.9592	5.4433	-	-0.0001	0	9.3007	0.0891
	MM+	1.0068	4.2728	-8.4283	8.1713	-0.0481	-	0.3402	5.3149	0.0907

Table 3b. Geometry optimization properties of 5,7-dihydroxy-4-methylcoumarin at 300 K with MM+ and AMBER after LD and MD simulations in water (254 TIP3P water molecules)

	Force field	Bond	Angle	Dihedral	VDW	Stretch-bend	H-bond	Electrostatic	Energy	Gradient
Method		(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/mol)	(Kcal/ Å.mol)
MD -	AMBER	0.5762	1.2860	2.0122	5.4313	-	-0.0001	0	9.3056	0.0964
	MM+	1.0118	4.2896	-8.4264	8.1508	-0.0484	-	0.3396	5.3170	0.0911
LD	AMBER	0.5815	1.3126	2.0039	5.4086	-	-0.0001	0	9.3066	0.0960
	MM+	1.0197	4.2836	-8.4268	8.1488	-0.0480	-	0.3397	5.3170	0.0914

3.2 Dynamic Simulation

Dynamic simulations were accomplished at constant temperature, using two force fields: AMBER and MM+. The conjugate gradient (Polak-Ribiere) algorithm was preferred because for the constant dielectric. We have calculated the coumarin dynamic properties in both gas and water. We have employed the TIP3P water molecules model, and chosen the bath relaxation time equals to 0.1 ps. The simulation temperature was fixed to 300 K and the step at 20 K. The run time was 0.5 ps, the step size was 0.0001 ps and the heat time was 0.1 ps. Here, we study the evolution of: i) energies, ii) dihedrals angles, iii) temperatures, and iv) energy deviations in both vacuum and water environments.

At first, we have calculated all energies for the two force fields of AMBER and MM+ simulation then presented them in Figures 2, 3, 4, and 5. The average energies are in good agreement with the simulation accuracy. The sampling results of step-size of MD and LD methods are presented in Figures (a) and (b), respectively.

With AMBER, the total energies (ETOT) are the same in MD and LD simulations in vacuum environment. After equilibration, the MD simulation becomes more stable (respectively 45.25; 45.06 Kcal/mol). The potential energy (EPOT) and the kinetic energy (EKIN) illustrate fluctuations in both methods and vary around (24.77, 20.47) Kcal/mol for MD, and (23.19, 21.86) Kcal/mol for LD (Figure 2). In LD simulation, we observe that the trajectory of potential energy has attained a minimum around 0.3 ps (18.80 Kcal/mol). In the same time, the kinetic energy provides an opposite situation i.e. its reach a maximum (27.10 Kcal/mol).



Figure 2. Evolution of energies for 5,7-dihydroxy-4-methylcoumarin in vacuum with AMBER: (a)MD, (b)LD

With MM+ (in gas), total energy in MD simulation is little high than in LD and much stable after equilibrium time (MD: 42.29, LD: 40.34 Kcal/mol). The kinetic energies have the same value in the two simulations (MD: 16.16, 16.51 Kcal/mol). However, potential energies is high for MD than for LD (MD: 26.12, LD: 23.82 Kcal/mol) (Figure 3).



Figure 3. Evolution of energies for 5,7-dihydroxy-4-methylcoumarin in vacuum with MM+: (a)MD, (b)LD

Under the force field AMBER in water, the evolution of total energy in MD simulation is much stable and little big than in LD one. In MD EPOT energy is greater than in LD. The situation is reversed for kinetic energies (MD: ETOT 41.70, EPOT 23.08, EKIN 18.61 Kcal/mol; LD: ETOT 41.38, EPOT 22.90, EKIN 18.48 Kcal/mol) (Figure 4).



Figure 4. Evolution of energy for 5,7-dihydroxy-4-methylcoumarin in TIP3P water molecules with AMBER: (a)MD, (b)LD

Under the force field MM+ in solution, the position of all energies stay unaltered (MD: ETOT 40.12, EKIN 22.62, EPOT 17.50 Kcal/mol; LD: ETOT 39.77, EKIN 21.88, EPOT 17.89, Kcal/mol). All energies are very stables at time greater than 0.1 ps for both simulations. Here, the fluctuations disappear in spite of interactions between molecules (Figure 5).



Figure 5. Evolution of energies for 5,7-dihydroxy-4-methylcoumarin in TIP3P water molecules with MM+: (a)MD, (b)LD

Comparing our results with analysis performed by two semi empirical (Kancheva et al., 2010), we note that the total energies obtained by MD and LD simulations (using AMBER in water) (respectively 41.70 and 41.38 Kcal/mol) approach that obtained with MNDO (44.57 Kcal/mol). With MM+ in water, MD and LD total energies (40.12, 39.77 Kcal/mol) are lightly up to the PM3 one (36.08 Kcal/mol). In gas, the total energies (MD 45.25, LD 45.06 Kcal/mol; under AMBER) are analogous to that with MNDO (44.57 Kcal/mol). With MM+ in gas, (MD: 42.29, LD: 40.34 Kcal/mol) the situation is similar with those by MM+ in solution. We conclude that our two dynamic methods are suitable for the coumarin system, converge to the MNDO semi-empirical technique and give good results for MM+ force field.

