This discussion paper is/has been under review for the journal Atmospheric Measurement Techniques (AMT). Please refer to the corresponding final paper in AMT if available.

Flux correction for closed-path laser spectrometers without internal water vapor measurements

R. V. Hiller¹, C. Zellweger², A. Knohl³, and W. Eugster¹

¹Institute of Agricultural Sciences, ETH Zurich, Zürich, Switzerland

²Empa, Materials Science & Technology, Dübendorf, Switzerland

³Institute of Bioclimatology, Georg-August University of Göttingen, Göttingen, Germany

Received: 13 December 2011 – Accepted: 14 December 2011 – Published: 10 January 2012

Correspondence to: R. V. Hiller (hillerr@ethz.ch)

Published by Copernicus Publications on behalf of the European Geosciences Union.

351

Abstract

Recently, instruments became available on the market that provide the possibility to perform eddy covariance flux measurements of CH₄ and many other trace gases, including the traditional CO₂ and H₂O. Most of these instruments employ laser spectroscopy, where a cross-sensitivity to H₂O is frequently observed leading to an increased dilution effect. Additionally, sorption processes at the intake tube walls modify and delay the observed H₂O signal in closed-path systems more strongly than the signal of the sampled trace gas. Thereby, a phase shift between the trace gas and H₂O fluctuations is introduced that dampens the H₂O flux observed in the sampling cell. For instruments that do not provide direct H₂O measurement in the sampling cell, transfer functions from externally measured H₂O fluxes are needed to estimate the effect of H₂O on trace gas flux measurements. The effects of cross-sensitivity and the damping are shown for an eddy covariance setup with the Fast Greenhouse Gas Analyzer (FGGA, Los Gatos Research Inc.) that measures CO₂, CH₄, and H₂O fluxes. This instrument is technically identical with the Fast Methane Analyzer (FMA, Los Gatos Research Inc.) that does not measure H₂O concentrations. Hence, we used measurements from a FGGA to derive a modified correction for the FMA accounting for dilution as well as phase shift effects in our instrumental setup. With our specific setup for eddy covariance flux measurements, the cross-sensitivity counteracts the damping effects, which compensate each other. Hence, the new correction only deviates very slightly from the traditional Webb, Pearman, and Leuning density correction, which is calculated from separate measurements of the atmospheric water vapor flux.

1 Introduction

The eddy covariance technique (EC) is the established method to directly measure the exchange of trace gases, such as CO₂ and H₂O, between the land surface and the atmosphere (Baldocchi, 2003). Recently, new instruments became available that are

suitable to measure fluxes of additional trace gases, such as CH₄ or N₂O (e.g. Eugster and Plüss, 2010; Smeets et al., 2009; Neftel et al., 2010; Hendriks et al., 2008; Kroon et al., 2007; Eugster et al., 2007). However, by careful evaluation of the suitability of the instrument for the EC method, it becomes clear that some adaptations of the standard processing techniques, e.g. accounting for cross-sensitivity effects with H₂O, are needed (McDermitt et al., 2010).

Most new instruments apply laser spectroscopy to measure trace gases and two general types of instruments can be distinguished: (1) closed-path instruments pull air through a sampling cell where the air is measured, and (2) open-path instruments measure the composition of the air directly in the atmosphere. Both types of instruments measure the decay of a laser beam over a known distance (Baer et al., 2002; McDermitt et al., 2010). In closed-path instruments, the beam is reflected by mirrors to achieve a path length of several kilometers (Baer et al., 2002; Crosson, 2008). The definition of the spectral line assumes that the gas matrix does not change. Atmospheric water content can however vary considerably and thus broaden the spectral line. Hence, the instrument or the post-processing of the data have to account for this effect (Rella, 2010). This phenomenon was described for a Quantum Cascade Laser (QCL, Aerodyne Research Inc.) by Neftel et al. (2010), and for a Fast Methane Analyzer (FMA, Los Gatos Research Inc.) by Tuzson et al. (2010). Picarro Inc. discussed the issue for their laser spectroscopy-based instrument G1301 in a white paper (Rella, 2010). The same artifact also influences open-path instruments such as the LI-7700 (Li-Cor) where a combined correction accounts for the dilution as well as for the additional H₂O interference on the measured CH₄ flux (McDermitt et al., 2010). With closed-path instruments, one solution would be to dry the air prior to sampling. However, this tends to have a negative impact on the frequency response and high-frequency performance of the sensor (Griffis et al., 2008). Alternatively, a correction procedure could solve the cross-sensitivity issue in a more elegant way as will be discussed in this paper.

Since the FMA does not measure H₂O, the correction of the density effect caused by water vapor molecules is essential to improve the accuracy of the flux measurements.

353

Therefore, the H₂O flux measured by a separate system is typically used to apply this correction. The traditional density flux correction by Webb et al. (1980) employs the water vapor flux measured in the same sampling cell as the other gas of interest. The water vapor signal however gets more strongly dampened from the inlet to the sampling cell (Ibrom et al., 2007a) than a less polar gas such as CH₄. Thus, substituting the water vapor flux in the original correction term with the water vapor flux measured separately with a second instrument and at some distance from the first instrument might overcorrect the density effect. To find out whether this is a small (negligible) or large and important effect was the goal of this study.

We used a Fast Greenhouse Gas Analyzer (FGGA, Los Gatos Research Inc.) to assess the phase shift and damping of the H₂O signal with respect to the CH₄ signal and in relation to an externally measured H₂O flux. Based on these findings, two transfer functions will be suggested to (1) correct the H₂O flux measured by the FGGA to represent atmospheric conditions; and (2) its inverse function that allows to translate an externally measured H₂O flux to the conditions in the sampling cell which then can be used to correct the CH₄ flux for the density fluctuations. This correction procedure can be used to correct CH₄ flux measurements performed with a FMA which does not provide internal H₂O measurements.

2 Background

This study closely follows the concept of the density flux correction accounting for density effects due to heat and water vapor fluctuations that was introduced by Webb et al. (1980). We will shortly introduce the concept of the theory and highlight the importance in relation to this study. The original concept will be referred to as WPL correction thereafter.

For water vapor, the original correction for open-path analyzers reads (WPL Eq. 25):

$$E = (1 + \mu\sigma)\{E_{\text{raw}} + (\overline{\rho_v}/\overline{T})\overline{w'T'}\}, \quad (1)$$

354

with E [$\text{kg m}^{-2} \text{s}^{-1}$] the corrected water vapor flux, $\mu = \frac{m_a}{m_v} \approx 28.97/18$ [-] the ratio of molecular masses of dry air and water vapor, $\sigma = \frac{\rho_v}{\rho_a}$ [-] the ratio of the vapor density and the density of dry air, E_{raw} [$\text{kg m}^{-2} \text{s}^{-1}$] the measured water vapor flux, $\bar{\rho}_v$ [kg m^{-3}] the density of dry air, T [K] the air temperature, and $\overline{w'T'}$ [K m s^{-1}] the sensible heat flux in kinematic units.