After analyze the evolution of values of dihedrals angles C_{10} - O_1 - C_2 - O_{11} (coco) for coumarin structure, we have observed that the angle is instable in MD simulation, while it has three distinct states of stability in LD calculation with same time intervals (equal to 0.25 ps) in gas. During this simulation run, the torsion angle changes and becomes an anhedral one. The deviations have two states of stability: the first until 0.05 ps for MD and 0.08 ps for LD. The second state has about fully duration of the simulation run in MD and LD calculations (Figure 6).



Figure 6. Evolution of dihedral angle for 5,7-dihydroxy-4-methylcoumarin in vacuum with AMBER: (a)MD, (b)LD

The dihedrals are different for the two simulations in solution. In MD calculation, the angle is most anhedral for the period inferior to 0.25 ps, and then become more dihedral. In LD simulation, the structure acquires inverse situation. The coumarin undergoes deformations because of the cycle torsion (Figure 7). So, we think that the molecule carries an intermediary configuration structure imposed by the delocalization of $C_4=C_3-C_2=O_{11}$.



Figure 7. Evolution of dihedral angle for 5,7-dihydroxy-4-methylcoumarin in TIP3P water molecules with AMBER: (a)MD, (b)LD

In vacuum (under MM+), the dihedrals angles present instability, and oscillate between multiple states (Figure 8).



Figure 8. Evolution of dihedral for 5,7-dihydroxy-4-methylcoumarin in vacuum with MM+: (a)MD, b)LD

By analogy to calculations with MM+ field in vacuum, in solution the evolution of dihedrals is similar for the two simulations (Figure 9).



Figure 9. Evolution of dihedral for 5,7-dihydroxy-4-methylcoumarin in TIP3P water molecules with MM+: (a)MD, (b)LD

The similarities will appear in both methods (MD, LD) and under the two fields (AMBER and MM+) where the angle changes in these simulations due to the interaction of water molecules. In the actual conditions, the torsion angle turns quickly into anhedral angle for short periods (in six states). Consequently, deviations have an important state of stability while the angle has been affected by the collision. The effect is much represented under MM+ field. So, the structure has been disturbed while the calculation runs. At this transition period, the conformation geometry is conserved.

In the present simulation experiments, the temperature was kept at 300 K. In vacuum, the fluctuations are important and temperature is more stable in MD simulations than in LD ones and is situated near 300K (Figure 10a). In water, the situation is much better. We observe that the resulting trajectories present similarity with fewer fluctuations and the value is flat to 300K (Figure 10b).



Figure 10a. Evolution of temperature for 5,7-dihydroxy-4-methylcoumarin in vacuum : with AMBER (a₁)MD, (b₁)LD; with MM+ (a₂)MD, (b₂)LD



Figure 10b. Evolution of temperature for 5,7-dihydroxy-4-methylcoumarin in water : with AMBER (a₁)MD, (b₁)LD; with MM+ (a₂)MD, (b₂)LD

Furthermore investigations lead us to discuss the energy deviations. In the gas, the RMS Deviations of energies fluctuate around 0.15 ps with average value (5 Kcal/mol) for DEKIN and DEPOT and 9.5 Kcal/mol for DETOT using the AMBER field for the MD method. In LD calculation, the fluctuations were happened around 0.2 ps with resulting values of DETOT: 10.5 Kcal/mol, DEPOT: 6.5 Kcal/mol and DEKIN: 5.5 Kcal/mol (Figure 11a). Faraway the maximum, the two curves (DEKIN and DEPOT) converge to same value (5Kcal/mol) in LD simulation.

The same observations are made for DEKIN (MD: 5; LD: 6.5 Kcal/mol) under the MM+ field. Under the MM+ field, we observe that DEPOT deviation curve in MD and LD processes show the same value (5 Kcal/mol). DETOT deviations illustrate similar values (10 Kcal/mol) for both methods.



Figure 11a. Evolution of energy deviations for 5,7-dihydroxy-4-methylcoumarin in vacuum : with AMBER (a₁)MD, (b₁)LD; with MM+ (a₂)MD, (b₂)LD

In solution, the RMS Deviations of energies fluctuate around 0.15 ps in MD with average values of: DETOT (11 Kcal/mol), DEKIN and DEPOT (5.5 Kcal/mol); In LD, they fluctuate around 0.1ps with average values of DETOT (8.5 Kcal/mol), DEKIN and DEPOT (5 Kcal/mol) when using the AMBER field (Figure 11b). As consequence, the two trajectories (DEKIN and DEPOT) converge to the same value 4 Kcal/mol.

The situation will be unchanged when we used the MM+ field: DETOT (10.5 Kcal/mol), DEKIN (6.5 Kcal/mol) and DEPOT (4.5 Kcal/mol).



Figure 11b. Evolution of energy deviations for 5,7-dihydroxy-4-methylcoumarin in water: with AMBER (a₁)MD, (b₁)LD; with MM+ (a₂)MD, (b₂)LD

DEKIN and DEPOT will stay steady around constant value for both simulations under the two fields (AMBER and MM+) in all environments. In solution and under MM+, after the equilibration time (0.5 ps) the corresponding deviations become collinear. In solution, DETOT reach its greatest value 11 Kcal/mol for both fields. As consequence, the coumarin conformation is more stable in solution.

3.3 Conformation of Coumarin

To provide a better estimate of conformations, we should search the conformational space in reasonable computing time. So, we run the simulations at the constant temperature, and then we run a geometry optimization on each structure. In Table 4, we have grouping the following characteristics which are determined after geometry optimization (and dynamic simulations): the dipole moment (μ), the electrostatic potential (Ep) and the total energy (Et).

Table 4. Properties of 5,7-dihydroxy-4-methylcoumarin at 300 K with MM+	 and AMBER 	after LD a	and MD
simulations and geometry optimization in two environments: Dipole moment	, Electrostatic	potential an	nd Total
Energy			

		F	Properties					
State	State Method		μ (D)	Electrostatic potential (Kcal/mol)	Total Energy (Kcal/mol)			
	MD	AMDED	0.000	0.000	9.306			
Isolated	LD	AMBER	0.000	0.000	9.300			
	MD		0.908	0.340	5.316			
	LD	MM+	0.909	0.340	5.315			
	MD	AMDED	0.000	0.000	9.306			
Solvated	LD	AMBER	0.000	0.000	9.306			
	MD	MAL	0.903	0.339	5.317			
	LD	MM+	0.909	0.339	5.317			

From these results, we observe that the coumarin has the zero value for the dipole moment and the electrostatic potential when using AMBER field. However under MM+, the dipole moment and the electrostatic potential are unlike of zero and get the same value in both simulations ($\mu = 0.908$ D, Ep = 0.34 Kcal/mol).

So, we have plotted the electrostatic potential in 2D contours with the value obtained after geometry optimization (0.34 Kcal/mol). Finally, we have compared representation structures to determine low-energy conformations (Figures 12a, 12b).

The electrostatic potential representation clearly shows that the σ -withdrawing properties of the oxygen in the lactone reduce the π -donating and deactivate the ring. So, the molecule grows to be polarizable.

Combining these properties (Table 4) and the representations (Figures 12), we note that under MM+ the coumarin acquires its low-energy value (5.316 Kcal/mol). At the same time, the molecule has the dipolar characteristic ($\mu = 0.908$ D). On this situation, the molecule reaches its high stable conformation state.



(a₁)
 (b₁)
 (a₂)
 (b₂)
 Figure 12a. Electrostatic potential 2D representation of Coumarin structure after Dynamic Simulations and geometry optimization in vacuum: with AMBER (a₁)MD, (b₁)LD; with MM+ (a₂)MD, (b₂)LD



(a₁)
 (b₁)
 (a₂)
 (b₂)
 Figure 12b. Electrostatic potential 2D representation of Coumarin structure after Dynamic Simulations and geometry optimization in water: with AMBER (a₁)MD, (b₁)LD; with MM+ (a₂)MD, (b₂)LD

4. Conclusion

We have studied the evolution of the geometry optimized properties for 5,7-dihydroxy-4-methylcoumarin. We have chosen the MD and LD methods to predict much better the characteristics. The geometries, interaction energies, bonds, angles, stretch-bends, electrostatic and the VDW interactions were carried out in solution and in gas phase. We have calculated the thermodynamic and structural properties for both the two force fields Amber and MM+.

The total energy is more stable in MD simulations under the AMBER and MM+ fields. Under AMBER and MM+ in water, we note that the total energies obtained by MD and LD simulations are similar. In gas, the situation is similar. We conclude that our two dynamic methods are suitable for the coumarin system, and give good results for MM+ force field.

The deviations of energies stay steady around constant value for both simulations under the two fields (AMBER and MM+) in all environments. In solution, the total energy deviation reaches its greatest value for both fields. As consequence, the coumarin conformation is more stable in solution.

The torsion angle changes in these simulations due to the interaction of water molecules and turns quickly into anhedral angle for short periods. At this transition period, the conformation geometry is conserved. LD method is more effective for calculate the dihedral angle with MM+ field in solution environment.

We conclude that our simulation under MM+ field gives best results; the coumarin acquires the low-energy and has the dipolar characteristic. So, the molecule reaches its high stable conformation state in solution environment under MM+ field.

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