For trace gas fluxes, buoyancy effects from water vapor additionally influence the measured flux. Thus, Webb et al. (1980) introduced the following correction (WPL Eq. 24):

$$F = F_{\text{raw}} + \Delta F = F_{\text{raw}} + \underbrace{\mu(\bar{\rho}_c/\bar{\rho}_a)E}_{\text{water vapor}} + \underbrace{(1 + \mu\sigma)(\bar{\rho}_c/\bar{T})\overline{w'T'}}_{\text{heat}}, \quad (2)$$

where F [$\text{kg m}^{-2} \text{s}^{-1}$] is the corrected trace gas flux and F_{raw} the measured one. $\bar{\rho}_c$ [kg m^{-3}] denotes the density of the particular trace gas under consideration. The curly brackets summarize the correction terms for density fluctuations caused by the buoyancy effect of water vapor and by sensible heat. Taken together, they constitute ΔF , the WPL correction term.

The sensitivity of the WPL correction itself largely depends on the ratio between the concentration of the particular trace gas under consideration and the corresponding flux. Largest WPL corrections are expected for small fluxes at comparably high background concentrations (Ibrom et al., 2007b). If $\bar{\rho}_c$ is factored out in Eq. (2), it yields:

$$F = F_{\text{raw}} + \bar{\rho}_c \left\{ \underbrace{(\mu/\bar{\rho}_a)E}_{\text{water vapor}} + \underbrace{[(1 + \mu\sigma)/\bar{T}]\overline{w'T'}}_{\text{heat}} \right\}. \quad (3)$$

environment

The term in the curly bracket to the right of $\bar{\rho}_c$ is only influenced by the measurements of the environment and does not depend on the density of the trace gas nor

355

the trace gas flux. Therefore, the WPL correction (ΔF), the difference between the measured and the corrected flux, depends on the density of the trace gas multiplied by a term related to the environment conditions where the measurements were performed. This term includes the water content of the air, the water vapor flux, air temperature, and the sensible heat flux. The relative magnitude of the WPL correction compared to the measured flux can be assessed by the ratio $\Delta F/F_{\text{raw}}$. This is equal to $(\bar{\rho}_c \times \text{environment})/F_{\text{raw}}$ and hence the relative WPL correction is highest for high background concentrations ($\bar{\rho}_c$) combined with small fluxes F_{raw} . Detto et al. (2011) presented in their Fig. 6 the influence of the latent heat flux on CH_4 flux for a closed-path instrument, where the heat related component can be neglected based on the evidence that temperature fluctuations are almost completely removed in the closed-path instruments. However, since the WPL correction is an additive term and not a scaling factor, one should be careful with expressing the WPL correction as a percentage of the measured flux. This becomes most obvious at small fluxes, where the WPL correction may lead to a sign change of the flux.

For closed-path instruments, the environment relevant for the WPL correction relates to the environment in the sampling cell, not to the ambient environment where the intake of the gas inlet was placed. Therefore, the water vapor and sensible heat flux in the cell define the magnitude of the WPL correction. In case of the sensible heat flux, the tubing from the inlet to the measurement cell and the waste heat of the instrument cancel the natural short time temperature fluctuations. Hence, $\overline{w'T'}$ gets very small and the density effects due to heat transfer is negligible (Leuning and Judd, 1996). The correction term reduces to density effects due to water vapor. This term depends on various factors. If the air is dried prior to measurement or if concentrations are measured in units of dry mole fractions, this correction can be dismissed too. Otherwise, the term applies, but the instrumental setup may influence the measured H_2O flux and consequently also the magnitude of the correction to be applied. Damping of the H_2O signal reduces the H_2O flux in a closed-path analyzer and thus the WPL correction would be overestimated by using H_2O flux estimates taken from a second instrument

356

(Ibrom et al., 2007b). The investigation of such effects is important to obtain the highest accuracy of trace gas flux measurements (Smeets et al., 2009). Signal damping due to amplitude effects (Massman and Clement, 2005) as well as phase shifts (Ibrom et al., 2007b) should be attributed.

5 The attenuation of the water vapor flux measured by a closed-path instrument depends on the relative humidity (Ibrom et al., 2007a; Mammarella et al., 2009). This however does not apply to CO₂ and CH₄ fluxes, since the molecules are not as sticky as H₂O and hence do not strongly interact with the intake hose's inner wall. Thus, the transportation of the air through the instrument leads to a decoupling of the CO₂ or CH₄ and H₂O signal (Ibrom et al., 2007b). This decoupling consequently results
10 in a reduction of the water vapor flux at the time lag that was determined for the CO₂ flux. Therefore, the WPL correction needs to be adjusted in a way that the H₂O flux in the cell is approximated, as for example shown by Ibrom et al. (2007b), Smeets et al. (2009), and Massman (2005).

15 3 Methods

3.1 Instruments

The Fast Methane Analyzer (FMA) and the Fast Greenhouse Gas Analyzer (FGGA, both Los Gatos Research Inc., Mountain View, CA, USA) are capable of measuring CH₄ mole fractions. Both instruments operate at high measurement frequencies up
20 to 20 Hz as required for eddy covariance measurements. The construction of the two instruments is identical except for a second laser in the FGGA that measures CO₂. H₂O is also only measured by the FGGA, but retrieved from the same laser absorption spectrum as CH₄. With the help of measurements with the FGGA, the influence of H₂O in the sampling cell of the FMA can be assessed.

3.2 Dilution and cross-sensitivity experiment

Linearity and precision of the FMA and FGGA were tested in the laboratory. The instruments were connected to a common air source. Eleven concentration levels at regular intervals between 0 and 4 ppm for CH₄ and 16 between 0 and 1500 ppm for
5 CO₂ were produced by diluting CH₄ and CO₂ with CH₄ and CO₂ free air, respectively. The concentration was kept constant for 20 min in case of CH₄ and 15 min for CO₂. After verifying the linearity of the measurements, the precision of the CH₄ measurements was checked against two National Oceanic & Atmospheric Administration (NOAA) standards (CA04580 and CA05316). H₂O was calibrated against a LI-610
10 dew point generator (Li-Cor, Lincoln, NB, USA). Further, dry air with a known CH₄ and CO₂ concentration was humidified to check for dilution effects and cross-sensitivity of CH₄ and CO₂ measurements to H₂O. H₂O concentrations were set to 9 different levels between 0 and 25 000 ppm and kept constant for 20 min. The air temperature in the laboratory was controlled during all experiments.

15 3.3 Flux measurements

3.3.1 Site

The flux measurements presented in this study were acquired at the ETH agricultural research station Chamau (8° 24' 38" E, 47° 12' 37" N, 400 m a.s.l., Switzerland) in summer 2009. The station is located in the broad prealpine Reuss Valley. Two flux towers, a
20 mobile and a long-term one, were situated on a temperate, intensively managed grassland with mixed ryegrass-clover vegetation raised for forage (Zeeman et al., 2010).

3.3.2 Instruments and setup

The mobile tower consisted of a Fast Greenhouse Gas Analyzer (FGGA, Los Gatos Research, Inc., Mountain View, CA, USA) measuring concentrations of CO₂, CH₄, and
25 H₂O and a Solent R2A ultrasonic anemometer (Gill Instruments Ltd., Lymington, UK),

thereafter abbreviated as sonic. The sonic was installed at a height of 1.7 m with the inlet for the FGGA attached to the sonic boom 25 cm below the instrument head. The sample was drawn through a 6.7 m long tube with 8 mm inner diameter (Synflex-1300, Eaton Performance Plastics, Cleveland, OH, USA) followed by two serially mounted
 5 particle filters with integrated droplet separator (SMC, Japan, model AF30-F03 with 5 μm filter and model AFM30-F03 with 0.3 μm filter, respectively). The FGGA was connected to an external dry vacuum scroll pump (BOC Edwards XDS-35i, Crawley, UK) drawing 30.5 slpm to ensure turbulent flow in the sampling tube. An embedded computer (Advantech ARK-3381, Taipei, Taiwan) running a Linux system recorded the
 10 raw data at 20.8 Hz. Eugster and Plüss (2010) and Tuzson et al. (2010) provide a more detailed description of the system.

The long-term tower measured CO_2 and H_2O fluxes by an open-path Infrared Gas Analyzer (IRGA, Li-7500, Li-Cor, Lincoln, NB, USA) and a Solent R3 ultra sonic anemometer (Gill Instruments Ltd., Lymington, UK) at 2.40 m above ground, about
 15 5 m separated from the mobile tower. A more detailed description of this system can be found in Zeeman et al. (2010).

Additionally, several meteorological variables were measured: air temperature and relative humidity of the air (HygroClip S3, Rotronic AG, Bassersdorf, Switzerland), and atmospheric pressure (LI-7500, Li-Cor, Lincoln, NB, USA).

20 Even though measured by the HygroClip, relative humidity values used in the evaluations were derived from the LI-7500 measurements (RH_{IRGA}) due to higher accuracy of this instrument, whereas the HygroClip data were used for cross-checking the results.

3.3.3 Flux processing

The statistical software R (Version 2.12.1, R Development Core Team, 2010) was used
 25 for all analyses. This includes the processing of the flux data from the mobile system with the R package *Reddy* (under development and available from the authors upon request). CO_2 and H_2O fluxes of the long-term tower were calculated by the in-house

software *eth-flux* (Mauder et al., 2008) which is also available from the authors. In *eth-flux* and *Reddy*, a two-way rotation of the wind vector was performed for each 30-min interval (McMillen, 1988). The time lag between the vertical wind speed w and the scalar was determined for each 30-min block such that the correlation maximized. If not
 5 stated differently, this time lag was applied before calculating the respective flux. Fluxes were calculated as the covariance between the two time series following the rules of Reynold's averaging ($\overline{w'c'}$) (Baldocchi, 2003). The CH_4 fluxes at our site were often very small and close to the detection limit of the system. Thus the time lag between w and the CH_4 concentration measurements was difficult to detect since no clear peak
 10 in the cross-correlation function was found. For periods with a detectable CH_4 time lag, it was very similar to the one of CO_2 . Thus, we used the time lag found for CO_2 to calculate the CH_4 fluxes. The H_2O flux measured by the FGGA was calculated in two ways: (1) the time lag maximizing the correlation between w and the H_2O signal ($\overline{w'\text{H}_2\text{O}'_{\text{FGGA}}}$) and (2) with the time lag found for CO_2 ($\overline{w'\text{H}_2\text{O}'_{\text{FGGA}, \text{lag}(w'\text{CO}_2')}}}$). Prior to
 15 the flux calculations, all scalars of the FGGA were post-calibrated as described below.

Flux data were filtered according to the following criteria. For the mobile system, fluxes were rejected if the cell pressure dropped below 100 Torr or the number of available trace gas measurements was below 20 000 of 36 000 within the 30-min averaging window. Further, flux records were removed if *Reddy* automatically detected the maximal
 20 or minimal time lag within the lag detection window, namely 10 and 40 records for $\overline{w'\text{CH}_4'}$, 15 and 40 records for $\overline{w'\text{H}_2\text{O}'}$, and 10 and 30 records for $\overline{w'\text{CO}_2'}$. For the long-term station, fluxes were dismissed if the automated gain control (AGC) value by the LI-7500, the so called window dirtiness, exceeded 70 % or was not a value dividable by 6.25. Both criteria are indicators for rain or dew disturbing the measurement path.
 25 Since the LI-7500 reports the AGC as a four-bit variable, 16 different values between 0 and 100 can be reported which results in incremental steps of 6.25. Thirty-minute averages between these steps indicate temporal changes of the AGC value within the averaging period, e.g. introduced by slight rain or an insect passing the sensor path during a short part of the averaging interval. Further, fluxes were rejected when the

friction velocity dropped below 0.08 m s^{-1} as suggested by Zeeman et al. (2010) for this site.

3.3.4 Cospectra and damping loss correction

Before the calculation of the cospectra, a two way rotation was applied to the wind vectors (McMillen, 1988) and a time lag was implemented to maximize the correlation of the two time series. The scalar concentrations were converted to dry mole fractions and a correction for additional cross-sensitivity to H_2O effects (see Sect. 4.1) was applied. Since no cospectral analysis function exists for R , we wrote a new function downloadable from <http://www.swissfluxnet.ch/R/>. The procedure follows the suggestion by Stull (1988, p.310ff) and was double-checked against the algorithm for Fast Fourier Transformation suggested by Press (1988). Thereafter, a bandwidth averaging was performed on the output (Kaimal and Gaynor, 1983) and the high frequency flux losses were quantified following Eugster and Senn (1995). A slight modification of the theoretical cospectral model introduced by Kaimal et al. (1972) is damped according to the inductance $L [\text{s}^{-1}]$ in an alternating-current circuit to match the calculated cospectra. This corresponds to the standard first-order damping known from similar approaches, e.g. by Horst (1997). For unstable conditions, the cospectral approximation reads (Eq. 29 in Eugster and Senn, 1995):

$$\frac{fCo_{w,c_L}(f)}{w'c'_{L=0}} = \begin{cases} 10.53n / \left[(1 + 4\pi^2 f^2 L^2)(1 + 13.3n)^{1.72} \right], & n \leq 1.0 \\ 4.21n / \left[(1 + 4\pi^2 f^2 L^2)(1 + 3.8n)^{2.4} \right], & n \geq 1.0 \end{cases}, \quad -2 \leq \zeta \leq 0, \quad (4)$$

and for stable conditions (Eq. 30 in Eugster and Senn, 1995):

$$\frac{fCo_{w,c_L}(f)}{w'c'_{L=0}} = \frac{1}{1 + 4\pi^2 f^2 L^2} \times \frac{0.81(n/n_0)}{1 + 1.5(n/n_0)^{2.1}}, \quad 0 < \zeta \leq 2 \quad (5)$$

361

with

$$n_0 = \begin{cases} 0.23, & -2 \leq \zeta \leq 0 \\ 0.23(1 + 6.4\zeta)^{3/4}, & 0 < \zeta \leq 2 \end{cases}. \quad (6)$$

Here, $\zeta [-]$ denotes the Monin-Obukhov stability parameter and n the normalized frequency $n = fz/\bar{u}$ with f [Hz] the natural frequency, z [m] the measuring height above zero-plane displacement, and $\bar{u} [\text{m s}^{-1}]$ the mean horizontal wind-speed along the mean wind direction.

Cospectra were calculated from 1-h blocks of high frequency data for the whole measurement period. The damping factors L were retrieved by a non-linear least squares fit (function `nls()` in R) to the individually calculated cospectra. For further analyses, L was only used for periods when a detectable time lag was found within the predefined window and the fit of the idealized damped cospectrum on the calculated cospectrum was significant at $p < 0.1$. L was originally set constant over periods without system maintenance where the measured tracer was NO_2 (Eugster and Senn, 1995). Especially for the water vapor flux, additional dampening of the signal is however observed depending on the relative humidity of the air (Ibrom et al., 2007a; Mammarella et al., 2009). Therefore, L was adjusted according to the following empirical model with relative humidity RH as predictor:

$$L = a + b \times \text{RH} + c \times \text{RH}^2, \quad (7)$$

where a , b , and c were obtained by conditional linear regression fitting. To reduce the scatter in the data, the automatically determined L values were binned into relative humidity classes of 5% width. The model was fitted through the medians of each class. In contrast to L for the H_2O fluxes, analyses of the empirically determined L for the CH_4 and CO_2 fluxes resulted in a constant value. The measured fluxes were corrected with the derived L for damping losses following the method by Eugster and Senn (1995).

4 Results

4.1 Dilution and cross-sensitivity experiment

The linearity of the instruments was tested in a dilution experiment (data not shown). Comparing the median of the measurements at each level against the given concentrations, R^2 yields 0.999997 for the FMA and 0.999995 for the FGGA. This implies that a two-point calibration is sufficient to calibrate the instrument. CH_4 was calibrated against two NOAA standards (1.713 ppm and 1.905 ppm). The linear relationship through these two points (in ppm) yielded:

$$\begin{aligned}\text{CH}_4 &= -0.0147 + 1.0002 \times \text{CH}_{4,\text{FMA}} \\ \text{CH}_4 &= -0.0024 + 0.9774 \times \text{CH}_{4,\text{FGGA}}\end{aligned}\quad (8)$$

- 10 The calibration of the FGGA H_2O signal determined the following polynomial of the second order (in ppm):

$$\begin{aligned}\text{H}_2\text{O}_{\text{ref}} &= (-52.36 \pm 8.01) + \\ & \quad (1.011 \pm 0.002) \times \text{H}_2\text{O}_{\text{FGGA}} + \\ & \quad (0.00000239 \pm 0.00000007) \times \text{H}_2\text{O}_{\text{FGGA}}^2,\end{aligned}\quad (9)$$

- 15 where $\text{H}_2\text{O}_{\text{FGGA}}$ refers to the H_2O humid mole fraction measured with the FGGA and $\text{H}_2\text{O}_{\text{ref}}$ to the dry mole fraction of the dew point generator, both measured in or converted to ppm. The uncertainty in the parameter estimate is given with $\pm$ one standard error and the R^2 yields 0.999. The FGGA overestimates H_2O concentrations at low concentrations up to ≈ 5000 ppm and thereafter overestimates them. For the following analysis, the water vapor signal was post-calibrated according to the parameter fit in Eq. (9).

- 20 The dilution experiment described in Sect. 3.2 showed a drop of the CO_2 and CH_4 humid mole fractions that was triggered by the changes of humidity. Most of the drop can be explained by the dilution of the air by water vapor (difference of the black and dark grey points in Fig. 1). The dark grey points correspond to the dry mole fraction

363

reported by the FGGA. However, a small part of the drop, the difference between dark and light grey points in Fig. 1, was not caused by dilution and hence must result from different processes. Rella (2010) suggested the following equation to calculate the dry mole fraction for the Picarro G1301 analyzer, here adapted for the FGGA instrument:

$$5 \quad c = c_{\text{FGGA}} / (1 + a \times \text{H}_2\text{O} + b \times \text{H}_2\text{O}^2), \quad (10)$$

where H_2O is the corrected H_2O dry mole fraction measured by the FGGA and c_{FGGA} the measured humid mole fraction of the issued dry mole fraction c of CO_2 or CH_4 . Table 1 shows the empirical coefficients for Eq. (10) which combines the dilution effect as well as other water dependent processes influencing the measured concentration.

- 10 The additional water dependent processes affect the measured concentration in a similar way as the dilution effect. To correct the effect of density fluctuations caused by the dilution, the WPL correction (Eq. 2) can be applied. The dilution is a linear process while Eq. (10) additionally includes the square of H_2O . The derivation of the WPL correction would be severely complicated if the square term is included. Since the performance of a simple linear model, same as Eq. (10) but without term $b \times \text{H}_2\text{O}^2$, is similar (compare corresponding R^2 in Tables 1 and 2), we decided to use the latter for the further analysis. The details for the individual gases and instruments can be found in Table 2.

- 15 An additional fit of Eq. (10) was determined for CH_4 measured by the FMA with H_2O taken from the measurements of the FGGA. The results only marginally differs from Fig. 1b (not shown). The results for the fit can be found in Table 1 and 2.

4.2 Damping loss correction

- The automatically fitted damping factor L for $\overline{w' \text{H}_2\text{O}'}_{\text{FGGA, lag}(w' \text{CO}_2')}$ shows a clear dependency on the relative humidity (see Fig. 2). Such dependency of L on the relative humidity would also be found if the L for $\overline{w' \text{H}_2\text{O}'}$ at the time lag maximizing $\overline{w' \text{H}_2\text{O}'}$ were used (not shown). The increase in L with increasing relative humidity can by

364

explained by sorption and desorption effects. The variability of the $\overline{w'H_2O'}$ parallel increases with increasing L . This can be explained by lower water vapor fluxes at high relative humidity and hence a less well defined cospectrum. The automatic fits of L are less robust and result in a higher scattering. Since outliers would strongly influence the parameter fit for the relationship between L and the relative humidity, L was binned into relative humidity classes of 5 % width. A polynomial of the 2nd order was fitted through the medians of the individual classes, which yields:

$$L_{\overline{w'H_2O'}_{\text{FGGA, lag}(w'CO_2')}} = 1.07 - 2.99 \times 10^{-2} \times RH_{\text{IRGA}} + 3.62 \times 10^{-4} \times RH_{\text{IRGA}}^2, \quad (11)$$

with RH_{IRGA} [%] being the relative humidity obtained from the air temperature measured by the HygroClip and the H_2O concentration measured by the LI-7500 IRGA. This relationship between L and the relative humidity was used in the flux processing to determine the damping factor for the water vapor flux and correct the flux there-with for damping losses. No dependency on relative humidity was found for the CO_2 and CH_4 fluxes (not shown). For these two trace gases, L seems to be a constant as proposed in the original paper by Eugster and Senn (1995).

4.3 Signal desynchronisation

Different time lags to maximize the covariance $\overline{w'H_2O'}$ and $\overline{w'CO_2'}$ were found. Based on the assumption that no time lag would be observed under ambient conditions, we guess that the signal of CO_2 and H_2O was desynchronized on the way from the inlet to the measurement cell. Figure 3 shows the histogram of the determined time lags between w and the particular scalar concentration. In general, the peak of $\overline{w'H_2O'}$ moved towards longer lags compared to $\overline{w'CO_2'}$. Also the distribution of $\overline{w'H_2O'}$ lags was wider than the one for $\overline{w'CO_2'}$. For $\overline{w'CH_4'}$, the distribution of time lags was almost uniform because the flux at our site was very low. However, the comparison of the time lag distribution of $\overline{w'CH_4'}$ in the two panels in Fig. 3 reveals that there is a slight

365

difference seen in the range where the $\overline{w'H_2O'}$ time lag was most frequent (20–30 records). For CH_4 humid mole fractions, more cases with a lag in this range were observed than for dry mole fractions. Since the CH_4 flux was very low at our site, the water-introduced fluctuations in the CH_4 signal were sometimes stronger than the natural variations of CH_4 , which resulted in lags for CH_4 being identical to those of $\overline{w'H_2O'}$.

For $\overline{w'H_2O'}$, the time lag clearly depended on the relative humidity (Fig. 4a) while it stayed constant for $\overline{w'CO_2'}$. This explains the wider distribution of the peak for time lags of $\overline{w'H_2O'}$ compared to $\overline{w'CO_2'}$ in Fig. 3. Nevertheless, more records lacked a detectable time lag during conditions of high relative humidity for $\overline{w'CO_2'}$, and lags of $\overline{w'H_2O'}$ were hardly detectable at $RH_{\text{IRGA}} > 95\%$.

The ratio between $\overline{w'H_2O'}_{\text{FGGA, lag}(w'CO_2')}$ and $\overline{w'H_2O'}_{\text{FGGA}}$ decreased with increasing relative humidity as a result of higher temporal separation of the optimal lags of CO_2 and H_2O (see Fig. 4b). The median observed $\overline{w'H_2O'}$ flux losses were increasing up to 30 % at highest relative humidity. However, $\overline{w'H_2O'}$ was smaller in these conditions than during dryer periods, and thus the absolute flux difference was minimal. Looking at the cospectra of $\overline{w'H_2O'}$, the flux difference appears visible as an additional damping of the signal (see Fig. 5). This observation agrees well with the relative humidity dependent damping losses found for $\overline{w'H_2O'}$ (Eq. 11).

The H_2O concentration measurements of the FGGA tend to be lower than those measured by the LI-7500 IRGA. This ratio between H_2O_{FGGA} and H_2O_{IRGA} decreased with increasing relative humidity (see Fig. 6). For the comparison, the H_2O measurements from the FGGA were converted from dry mole fractions to mmol m^{-3} , the unit recorded by the IRGA. Points were excluded from the analysis if the H_2O concentration measured by the two instruments differed by more than 20 % (7 out of 933 points). A non-linear model fits the medians of the 5 % relative humidity bin very well. The

relationship yields:

$$\frac{H_2O_{FGGA}}{H_2O_{IRGA}} = 1.072 - 2.528 \times 10^{-3} \times RH_{IRGA} + 1.146 \times 10^{-5} \times RH_{IRGA}^2. \quad (12)$$

4.4 Flux comparison

Putting the previous findings together, the $\overline{w'H_2O'}_{\text{FGGA,lag}(\overline{w'CO_2'})}$ should be scalable to match the atmospheric $\overline{w'H_2O'}_{\text{IRGA}}$, e.g. measured by the LI-7500. Therefore, $\overline{w'H_2O'}_{\text{FGGA,lag}(\overline{w'CO_2'})}$ was post-calibrated according to Eq. (9) and the damping loss correction applied with the damping factor calculated from the relative humidity (Eq. 11). Compared to the raw flux, the corrected $\overline{w'H_2O'}_{\text{FGGA,lag}(\overline{w'CO_2'})}$ compares better with the $\overline{w'H_2O'}_{\text{IRGA}}$. In addition to the post-calibration of the water concentration and the application of the damping correction, the flux was also corrected for the observed signal reduction with increasing relative humidity (Eq. 12). The combination of all corrections reduced the difference between $\overline{w'H_2O'}_{\text{FGGA,lag}(\overline{w'CO_2'})}$ and $\overline{w'H_2O'}_{\text{IRGA}}$ considerably, where the latter contributed 5 % to the total improvement of 47 % (see Fig. 7).

15 4.5 Adapted WPL correction

The information about the damping of $\overline{w'H_2O'}$ by the instrumental setup and the cross-sensitivity of other trace gas measurements to H_2O provides the basis for the adapted WPL correction. This modified correction can be applied to fluxes measured by the FMA where H_2O is not measured. The heat term of the original WPL correction (see Eq. 2) is negligible for the FMA since it is a closed-path instrument (Leuning, 2005). The adapted WPL correction yields:

$$F = a_{\text{slope}} F_{\text{raw}} + \mu[(a_{\text{intercept}} + a_{\text{slope}} \overline{\rho_c}) / \overline{\rho_a}] \frac{bE}{de} \quad (13)$$

The factors a_{slope} and $a_{\text{intercept}}$ apply the post-calibration of the concentration measurements as described in Sect. 4.1, Eq. (8). Since means are removed when calculating fluxes, the intercept can be omitted for the flux, but must be included in the post-calibration of the density of the trace gas.

The additional water dependent cross-sensitivities affect the measurements in a similar way as the dilution effect and can be expressed as an amplification of the dilution effect. The linear correction approach (Eq. 10 without term $b \times \text{H}_2\text{O}^2$, see Table 2) implies that, in case of CH_4 , the dilution effect is amplified by 18 %. Therefore, the water vapor flux E must be multiplied with 1.18, denoted as factor b in Eq. (13).

Since $\overline{w'H_2O'}$ observed in the measurement cell controls the density flux effects caused by H_2O , the atmospheric $\overline{w'H_2O'}$ needs to be modified such that it fits the observed flux in the cell. Therefore, the corrections derived for $\overline{w'H_2O'}$ $_{FGGA,lag(\overline{w'CO_2})}$ were reverted and applied to the atmospheric $\overline{w'H_2O'}$ to reconstruct the conditions in the cell. The damping factor L was calculated according to Eq. (11) and the damping effect assessed according to Eqs. (4) or (5), depending on the stability of the atmosphere. E in Eq. (13) was divided by the calculated damping effect d (Eq. 11) and by e that accounts for the reduced H_2O concentration in the sampling cell (Eq. 12) to simulate $\overline{w'H_2O'}$ $_{FGGA,lag(\overline{w'CO_2})}$.

The damping of the water vapor flux reduces the correction by about 20 % compared to the original correction (see Fig. 8). However, this reduction is compensated by the spectral broadening effect due to H₂O. The final result compares well with the flux calculated from dry mole fractions. In total, the clear CH₄ uptake is reduced to a near zero flux after applying all corrections.

5 Discussion

²⁵ The manufacturer's calibration of the FGGA and FMA agreed very well with the measurements against CH₄ calibration standards as also noted by Detto et al. (2011).

They found very small instrumental drifts for their FMA and FGGA and suggest that no post-calibration is needed for eddy covariance applications. However, CH₄ and CO₂ measurements depend on H₂O. Additional to the dilution effect, a cross-sensitivity on H₂O was found. This occurrence was already observed for a FMA (Tuzson et al., 2010) and other laser spectrometers of various brands (Chen et al., 2010; Tuzson et al., 2010; Nefftel et al., 2010; McDermitt et al., 2010). Such an interference is no problem as long as H₂O is measured simultaneously in the same measurement cell as the other trace gas. The effect can be quantified in a dilution experiment as described in Sect. 3.2. Picarro Inc. and Licor Inc. published the factors for their instruments Picarro G1301 and Licor Li-7700 (Rella, 2010; McDermitt et al., 2010). Meanwhile, new instruments by Los Gatos Research Inc. already correct for cross-sensitivity effects (FMA with serial numbers starting from 11-0049 and above and FGGA starting from 11-0018) and hence no additional correction is needed. FMA and FGGA with serial numbers 10-0150 and below do not include water correction, but the instrument can be upgraded by the manufacturer. Instruments with serial numbers inbetween might correct or are user upgradable and Los Gatos Research Inc. will provide instrument specific information upon request (Robert Provencal, personal communication, 13 September 2011). Instruments that do not measure H₂O perform best if additional measurements of $w'H_2O'$ are available to correct for dilution and cross-sensitivity effects of H₂O. The needed dilution factors can be derived from a dilution experiment added to the traditional WPL correction. Another option would be to dry the air sample before the analysis which might however introduce negative effects to the frequency response and high-frequency performance of the sensor (Griffis et al., 2008).

Further, the conditions in the sampling cell are decisive for the WPL correction. The water vapor flux is damped by the measurement system and hence is reduced in the cell. Therefore, atmospheric water vapor flux would overcorrect the dilution effect. The damping strength of the system depends on the tubing, filters, etc. and is therefore different for individual setups. Additionally, environmental variables such as wind speed and atmospheric stability influence the damping strength (see Sect. 3.3.4). For water

369

vapor, the damping also depends on relative humidity (Mammarella et al., 2009; Ibrom et al., 2007a,b). Smeets et al. (2009) already adapted the WPL correction following the suggestions by Ibrom et al. (2007b) for this issue. They used a FMA that does not measure H₂O and thus had no possibility to validate their assumption. Moreover, they did not address the additional cross-sensitivity effects of CH₄ humid mole fractions on H₂O. Detto et al. (2011) stated that they found no difference between the flux calculated from dry mole fractions and the flux corrected with the water vapor flux derived from an external sensor (LI-7500) for their FGGA. This agrees with our findings. Nevertheless, the original WPL term accounting for the dilution of the air sample by H₂O is reduced. This reduction is however compensated by the cross-sensitivity of the measurements to H₂O.

The dependency of the damping on relative humidity can be at least partially attributed to the likeliness that the tube walls and filters on the way to the cell absorb and desorb water molecules. For CO₂ and CH₄, no dependency was found and the system seems to be inert to these gases. Compared to other trace gases, the sorption processes of water prolong the travelling time of H₂O into the sampling cell. The different time lag introduces a phase shift between H₂O and e.g. CO₂ which leads to an additional damping of $w'H_2O'_{FGGA,lag(w'CO_2)}$ at the time lag of $w'CO_2'$.

We additionally observed a reduction of measured water vapor concentration with increasing relative humidity. The used particle filters are designed to retain liquid water, e.g. if the instrument is operated under foggy conditions. Unfortunately, we found that the filter material itself is not hydrophobic enough and thus introduces a memory effect in the H₂O signal. If the humidity of the atmosphere increases, the filter retains some water, even though the air is not saturated. Reversely, water in the filter reevaporates if the air dries. This might be a reason for the observed reduction in H₂O concentrations measured by the FGGA.

370

6 Conclusions

Our results show that the damping of the water vapor signal as well as cross-sensitivity effects to water affect the measured trace gas flux. While the damping reduces $w'H_2O'$ observed in the cell, the cross-sensitivity effects counteract the damping. For our setup, the two processes coincidentally compensated each other. Hence, the original WPL correction would lead to the same result as our adapted approach which accounts for the individual processes that affect the water vapor flux. However, the damping of $w'H_2O'$ is system specific and depends on deployed filters etc. Therefore, our approach is not generally valid, but provides at least an estimate for the order of magnitude of the described effects.

Acknowledgements. This project was funded by the Competence Center Environment and Sustainability at the ETH Domain (CCES). We thank Hans-Ruedi Wettstein, manager of the agricultural research station Chamau for access to the field site and Sebastian Wolf for his valuable comments on the manuscript.

References

- Baer, D., Paul, J., Gupta, M., and O'Keefe, A.: Sensitive absorption measurements in the near-infrared region using off-axis integrated-cavity-output spectroscopy, *Appl. Phys. B*, 75, 261–265, 2002. 353
- Baldocchi, D. D.: Assessing the eddy covariance technique for evaluating carbon dioxide exchange rates of ecosystems: past, present and future, *Glob. Change Biol.*, 9, 479–492, 2003. 352, 360
- Chen, H., Winderlich, J., Gerbig, C., Hofer, A., Rella, C. W., Crosson, E. R., Van Pelt, A. D., Steinbach, J., Kolle, O., Beck, V., Daube, B. C., Gottlieb, E. W., Chow, V. Y., Santoni, G. W., and Wofsy, S. C.: High-accuracy continuous airborne measurements of greenhouse gases (CO_2 and CH_4) using the cavity ring-down spectroscopy (CRDS) technique, *Atmos. Meas. Tech.*, 3, 375–386, doi:10.5194/amt-3-375-2010, 2010. 369

- Crosson, E.: A cavity ring-down analyzer for measuring atmospheric levels of methane, carbon dioxide, and water vapor, *Appl. Phys. B*, 92, 403–408, 2008. 353
- Detto, M., Verfaillie, J., Anderson, F., Xu, L., and Baldocchi, D.: Comparing laser-based open- and closed-path gas analyzers to measure methane fluxes using the eddy covariance method, *Agr. Forest Meteorol.*, 151, 1312–1324, doi:10.1016/j.agrformet.2011.05.014, 2011. 356, 368, 370
- Eugster, W. and Plüss, P.: A fault-tolerant eddy covariance system for measuring CH_4 fluxes, *Agr. Forest Meteorol.*, 150, 841–851, doi:10.1016/j.agrformet.2009.12.008, 2010. 353, 359
- Eugster, W. and Senn, W.: A cospectral correction model for measurement of turbulent NO_2 flux, *Bound.-Lay. Meteorol.*, 74, 321–340, 1995. 361, 362, 365
- Eugster, W., Zeyer, K., Zeeman, M., Michna, P., Zingg, A., Buchmann, N., and Emmenegger, L.: Methodical study of nitrous oxide eddy covariance measurements using quantum cascade laser spectrometry over a Swiss forest, *Biogeosciences*, 4, 927–939, doi:10.5194/bg-4-927-2007, 2007. 353
- Griffis, T. J., Sargent, S. D., Baker, J. M., Lee, X., Tanner, B. D., Greene, J., Swiatek, E., and Billmark, K.: Direct measurement of biosphere-atmosphere isotopic CO_2 exchange using the eddy covariance techniques, *J. Geophys. Res.*, 113, D08304, doi:10.1029/2007JD009297, 2008. 353, 369
- Hendriks, D. M. D., Dolman, A. J., van der Molen, M. K., and van Huissteden, J.: A compact and stable eddy covariance set-up for methane measurements using off-axis integrated cavity output spectroscopy, *Atmos. Chem. Phys.*, 8, 431–443, doi:10.5194/acp-8-431-2008, 2008. 353
- Horst, T. W.: A simple formula for attenuation of eddy fluxes measured with first-order-response scalar sensors, *Bound.-Lay. Meteorol.*, 82, 219–233, 1997. 361
- Ibrom, A., Dellwik, E., Flyvbjerg, H., Jensen, N. O., and Pilegaard, K.: Strong low-pass filtering effects on water vapour flux measurements with closed-path eddy correlation systems, *Agr. Forest Meteorol.*, 147, 140–156, 2007a. 354, 357, 362, 370
- Ibrom, A., Dellwik, E., Larsen, S. E., and Pilegaard, K.: On the use of the Webb-Pearman-Leuning theory for closed-path eddy correlation measurements, *Tellus B*, 59, 937–946, 2007b. 355, 357, 370
- Kaimal, J. C. and Gaynor, J.: The Boulder Atmospheric Observatory, *J. Climate Appl. Meteor.*, 22, 863–880, 1983. 361
- Kaimal, J. C., Wyngaard, J. C., Izumi, Y., and Coté, O. R.: Spectral characteristics of surface-

- layer turbulence, *Quart. J. Roy. Meteorol. Soc.*, 98, 563–589, doi:10.1002/qj.49709841707, 1972. 361
- Kroon, P. S., Hensen, A., Jonker, H. J. J., Zahniser, M. S., van 't Veen, W. H., and Vermeulen, A. T.: Suitability of quantum cascade laser spectroscopy for CH₄ and N₂O eddy covariance flux measurements, *Biogeosciences*, 4, 715–728, doi:10.5194/bg-4-715-2007, 2007. 353
- 5 Leuning, R.: Measurements of trace gas fluxes in the atmosphere using eddy covariance: WPL corrections revisited, in: *Handbook of Micrometeorology*, edited by: Lee, X., Massman, W., and Law, B., Vol. 29 of *Atmospheric and Oceanographic Sciences Library*, 119–132, Springer Netherlands, doi:10.1007/1-4020-2265-4_6, 2005. 367
- 10 Leuning, R. and Judd, M. J.: The relative merits of open- and closed-path analysers for measurement of eddy fluxes, *Glob. Change Biol.*, 2, 241–253, doi:10.1111/j.1365-2486.1996.tb00076.x, 1996. 356
- Mammarella, I., Launiainen, S., Gronholm, T., Keronen, P., Pumpanen, J., Rannik, Ü., and Vesala, T.: Relative humidity effect on the high-frequency attenuation of water vapor flux measured by a closed-path eddy covariance system, *J. Atmos. Ocean. Technol.*, 26, 1856–1866, doi:10.1175/2009JTECHA1179.1, 2009. 357, 362, 370
- 15 Massman, W.: Concerning the measurement of atmospheric trace gas fluxes with open- and closed-path eddy covariance system: The WPL terms and spectral attenuation, in: *Handbook of Micrometeorology*, edited by: Lee, X., Massman, W., and Law, B., 29 of *Atmospheric and Oceanographic Sciences Library*, 133–160, Springer, The Netherlands, doi:10.1007/1-4020-2265-4_7, 2005. 357
- 20 Massman, W. and Clement, R.: Uncertainty in eddy covariance flux estimates resulting from spectral attenuation, in: *Handbook of Micrometeorology*, edited by: Lee, X., Massman, W., and Law, B., 29 of *Atmospheric and Oceanographic Sciences Library*, 67–99, Springer, The Netherlands, doi:10.1007/1-4020-2265-4_4, 2005. 357
- 25 Mauder, M., Foken, T., Clement, R., Elbers, J. A., Eugster, W., Grünwald, T., Heusinkveld, B., and Kolle, O.: Quality control of CarboEurope flux data – Part 2: Inter-comparison of eddy-covariance software, *Biogeosciences*, 5, 451–462, doi:10.5194/bg-5-451-2008, 2008. 360
- 30 McDermitt, D., Burba, G., Xu, L., Anderson, T., Komissarov, A., Riensche, B., Schedlbauer, J., Starr, G., Zona, D., Oechel, W., Oberbauer, S., and Hastings, S.: A new low-power, open-path instrument for measuring methane flux by eddy covariance, *Appl. Phys. B*, 1–15–15, 2010. 353, 369

- McMillen, R. T.: An eddy correlation technique with extended applicability to non-simple terrain, *Bound.-Lay. Meteorol.*, 43, 231–245, 1988. 360, 361
- Neftel, A., Ammann, C., Fischer, C., Spirig, C., Conen, F., Emmenegger, L., Tuzson, B., and Wahlen, S.: N₂O exchange over managed grassland: Application of a quantum cascade laser spectrometer for micrometeorological flux measurements, *Agr. Forest Meteorol.*, 150, 775–785, 2010. 353, 369
- 5 Press, W. H.: *Numerical recipes in C*, Cambridge University Press, 1988. 361
- R Development Core Team: R: A language and environment for statistical computing, R Foundation for Statistical Computing, Vienna, Austria, available at: <http://www.R-project.org/>, ISBN 3-900051-07-0, 2010. 359
- 10 Rella, C.: Accurate greenhouse gas measurements in humid gas streams using the Picarro G1301 carbon dioxide/methane/water vapor gas analyzer, Picarro, Inc., Sunnyvale, California, 2010. 353, 364, 369
- Smeets, C. J. P. P., Holzinger, R., Vigano, I., Goldstein, A. H., and Röckmann, T.: Eddy covariance methane measurements at a Ponderosa pine plantation in California, *Atmos. Chem. Phys.*, 9, 8365–8375, doi:10.5194/acp-9-8365-2009, 2009. 353, 357, 370
- 15 Stull, R.: *An introduction to boundary layer meteorology*, Kluwer Academic Pub, 1988. 361
- Tuzson, B., Hiller, R. V., Zeyer, K., Eugster, W., Neftel, A., Ammann, C., and Emmenegger, L.: Field intercomparison of two optical analyzers for CH₄ eddy covariance flux measurements, *Atmos. Meas. Tech.*, 3, 1519–1531, doi:10.5194/amt-3-1519-2010, 2010. 353, 359, 369
- 20 Webb, E., Pearman, G., and Leuning, R.: Correction of flux measurements for density effects due to heat and water vapour transfer, *Quart. J. Roy. Meteorol. Soc.*, 106, 85–100, 1980. 354, 355
- Zeeman, M. J., Hiller, R., Gilgen, A. K., Michna, P., Plüss, P., Buchmann, N., and Eugster, W.: Management and climate impacts on net CO₂ fluxes and carbon budgets of three grasslands along an elevational gradient in Switzerland, *Agr. Forest Meteorol.*, 150, 519–530, doi:10.1016/j.agrformet.2010.01.011, 2010. 358, 359, 361
- 25

Table 1. Fitted parameters for Eq. (10) (CO_2 , CH_4 , and H_2O in ppm).

Parameter	Unit	Values for $\text{CO}_{2,\text{FGGA}}$	Values for $\text{CH}_{4,\text{FGGA}}$	Values for $\text{CH}_{4,\text{FMA}}$
a	$[\text{ppm}^{-1}]$	-1.219×10^{-06} ($\pm 2.169 \times 10^{-09}$)	-1.189×10^{-06} ($\pm 1.480 \times 10^{-09}$)	-1.219×10^{-06} ($\pm 1.182 \times 10^{-09}$)
b	$[\text{ppm}^{-2}]$	1.229×10^{-12} ($\pm 1.073 \times 10^{-13}$)	2.096×10^{-13} ($\pm 7.318 \times 10^{-14}$)	1.678×10^{-12} ($\pm 5.846 \times 10^{-14}$)
R^2		0.9928	0.9966	0.9992

Table 2. Fitted parameters for Eq. (10) without term $b \times \text{H}_2\text{O}^2$ (CO_2 , CH_4 , and H_2O in ppm).

Parameter	Unit	Values for $\text{CO}_{2,\text{FGGA}}$	Values for $\text{CH}_{4,\text{FGGA}}$	Values for $\text{CH}_{4,\text{FMA}}$
a	$[\text{ppm}^{-1}]$	-1.195×10^{-06} ($\pm 5.825 \times 10^{-10}$)	-1.184×10^{-06} ($\pm 3.952 \times 10^{-10}$)	-1.186×10^{-06} ($\pm 3.265 \times 10^{-10}$)
R^2		0.9927	0.9966	0.9991

Fig. 1. Results from the dilution experiment for CO₂ (a) and CH₄ (b). Moisture was added to the dry gas at 9 increasing levels, which were held constant over 20 min. The dilution of the gas explains most of concentration drop (difference between black and dark grey points). However, a small portion cannot be attributed to this process (difference between dark and light grey points) and must have a different source, which is also related to the water vapor concentration. The solid line indicates the mixing ratio of the supplied air.

377

Fig. 2. Damping factor L for $\overline{w'H_2O'}_{FGGA.lag(w'CO_2)}$ related to relative humidity. The grey band shows the interquartile range for 5% relative humidity bins with at least 8 observations. This corresponds to 50 % of all observed values. The fine line connects the medians of the individual bins and the bold line indicates the model fit through the medians.

378

Fig. 3. Histograms of the time lags for $\overline{w'CO_2'}$, $\overline{w'H_2O'}$, and $\overline{w'CH_4'}$ within the search window of 10–40 records (0.5–2 s). In (a), the time lags were determined from humid mole fractions whereas in (b), dry mole fractions were used. Only records where all three fluxes showed a detectable time lag were selected and thus the number of observations is equal for all gases of the individual histogram.

379

Fig. 4. (a): time lags of $\overline{w'H_2O'}$ depend on the relative humidity of the air whereas no such relation was found for $\overline{w'CO_2'}$. (b): the relative difference between the water vapor flux at the time lag of $\overline{w'CO_2'}$ and at the optimal lag maximizing $\overline{w'H_2O'}$.

380

Fig. 5. Cospectrum of $\overline{w'H_2O'}$ for 26 July 2009 11:40–13:25 CET (2^{17} records). The solid line shows the cospectrum of w and H_2O at the time lag maximizing the flux. The dashed line represents the cospectrum of $w'H_2O'$ at the time lag of $\overline{w'CO_2'}$. The total flux is slightly lower with considerable damping in the higher frequencies. The relative humidity of the air was around 54 % ($\approx 640 \text{ mmol m}^{-3}$).

381

Fig. 6. Water vapor concentration ratio between the FGGA and LI-7500 decreases with higher relative humidity of the air. The dots show the individual data points. The band represents the interquartile range for 5% relative humidity bins. The black line indicates the polynomial fit of the second order through the class medians whereas the white line connects the medians of the classes.

382

Fig. 7. Water vapor fluxes of the two systems. $\overline{w'H_2O'}_{IRGA}$ were binned into classes of $1 \text{ mmol m}^{-2} \text{s}^{-1}$. The symbols denote the median values for both variables within the class. The cross hairs span the interquartile range. The dashed line represents the values before correcting $\overline{w'H_2O'}_{FGGA, \text{lag}(\overline{w'CO_2'})}$ for damping losses, whereas the solid line denotes the results thereafter. The dotted line indicates the 1:1 relationship.

383

Fig. 8. Boxplots of the observed raw and corrected CH_4 fluxes. From left to right, the first one shows the raw flux. The second boxplot shows the flux after applying the original WPL term. In the third one, the water vapor flux is damped as it would be observed in the instrument whereas the forth also includes the spectral broadening effect of the instrument. The last boxplot is considered the reference flux: the flux when raw values are converted to dry mole fractions before the calculation of the covariance. The solid line is drawn at zero and the dashed line indicates the expected value, the median of the reference flux.

384